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news on the phase structure of the SMEFT

Gravitational Waves and the Early Universe · KMI Nagoya University | March 12th 2026

Luis Gil (Universidad de Granada)

Based on:

M. Chala, M. C. Fiore and **LG** [2505.14335]



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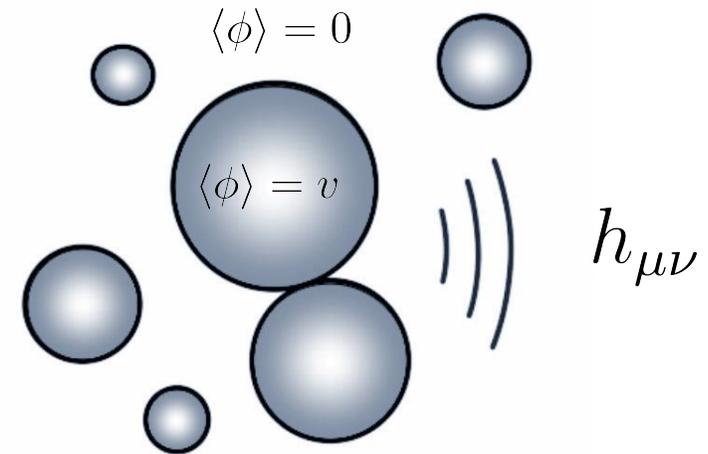
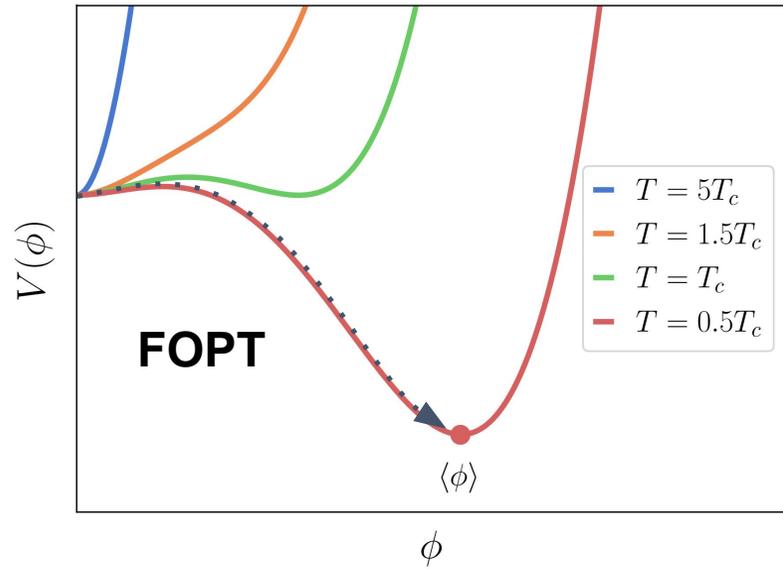


MINISTERIO
DE CIENCIA, INNOVACIÓN
Y UNIVERSIDADES



Introduction

Cosmological phase transitions



Microscopic scales

Thermal field theory
3D EFT
Phase transition (PT) parameters

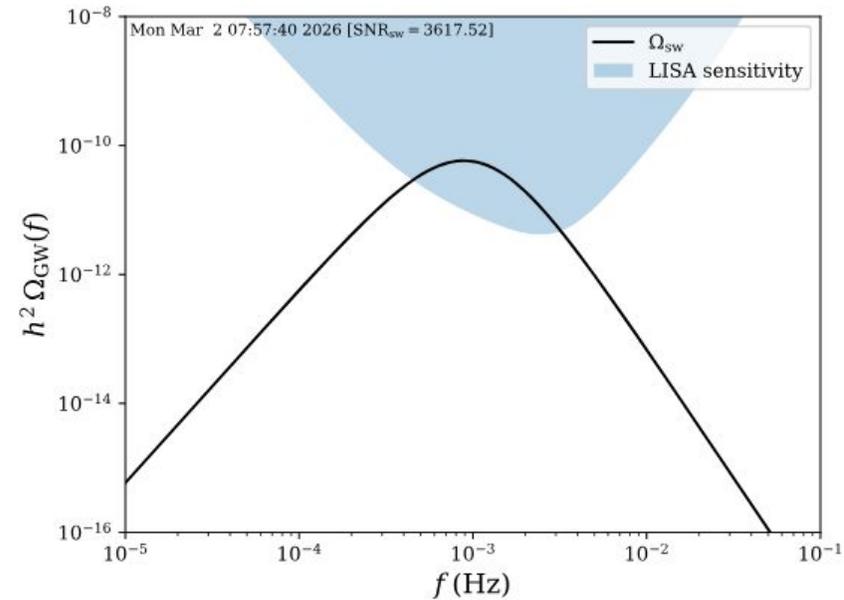
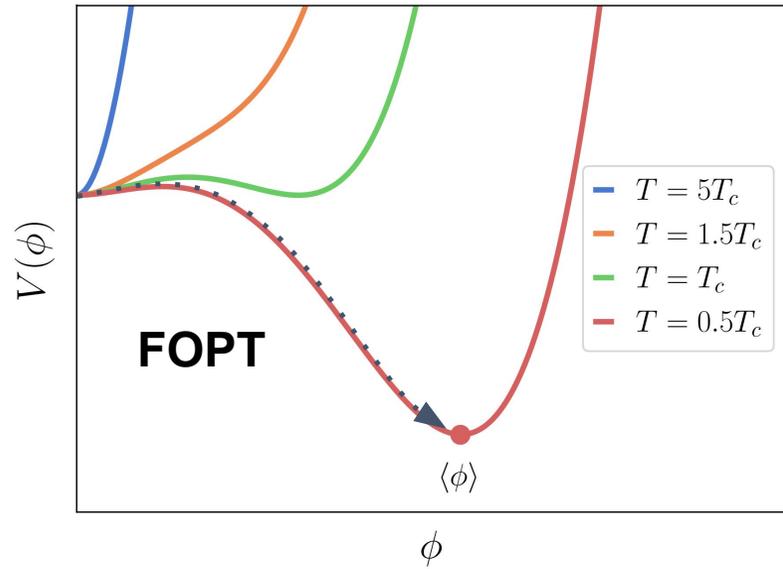
Meso- & macroscopic scales

Cosmological simulations
Bubble wall dynamics
Particle and GW production



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Cosmological phase transitions

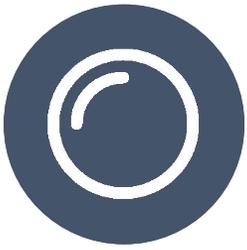


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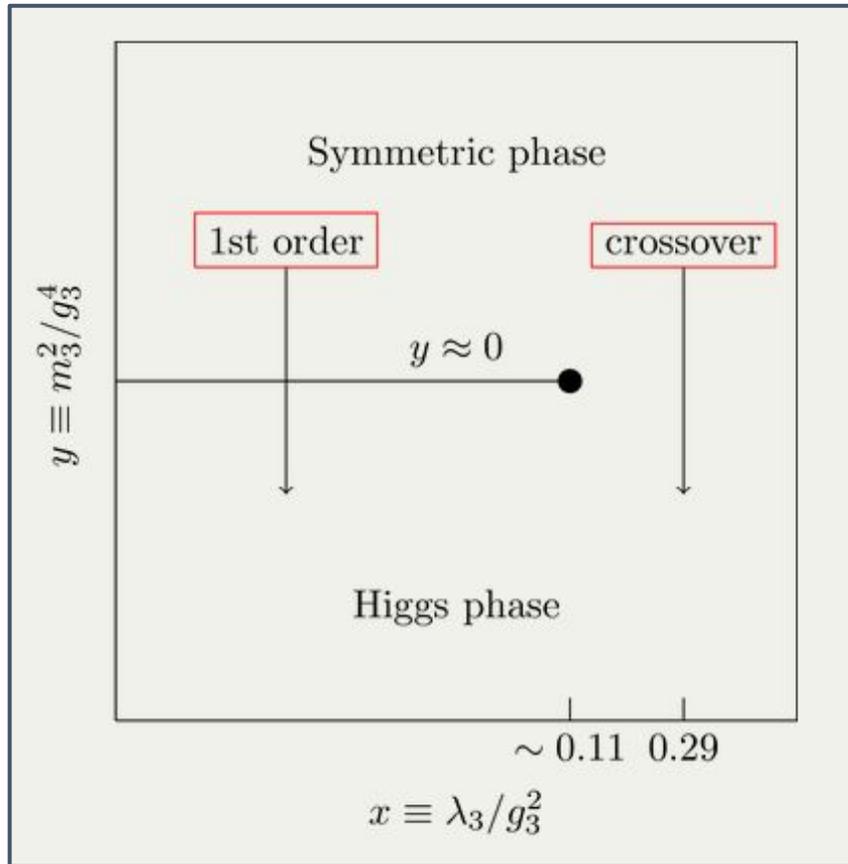
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Introduction

Higgs phase diagram

[Kajantie et al. - hep-ph/9605288]
[Gürtler et al. - hep-lat/9704013]
[Csikor et al. - hep-ph/9809291]

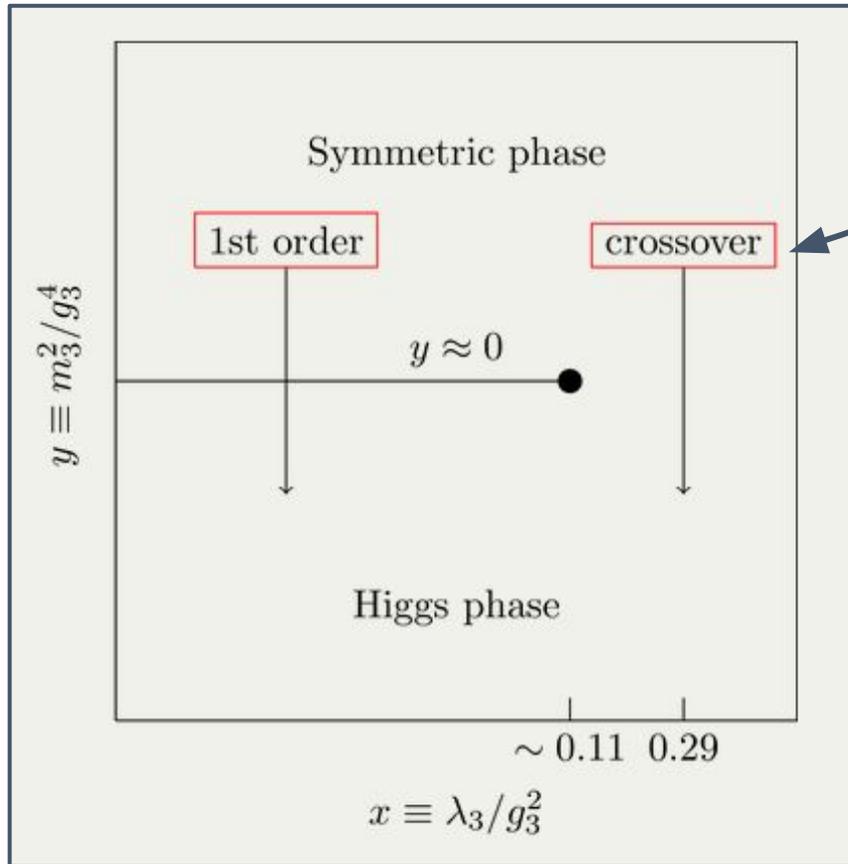




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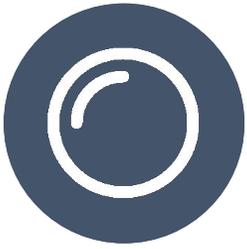
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No electroweak first-order phase transition (FOPT) in **SM** :(

Many unknowns:

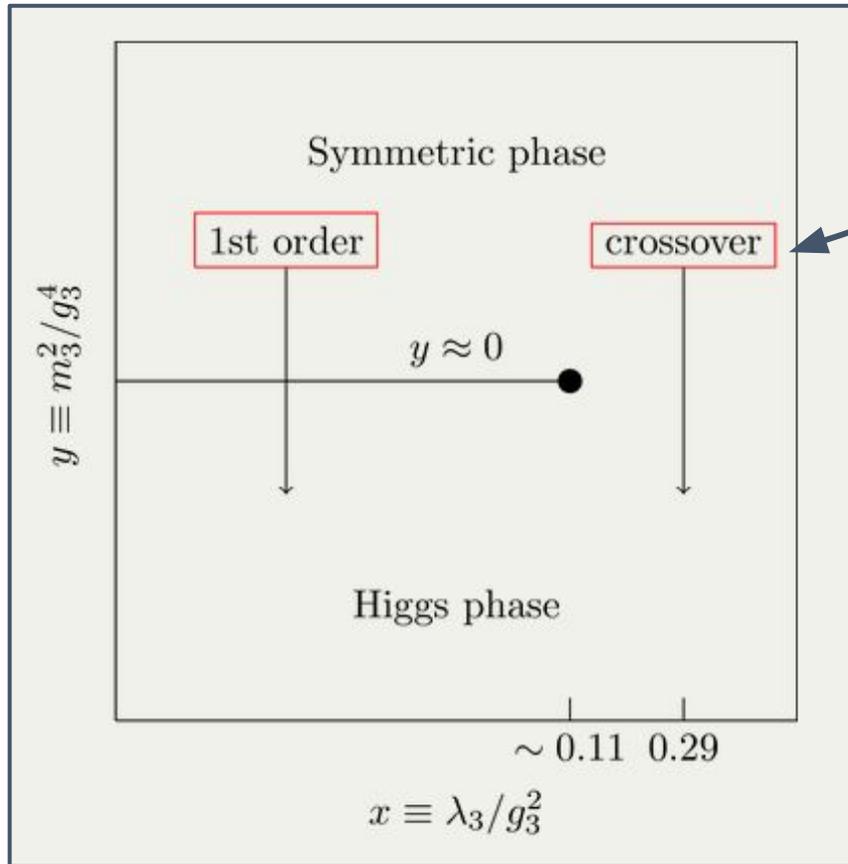
- What then caused baryogenesis?
- What BSM physics can trigger a FOPT?
- Can we expect to tell them apart through GWs?



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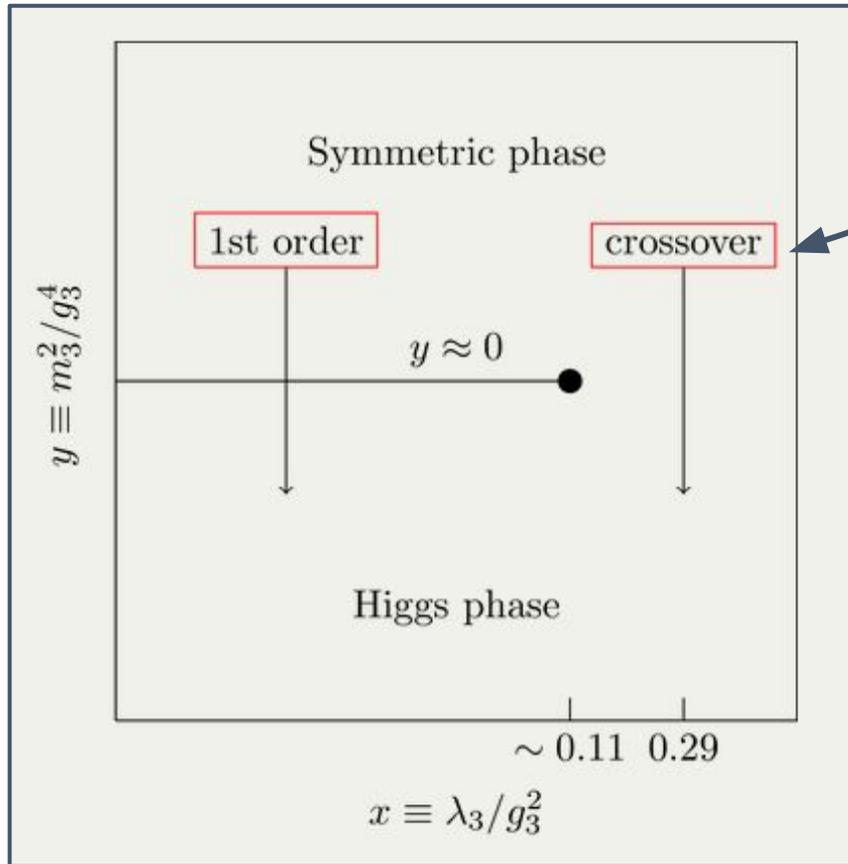
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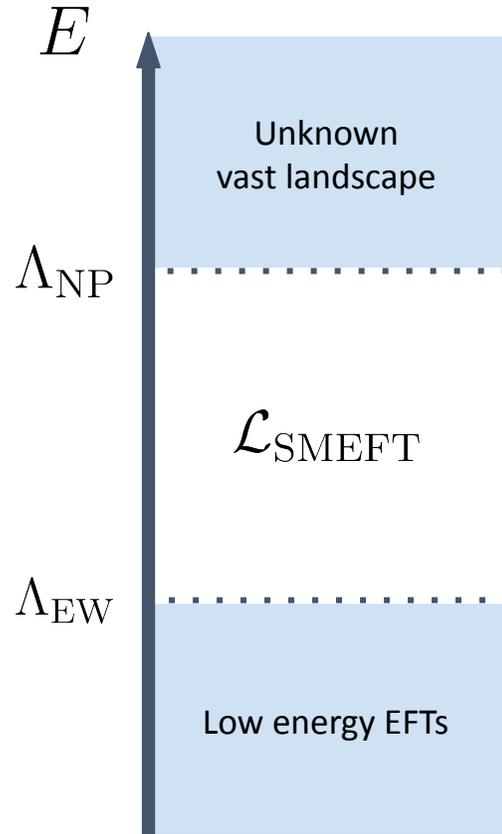
➡ **Let's take a look at the SMEFT!**



Introduction

What is SMEFT?

[Buchmuller, Wyler - NPB 268 (1986) 621-653]



- **SMEFT** captures all heavier-than- Λ_{EW} physics as tower of effective operators

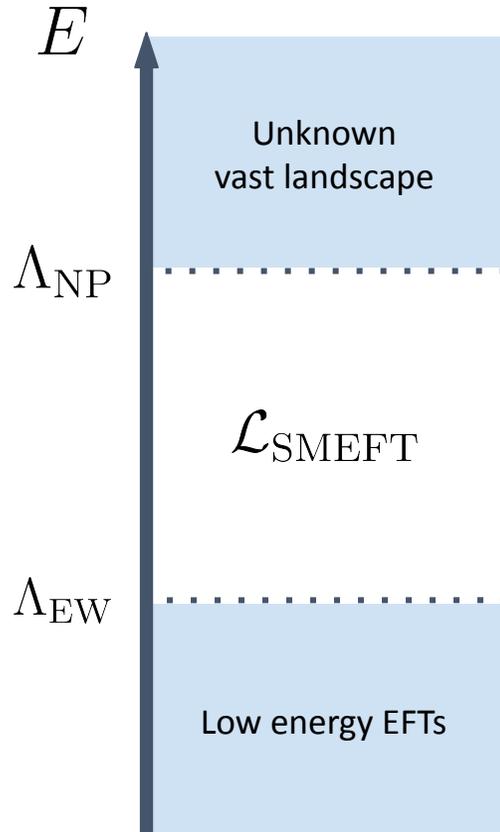
$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{\ell=0}^{\infty} \sum_{n=5}^{\infty} \sum_k \frac{c_{n,k}^{(\ell)}}{(4\pi)^{2\ell} \Lambda^{n-4}} \mathcal{O}_{n,k}$$



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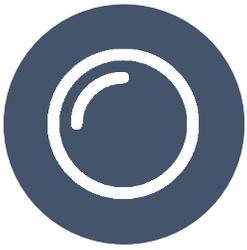
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PROS

- Whole classes of BSM models at once without new fields
- Universal framework to compare with experiments

CONS

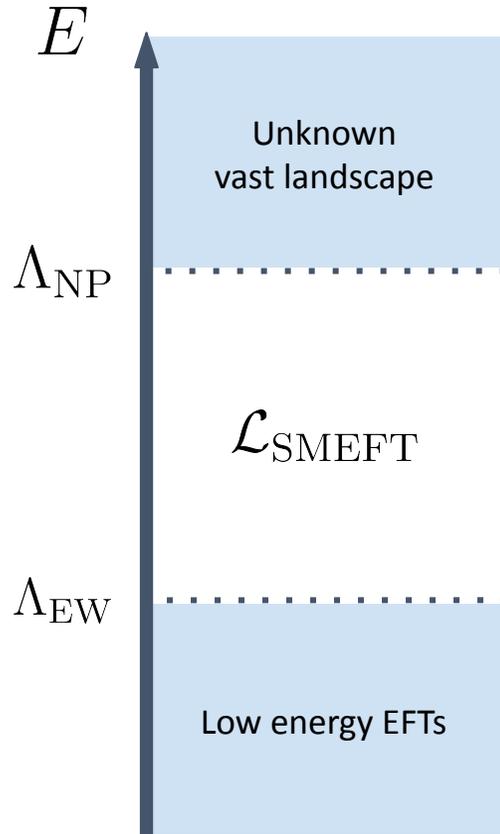
- No lighter-than- Λ_{EW} new physics
- Only Higgs phase transition



Introduction

What is SMEFT?

[Buchmuller, Wyler - NPB 268 (1986) 621-653]



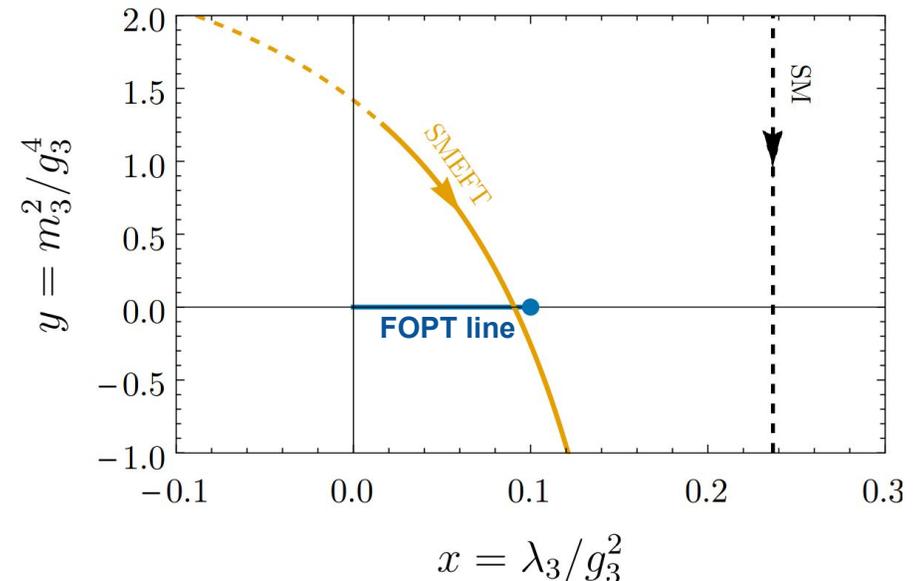
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$$x, y = f(c_i^{(SM)}; T)$$

$$x', y' = f(c_i^{(SMEFT)}; T)$$

↑
Modify how potential evolves with T





High-T limit of SMEFT

The EFT approach

DIMENSIONAL REDUCTION

At $T \sim 100 \text{ GeV}$, a hierarchy of thermal scales allows us to construct a **bosonic 3D EFT** for thermal equilibrium.

SM-like theories reduce to:

$$\mathcal{L}_{3D} = \frac{1}{4} W_{ij}^I W_{ij}^I + (D_i \phi)^\dagger (D_i \phi) + m_3^2 |\phi|^2 + \lambda_3 |\phi|^4$$

where now all couplings are **functions of T**.



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where now all couplings are **functions of T**.

- 1) Assign a **power counting** in powers of gauge coupling g
- 2) **Match** all 3D EFT parameters to $O(g^4)$

$$g_3(T), m_3^2(T), \lambda_3(T)$$

- 3) Compute parameters \mathbf{x} and \mathbf{y}

$$x \equiv \frac{m_3^2}{g_3^4} \quad y \equiv \frac{\lambda_3^2}{g_3^2}$$

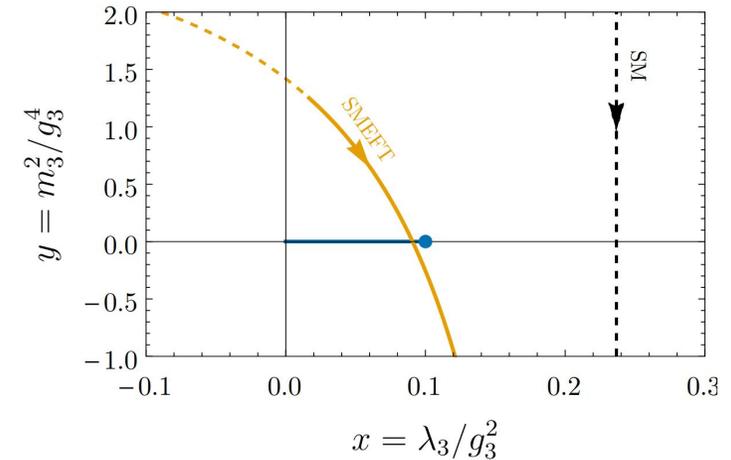
- 4) **Scan** SMEFT parameter space (compatible with colliders) [Giani et al. (SMEFIT) - 2302.06660]



High-T limit of SMEFT

Scanning the SMEFT

- Huge parameter space: **SMEFT @ dim-6** contains **59 independent operators (non B or L violating)**
- Which get us **closer to the FOPT line?**
 - Must make x smaller
 - Must be relatively unconstrained from collider data





High-T limit of SMEFT

Scanning the SMEFT

- Huge parameter space: **SMEFT @ dim-6** contains **59 independent operators (non B or L violating)**
- Which get us **closer to the FOPT line?**
 - Must make x smaller
 - Must be relatively unconstrained from collider data
- We find the following candidates:

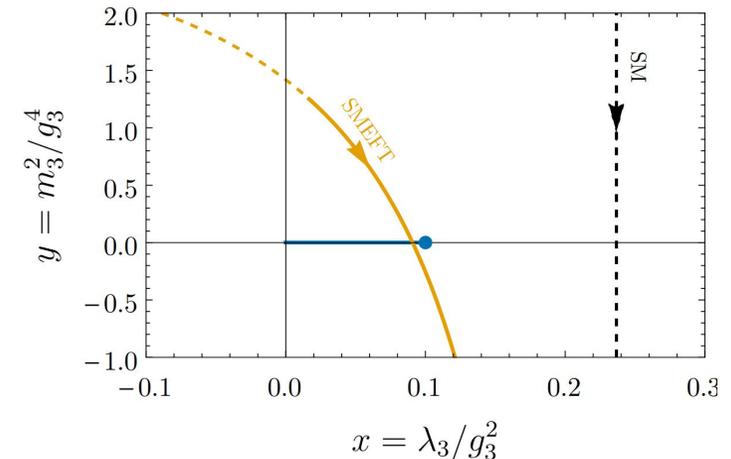
$$\mathcal{O}_\phi \equiv (\phi^\dagger \phi)^3$$

[Camargo-Molina et al. - 2103.14022]

$$\mathcal{O}_{\phi\Box} \equiv (\phi^\dagger \phi)\Box(\phi^\dagger \phi)$$

$$\mathcal{O}_{t\phi} \equiv \bar{t}_L \tilde{\phi} t_R (\phi^\dagger \phi)$$

(Hot) new!

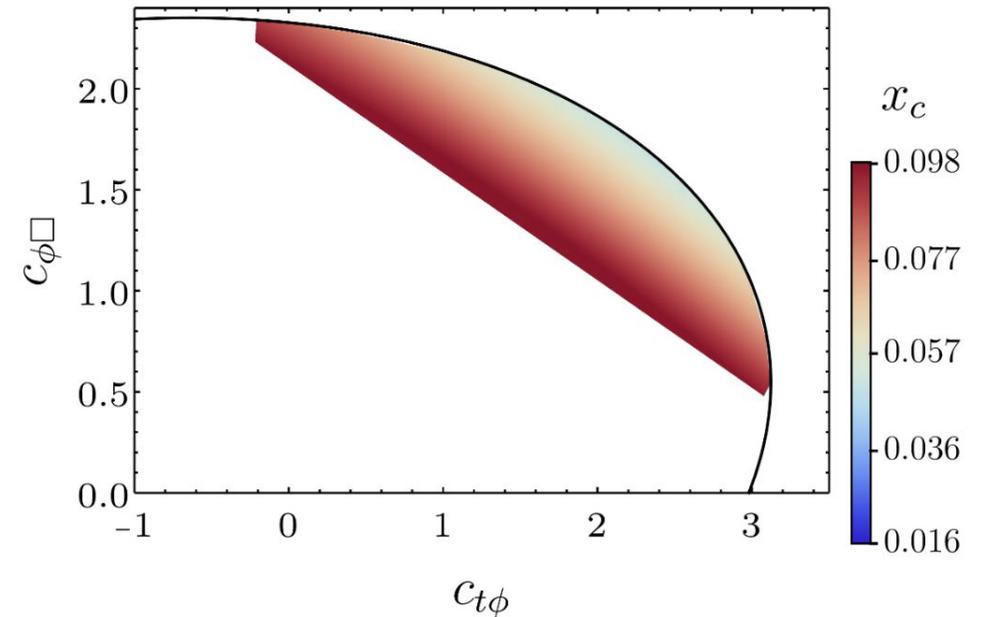
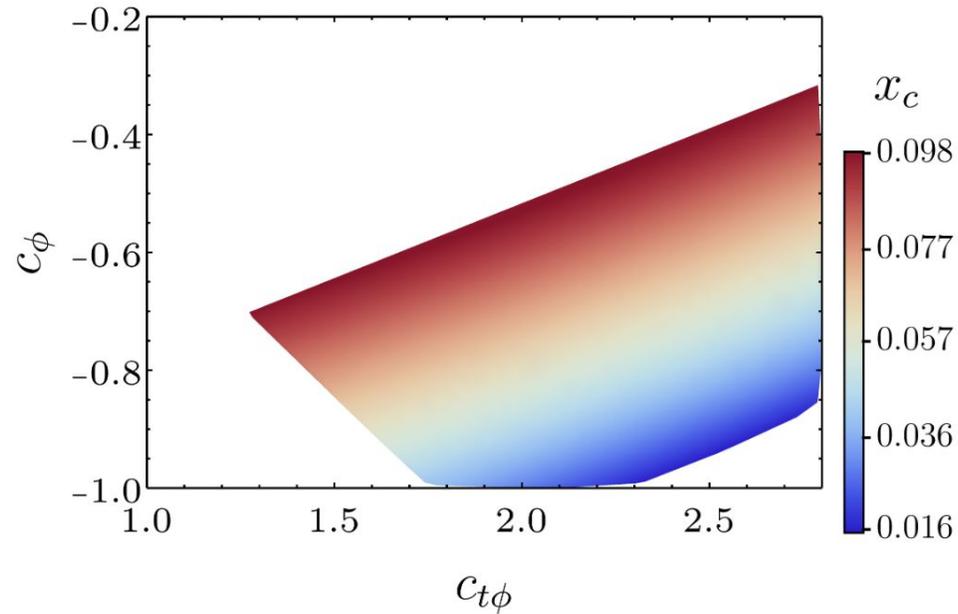




Results

Modified phase diagram

[Chala, Fiore, LG - 2507.16905]



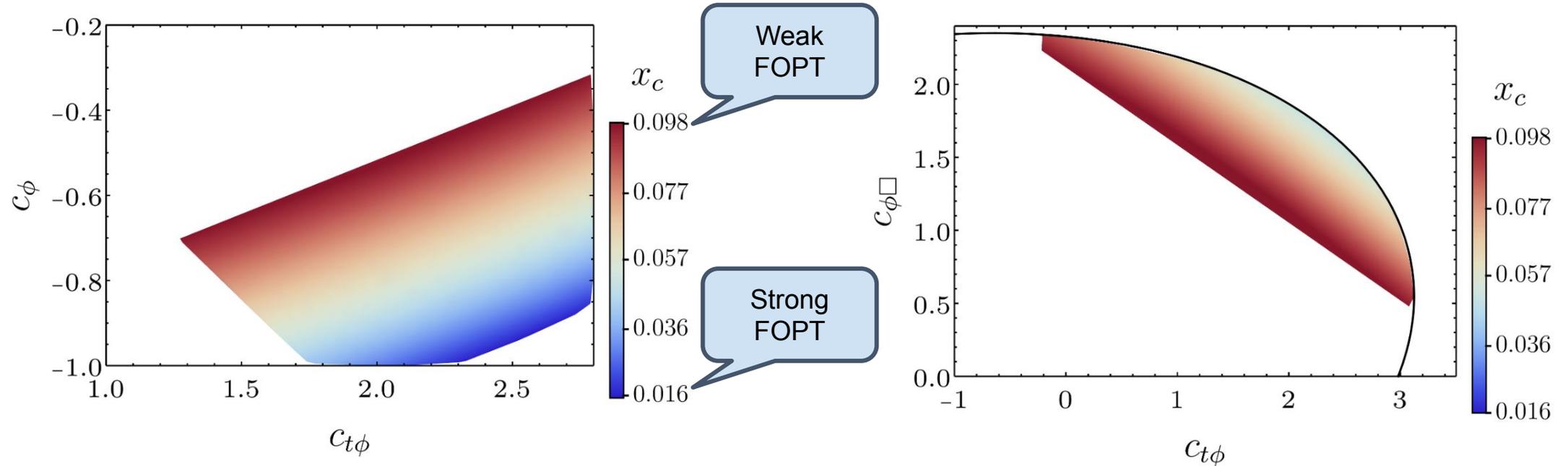
Figures: Regions of the parameter space (in TeV^{-2}) within collider bounds with a FOPT. In both panels, the non-appearing WC is set to zero.



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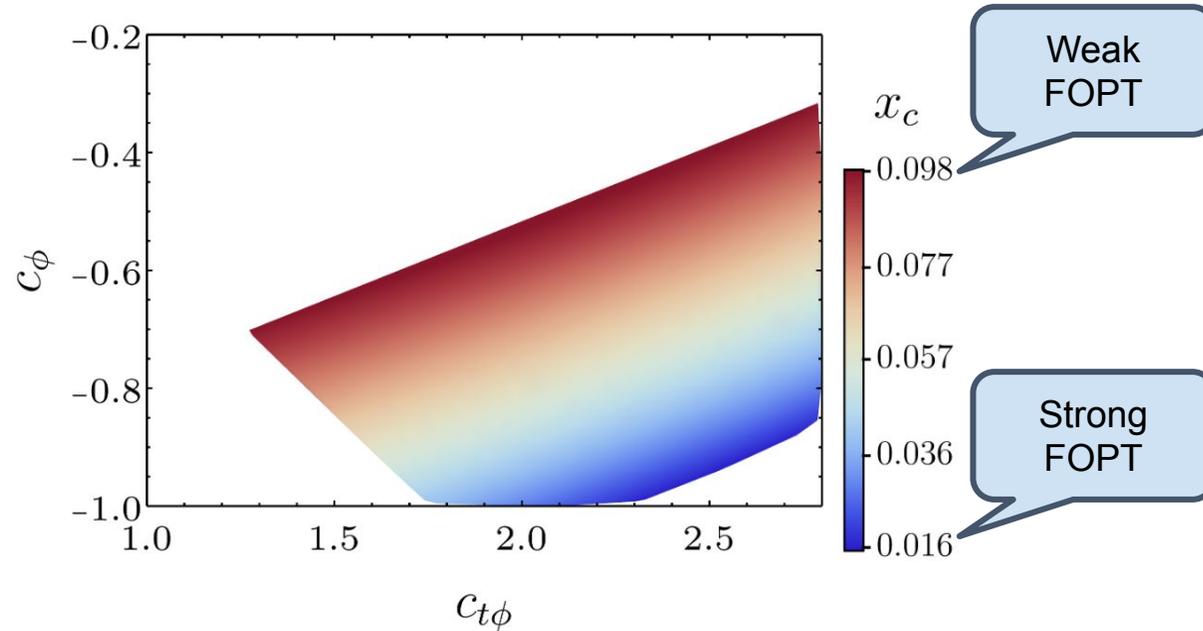
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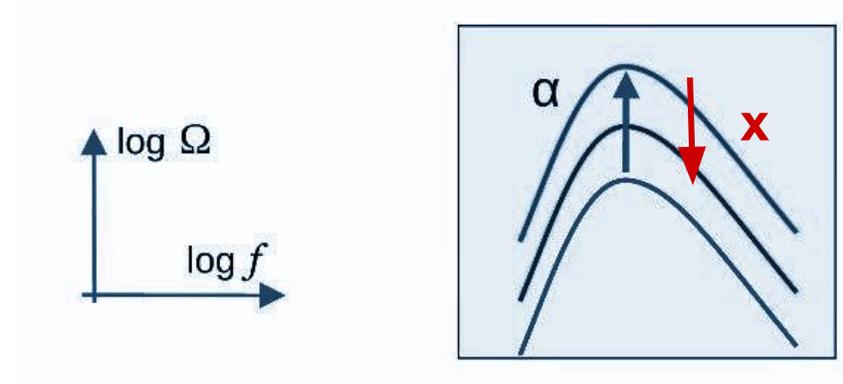
Results

Modified phase diagram

[Chala, Fiore, LG - 2507.16905]



- Small x is correlated with stronger transitions, with larger α (**strength parameter**):



Adapted from K. Rummukainen's talk at CERN August 2025

- Strong (observable via GWs) FOPTs require larger SMEFT coefficients.

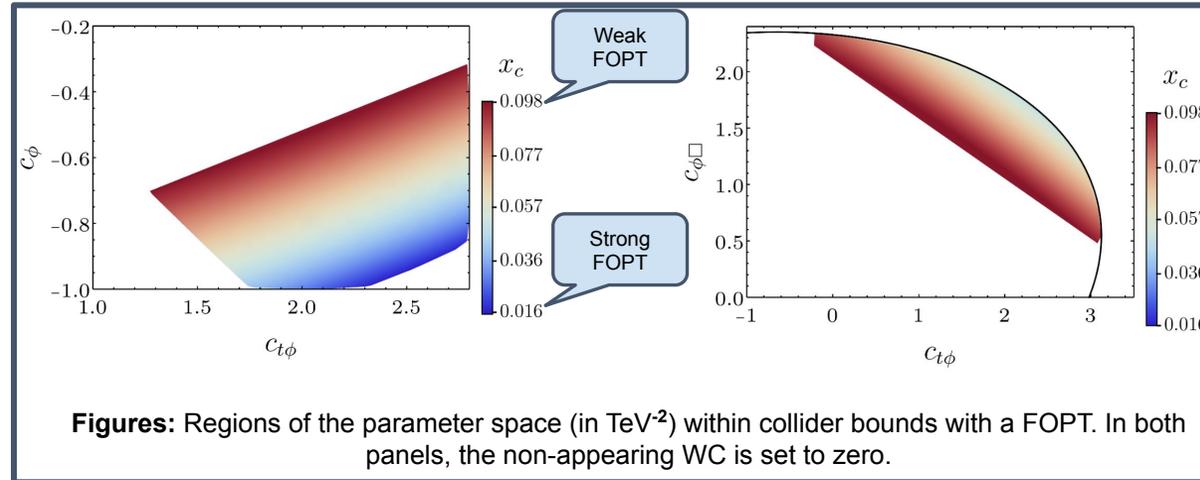
Equivalently, **BSM cannot be too heavy or too weakly-coupled** to produce a strong FOPT, but it might still be captured by SMEFT.



Results

[Chala, Fiore, LG - 2507.16905]

Modified phase diagram



- FOPTs possible without direct modification of the Higgs potential (**with $\mathbf{c}_\phi = 0$**).
- Strong FOPTs **with joint contributions** of the three dim-6 operators, with all of them carrying **at most $\mathbf{O(1) TeV}^{-2}$ coefficients**, which is within collider bounds.
- Non-observation of GWs can then put upper bounds on these operators



Conclusions

- Computed, for the first time, the full $O(g^4)$ high-temperature limit of the SMEFT.
- Explored new directions in its vast parameter space, constrained by collider bounds.
- Found that FOPTs are possible even without the sextic operator, but it is needed for strong FOPTs
 ➔ BSM models that generate $\mathcal{O}_{t\phi} \equiv \bar{t}_L \tilde{\phi} t_R (\phi^\dagger \phi)$ at low energies worth checking too!
- Few works have addressed FOPTs in the SMEFT so far:
 [Camargo-Molina et al '22] [Camargo-Molina et al '24] [Chala, Guedes '25] [Chala, Fiore, LG '25]

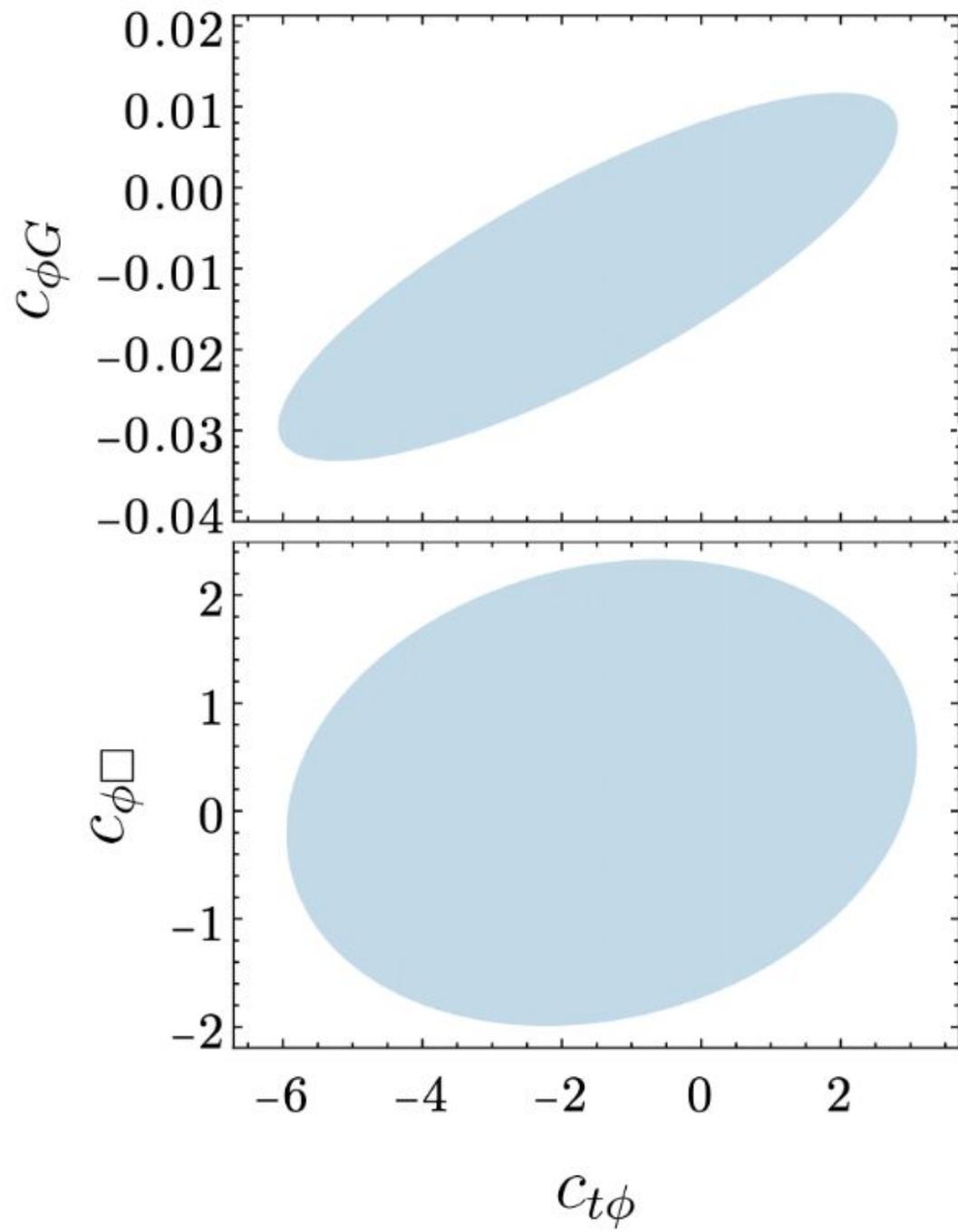
What's left to do? 🔍

Full nucleation study? SMEFT @ dim-8? HEFT? ...

Thank you for your attention!

ご清聴ありがとうございました。

$\mathcal{O}(g)$	$\mathcal{C}X^3$
$\mathcal{O}(g^2)$	$\mathcal{C}X^2\phi^4, \mathcal{C}\phi^4D^2, \mathcal{C}\psi^2X\phi, \mathcal{C}\psi^2\phi^2, \mathcal{C}\psi^4$
$\mathcal{O}(g^3)$	$\mathcal{C}\psi^2\phi^3$
$\mathcal{O}(g^4)$	$\mathcal{C}\phi^6$





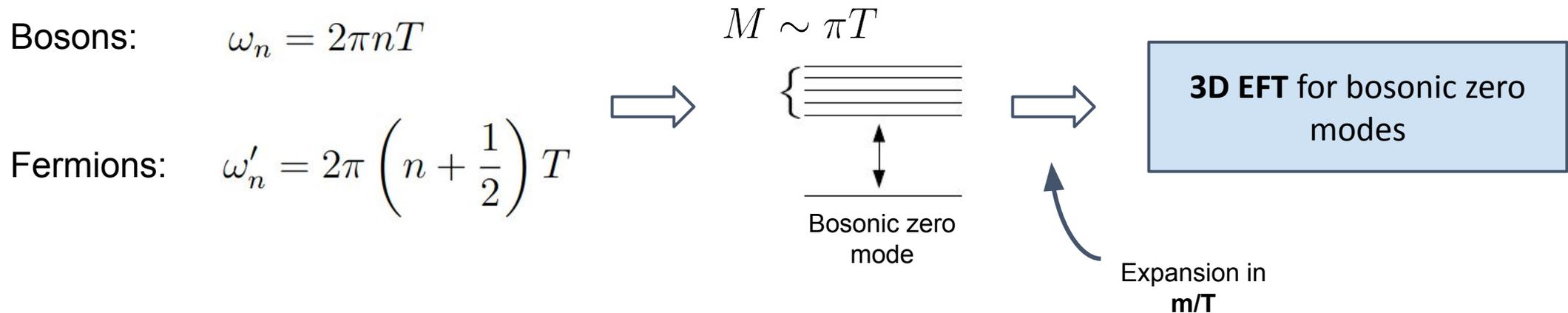
Thermal field theory in equilibrium

Dimensional reduction

- **Generating functional** in thermal field theory (= Euclidean QFT with periodic time):

$$\mathcal{Z}_{\text{th}} = \text{Tr} (e^{-\beta\mathcal{H}}) = \sum_q \langle q \ 0 | e^{-\beta\mathcal{H}} | q \ 0 \rangle = \mathcal{N} \int_{q(0)=q(-i\beta)} \mathcal{D}q \exp(-S_E)$$

- Fields decompose in tower of 3D **Matsubara modes** (~ Kaluza-Klein) with thermal masses:





Thermal field theory in equilibrium

[Gould, Tenkanen - 2309.01672]

Thermal hierarchy of scales

$$\underbrace{\pi T}_{\text{hard}} \gg \underbrace{\left(\frac{g}{4\pi}\right) \pi T}_{\text{soft}} \gg \underbrace{\left(\frac{g}{4\pi}\right)^{3/2} \pi T}_{\text{softer}}$$

Finite-temperature
4D SMEFT

$$\Lambda_{\text{NP}} \gg \pi T$$



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Integrate out heavy
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$$\pi T \gg \mu^2$$



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Finite-temperature
4D SMEFT

Dimensional
reduction

3D matching

Integrate out heavy
Matsubara modes

Integrate out temporal
gauge bosons

$$\Lambda_{\text{NP}} \gg \pi T$$

$$\pi T \gg \mu^2$$

$$m_{W_0}^2 \gg m_\phi^2$$

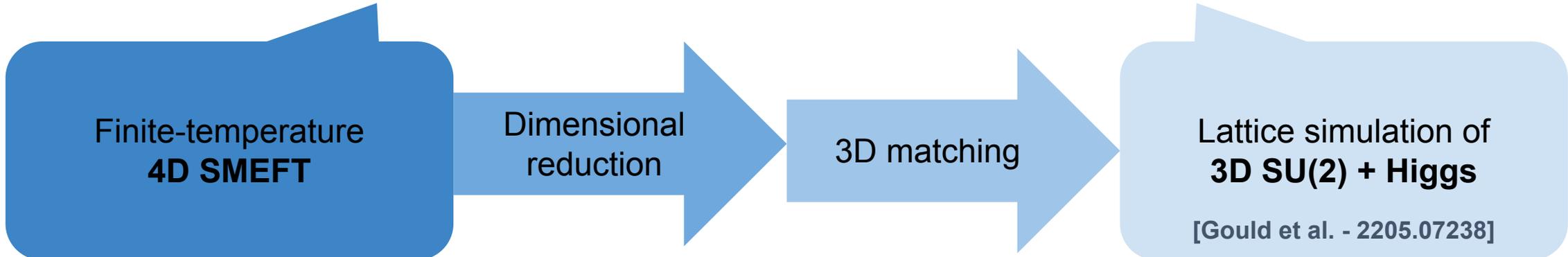


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