



Memory and soft modes in GW signals

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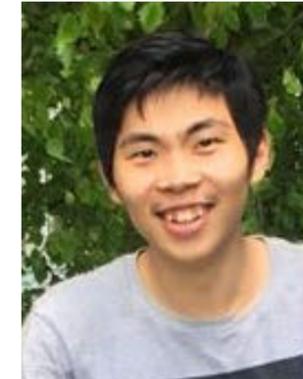
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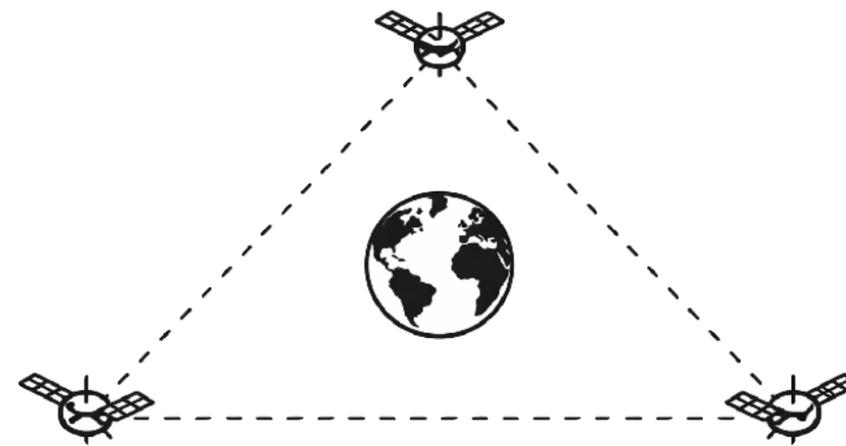
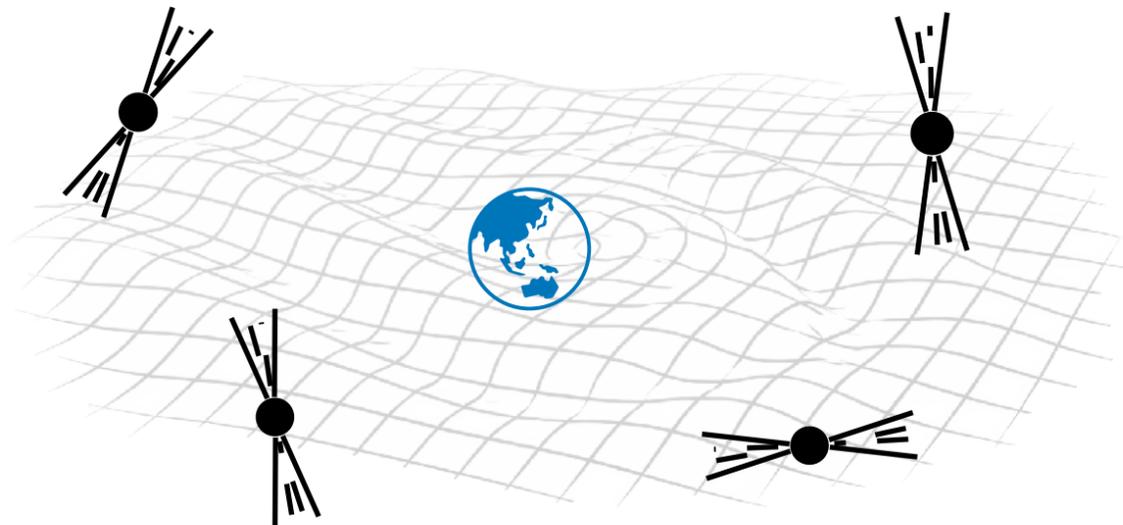
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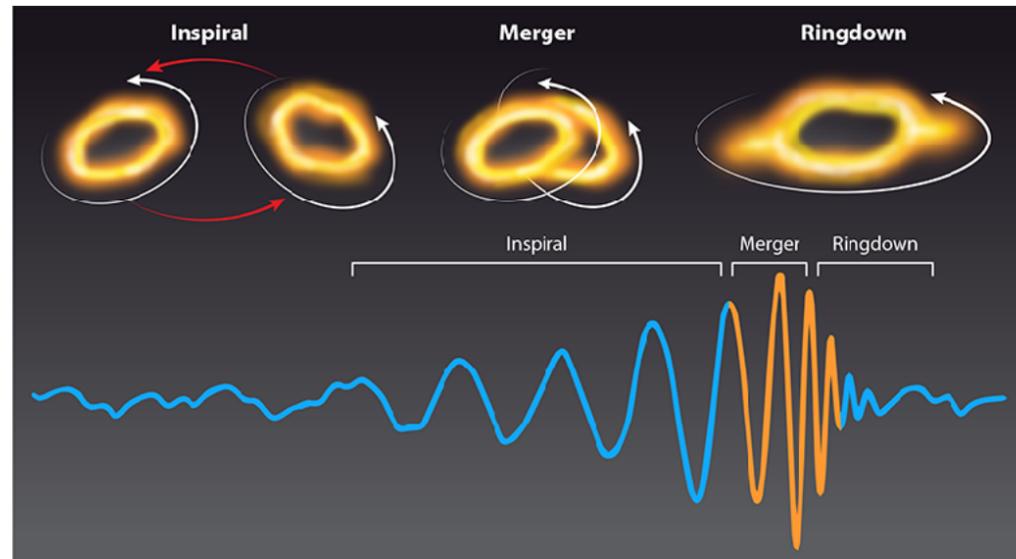
[De Luca, Khoury & **Wong**, 2412.01910]
[De Luca, Khoury & **Wong**, 2412.12273]
& work in progress

[Liang, Lin, Trodden & **Wong**, 2309.16666]



Two types of GW signals

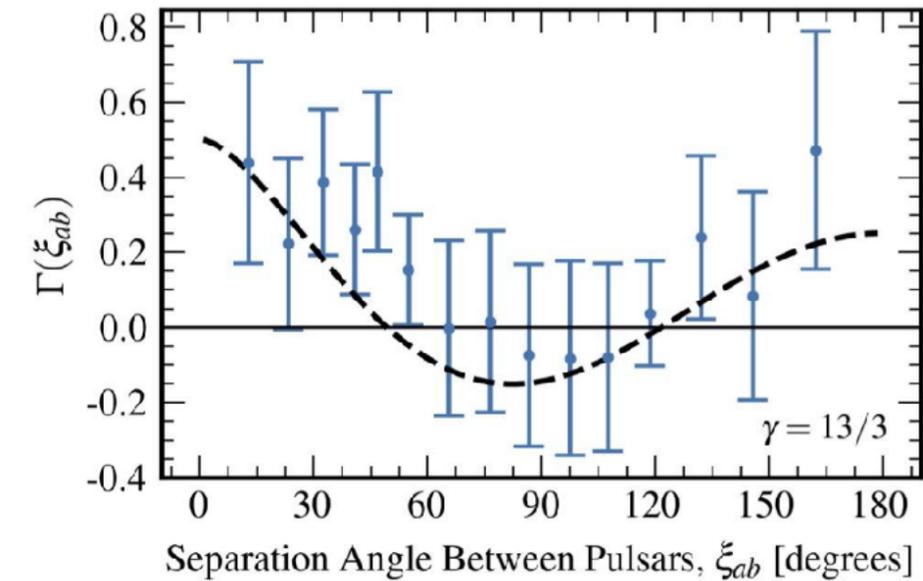
Single event signal



LIGO-VIRGO-KAGRA

Credit: <https://physics.aps.org/articles/v16/29>

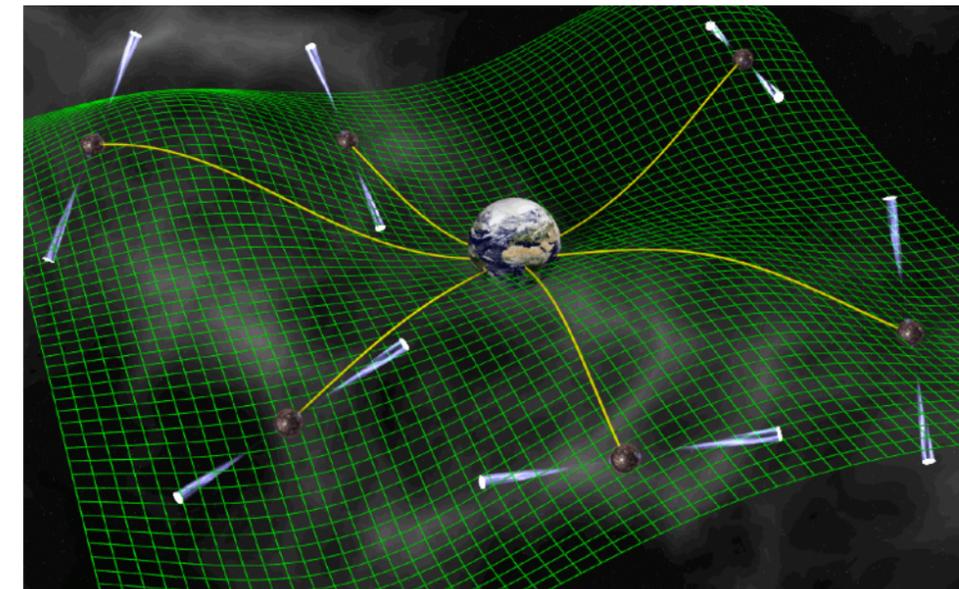
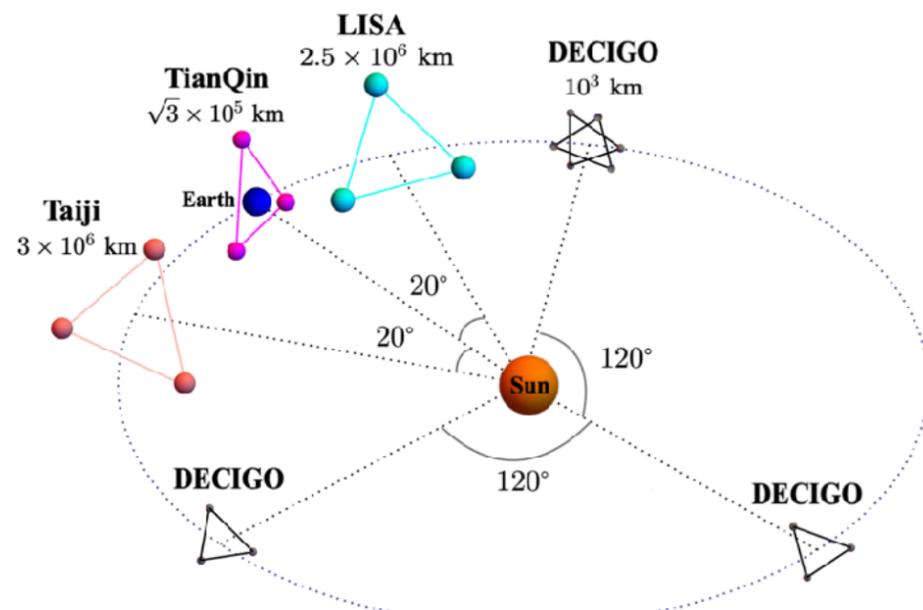
Stochastic



[NANOGrav 15-yr, 2306.16213]

Future: (Ground based) Cosmic Explorer, Einstein telescope

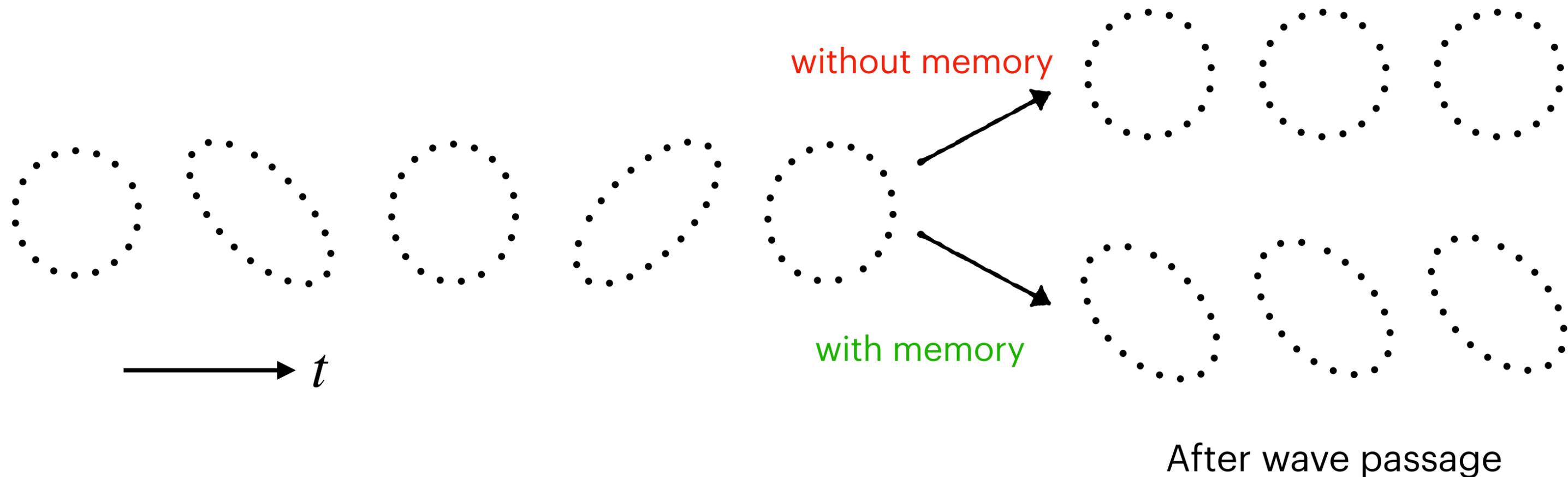
(Space based) LISA, Taiji/Tianqin, DECIGO



Gravitational memory

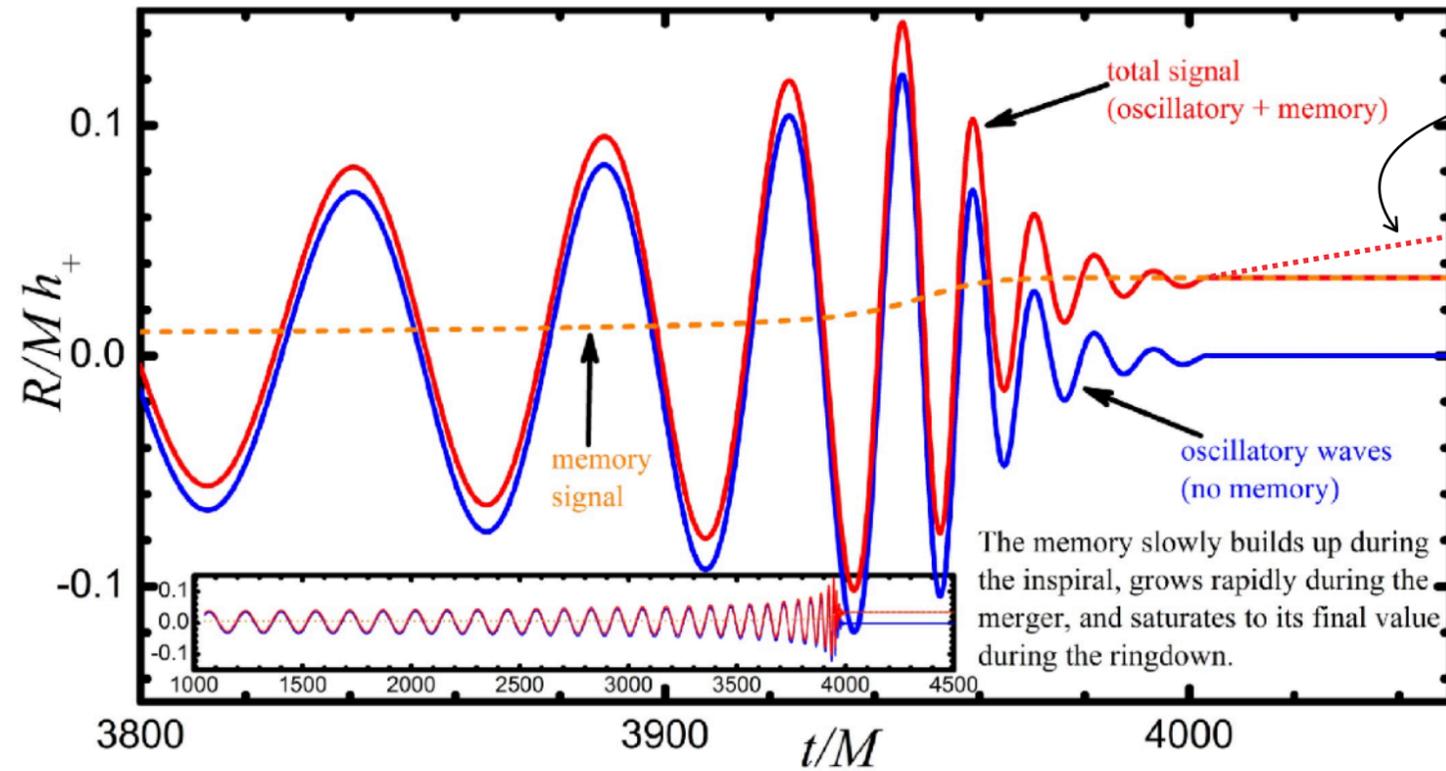
Gravitational memory (soft modes in single event signal)

- Permanent change of spacetime after wave trains have passed through
- Linear memory (caused by explicit matter source), nonlinear memory (gravity itself)
- For a pair of inertial detectors, the relative distance, relative velocity can be altered permanently
- Inertial detectors: Lisa, Tianqin/Taiji, DECIGO

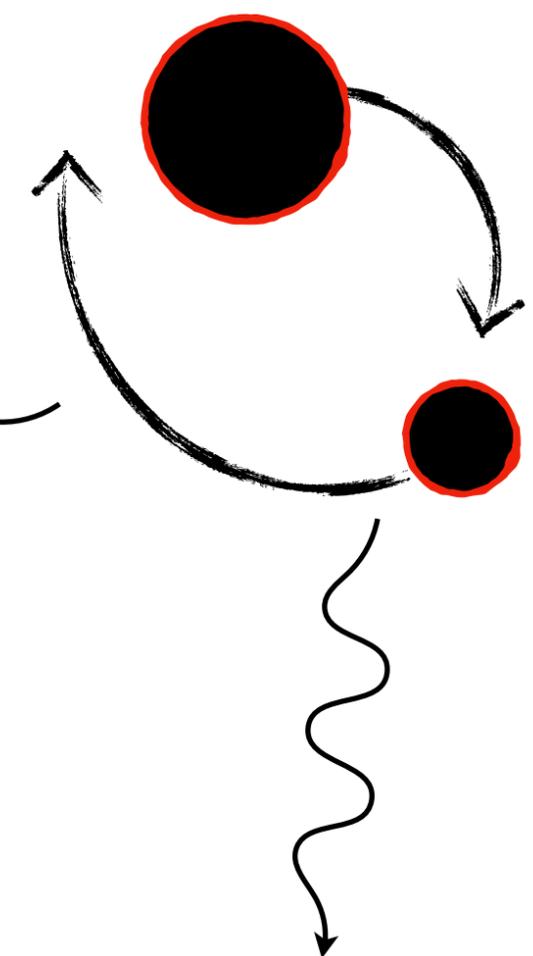
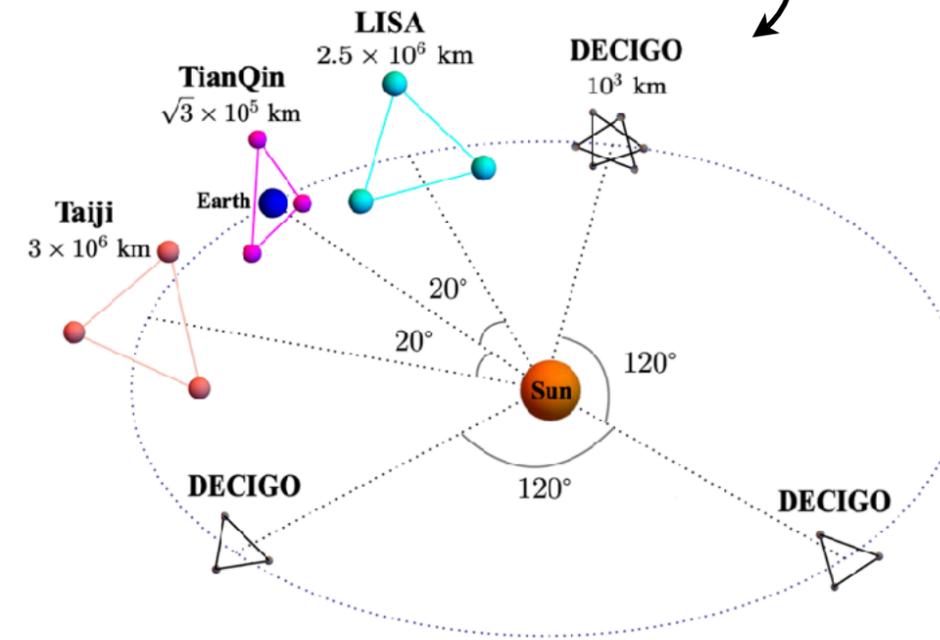


Gravitational meomry

- Examples of sources: black hole merger



Two mirrors flying apart



- Other known examples in simulations:

Supernova simulations and GRBs (due to ejection of mass, linear memory)

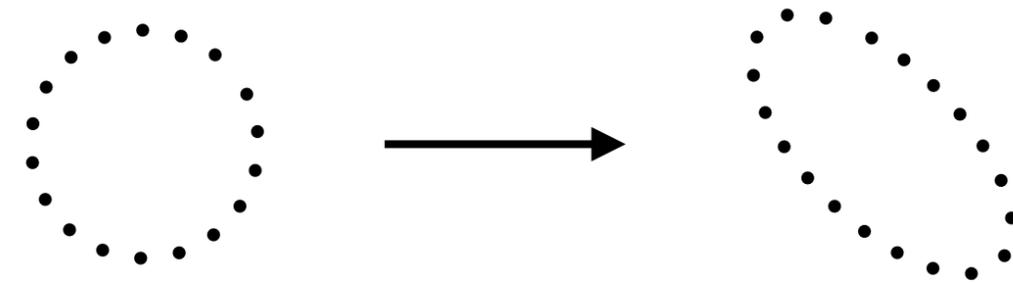
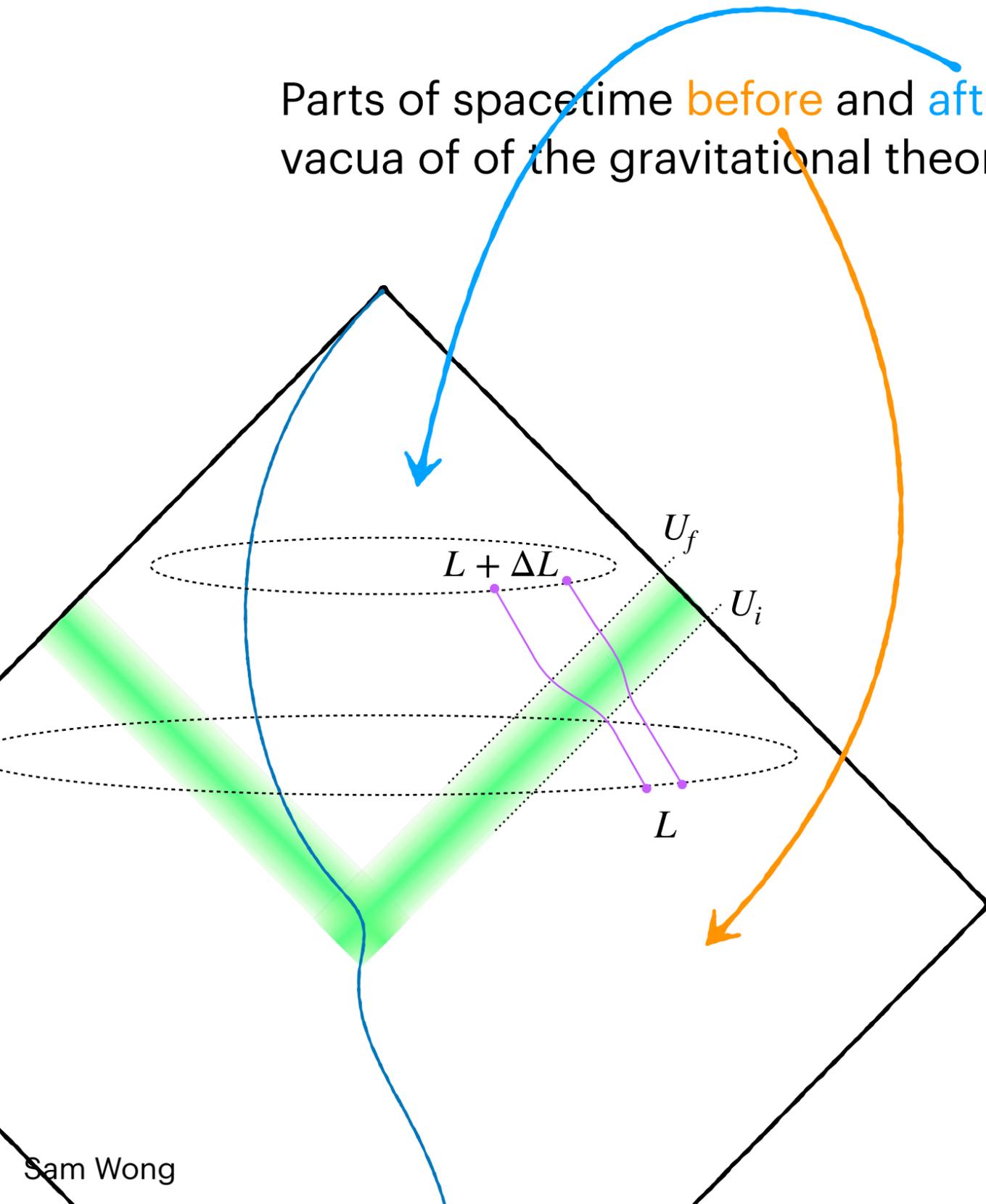
Soft modes and memory

Parts of spacetime **before** and **after** the passage of wave are described by different vacua of the gravitational theory.

These vacua are related to each other by diffs.

For instance: boosted black hole

Can be thought as “Rescaling by long wavelength GWs”



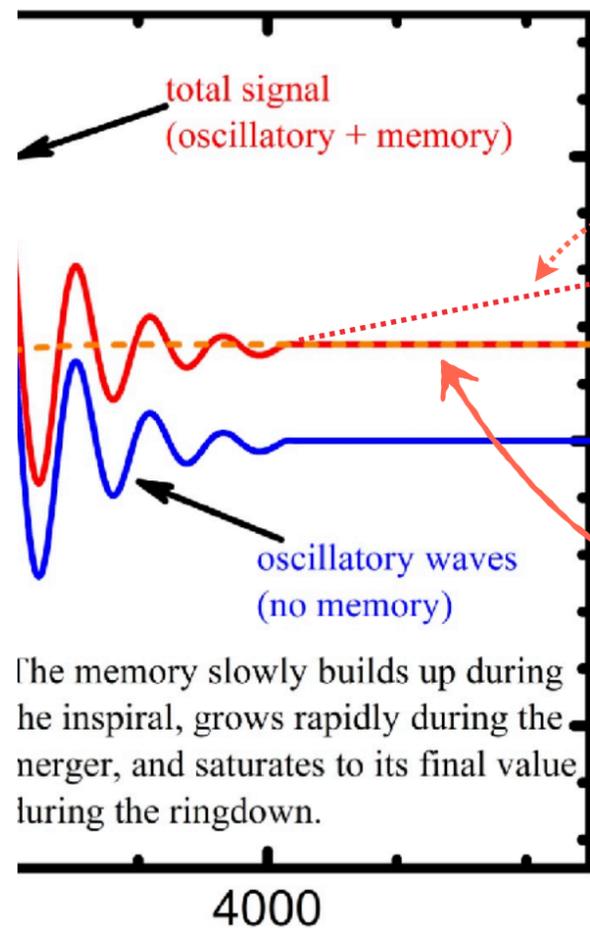
Memories can be characterized by symmetries!

(Caution: the part of the spacetime with the wave train cannot be described by diff)

Soft modes and memory

Local TT frame residual diffeomorphism $x^i \rightarrow x^i + \xi^i(x)$

The multipoles exactly describe the *long wavelength modes* of those memory terms:

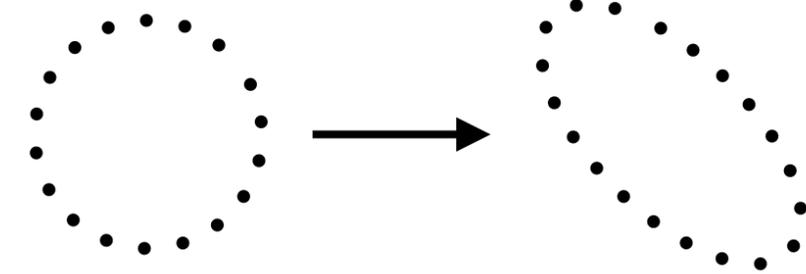
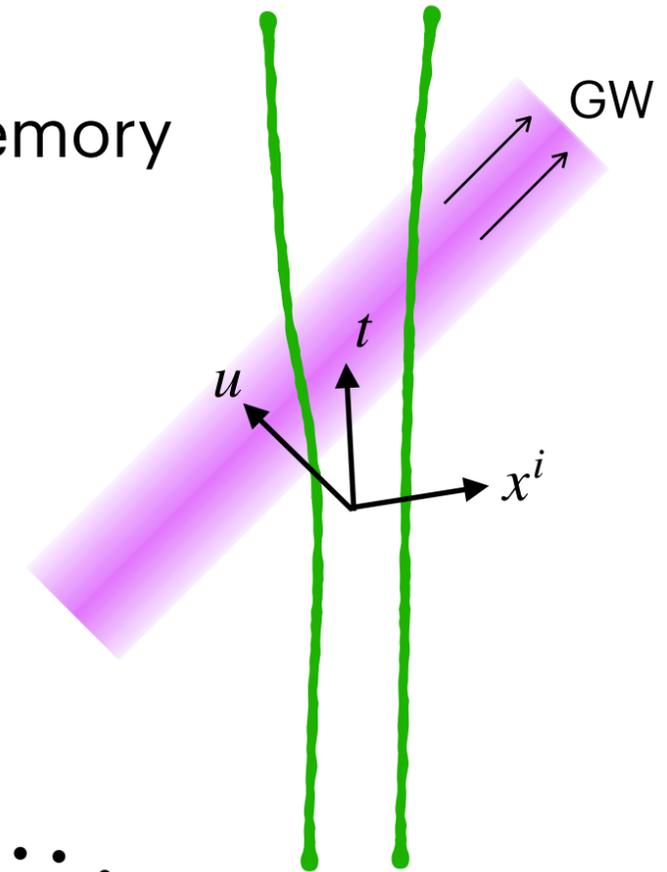


$$\xi_i = - \left(A_{ij} + B_{ij}u \right) x^j - \frac{1}{2} B_{jk} \bar{N}_i x^j x^k + \dots$$

Time dependent
(mirrors flying apart)

$$H_{ij}^{\text{TT}}(U, N) = A_{ij}(\bar{N}) + B_{ij}(\bar{N})u + \dots$$

time independent part
(constant displacement)



Anisotropic rescaling

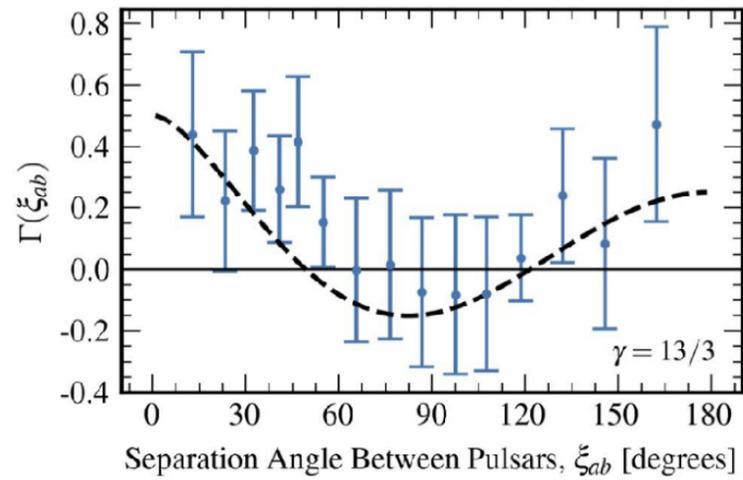
(The higher multipoles characterizes all possible memories)

[De Luca, Khoury & Wong, 2412.01910]

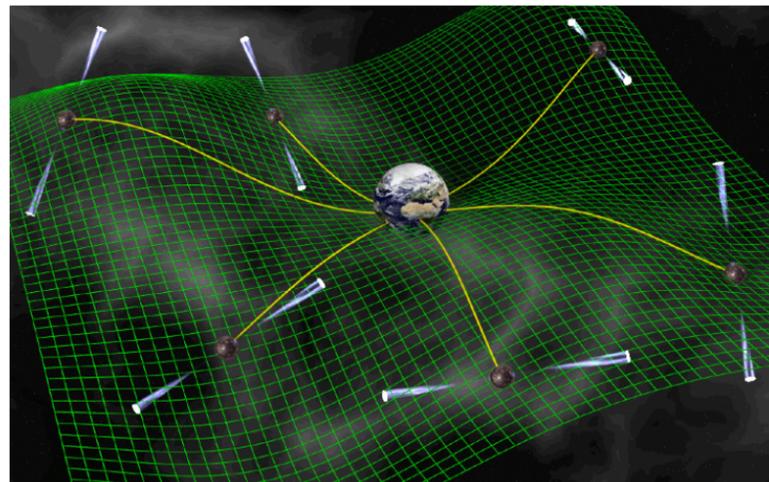
[De Luca, Khoury & Wong, 2412.12273]

Stochastic signals

SGWB (PTA & Astrometry)

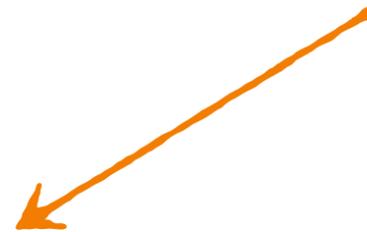


[NANOGrav 15-yr, 2306.16213]

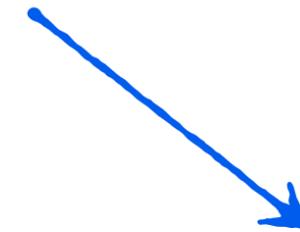


Photon geodesics:

$$x^\mu(\lambda) = \left(x^0(\lambda), x^i(\lambda) \right)$$



PTA

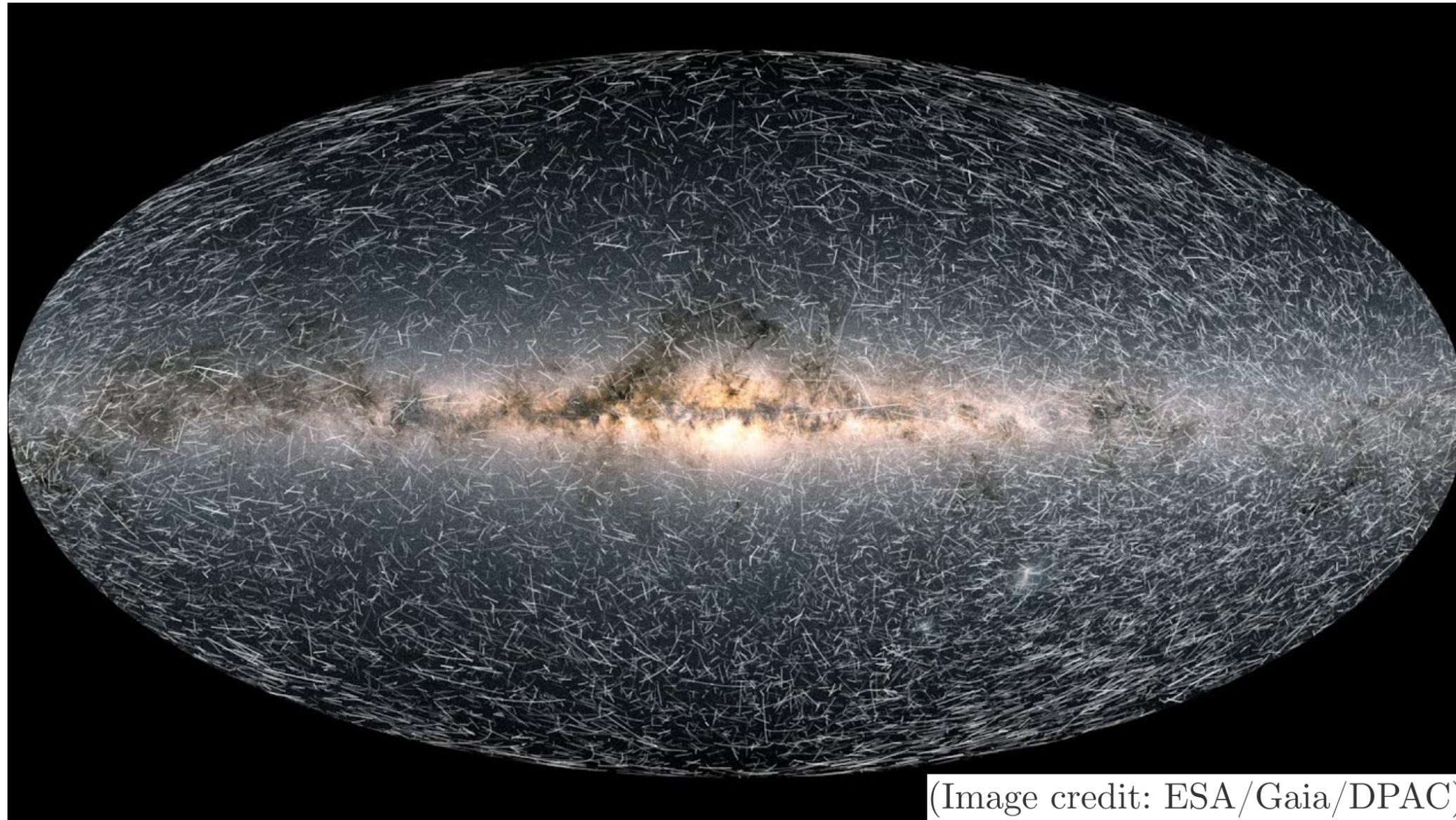


Astrometry



Deflection of star positions

GAIA: billions of stars, 10^6 quasars (for SGWB), $\sim 10\text{-}100\mu\text{as}$



(Image credit: ESA/Gaia/DPAC)

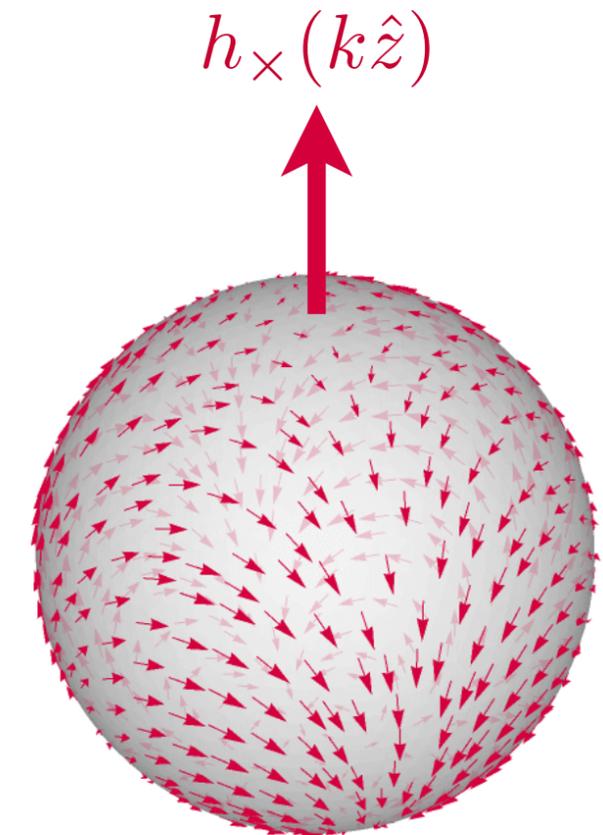
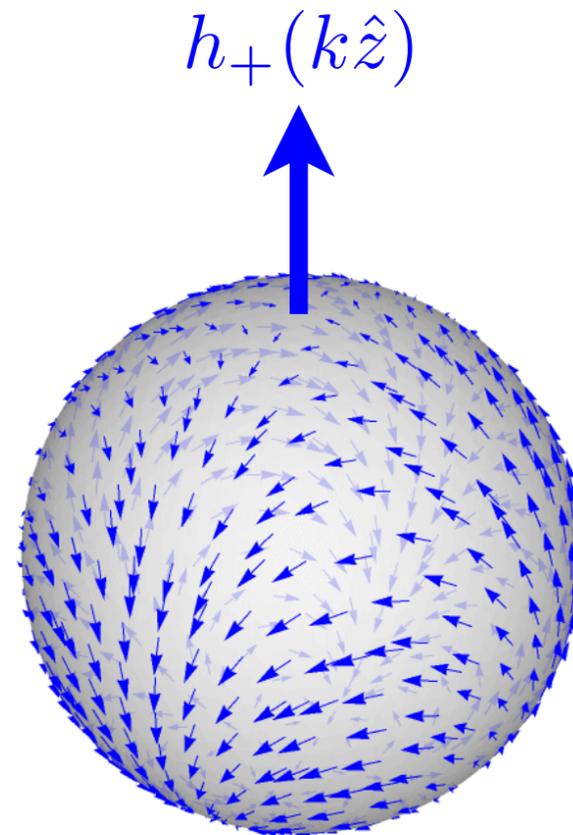
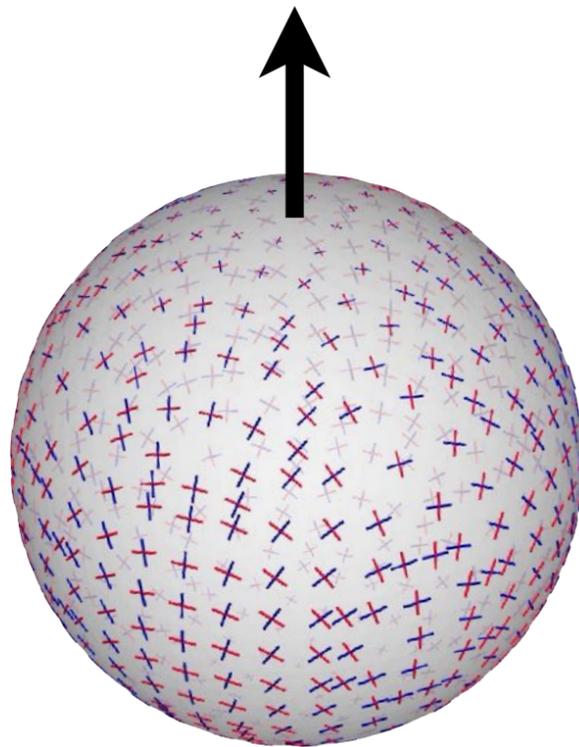
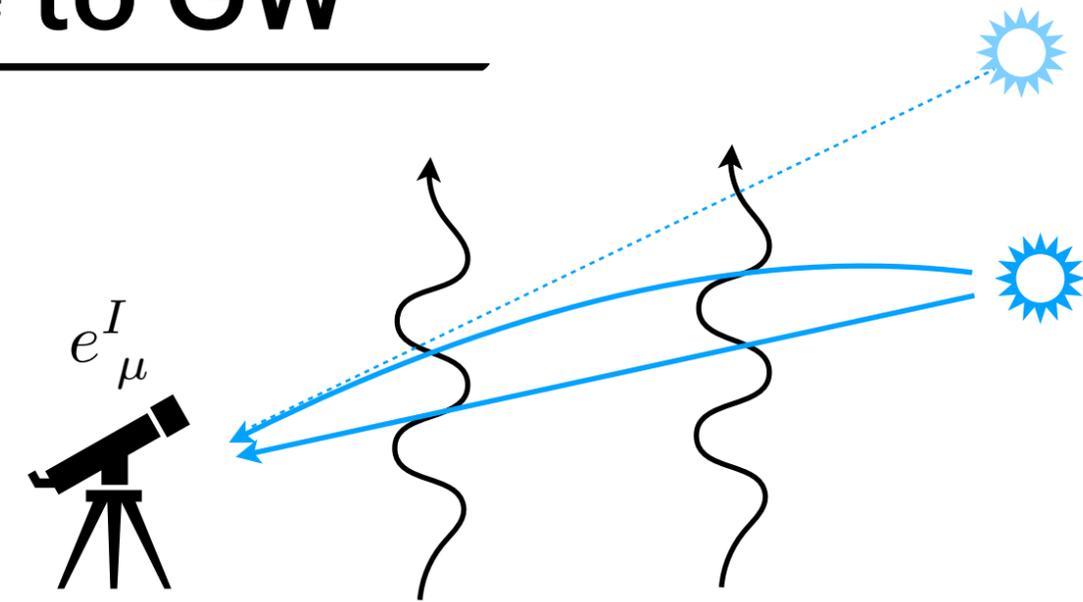
Deflection of star-lights due to GW

$$x^\mu(\lambda) = (x^0(\lambda), x^i(\lambda))$$

Defelction: $\delta n^I(t, \hat{n}) = e^I{}_\mu \delta n^\mu(t, \hat{n}) = \mathcal{R}^{IJK}(\hat{n}, \hat{k}) h_{JK}(\vec{k}) e^{-ikt}$

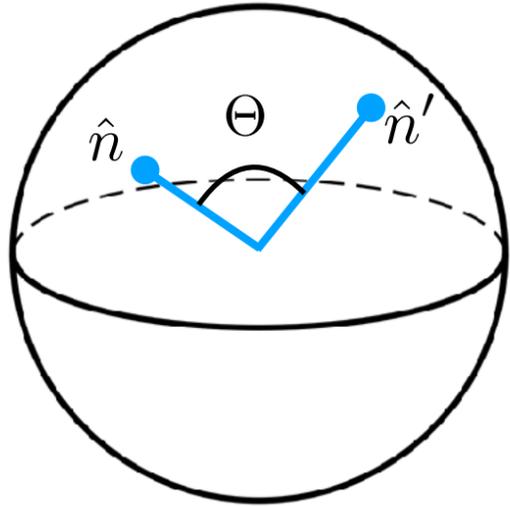
$$\mathcal{R}^{IJK}(\hat{n}, \hat{k}) = \frac{\hat{n}^I + \hat{k}^I}{2(1 + \hat{k} \cdot \hat{n})} \hat{n}^J \hat{n}^K - \frac{1}{2} \delta^{IJ} \hat{n}^K$$

[Book & Flanagan, 1009.4192]



Parity signal in astrometry

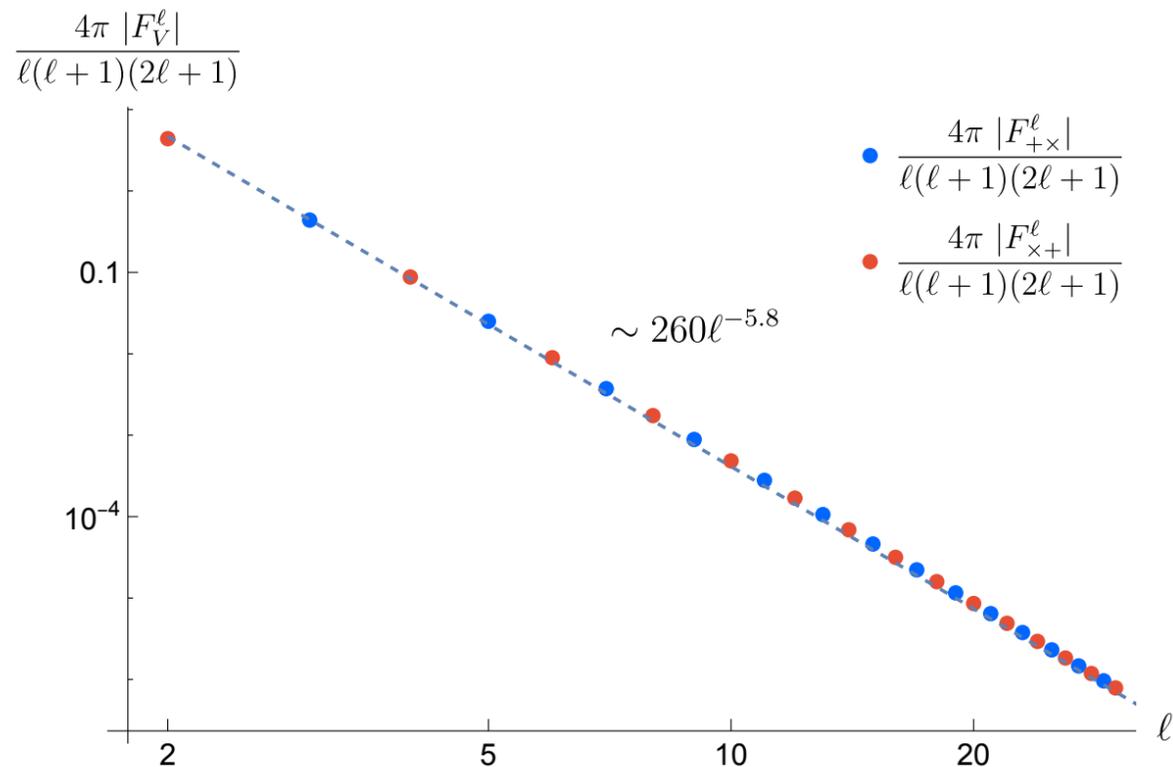
Vector spherical harmonics (E/B modes, cosmologist's favourite)



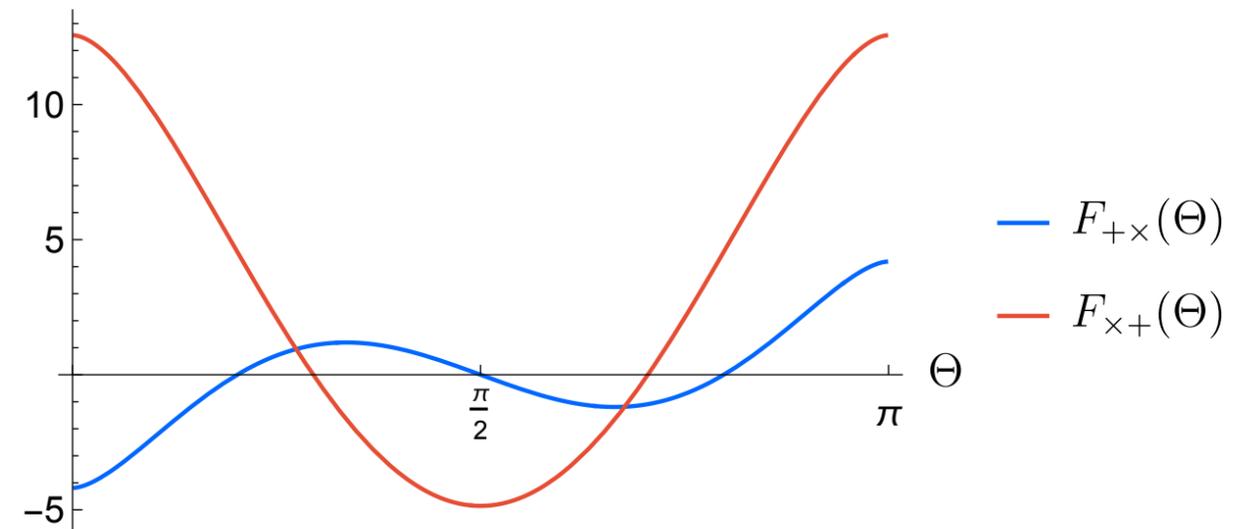
$$\delta n(t, \hat{n}) = \sum_{\ell m} \left[\delta n_{E\ell m}(t) \vec{Y}_{\ell m}^E(\hat{n}) + \delta n_{B\ell m}(t) \vec{Y}_{\ell m}^B(\hat{n}) \right]$$

EB correlations:

$$\begin{aligned} \langle \delta n_{E\ell m}(t) \delta n_{B\ell' m'}(t)^* \rangle &= \int d^2\Omega_{\hat{n}} d^2\Omega_{\hat{n}'} Y_{\ell m I}^{E*}(\hat{n}) Y_{\ell' m' J}^B(\hat{n}') \langle \delta n^I(t, \hat{n}) \delta n^J(t, \hat{n}') \rangle \\ &= \frac{1}{\ell(\ell+1)} \int d^2\Omega_{\hat{n}} d^2\Omega_{\hat{n}'} Y_{\ell m}^*(\hat{n}) Y_{\ell' m'}(\hat{n}') \epsilon_{JKL} \nabla_I \nabla'^L (\hat{n}'^K \langle \delta n^I(t, \hat{n}) \delta n^J(t, \hat{n}') \rangle) \\ &= \frac{\delta_{\ell\ell'} \delta_{mm'}}{\ell(\ell+1)} \frac{4\pi}{2\ell+1} (\mathcal{A}_{+\times} F_{+\times}^\ell + \mathcal{A}_{\times+} F_{\times+}^\ell) \end{aligned}$$



$$\epsilon_{JKL} \nabla_I \nabla'^L (\hat{n}'^K \langle \delta n^I(t, \hat{n}) \delta n^J(t, \hat{n}') \rangle) = \mathcal{A}_{+\times} F_{+\times}(\Theta) + \mathcal{A}_{\times+} F_{\times+}(\Theta)$$



[Liang, Lin, Trodden & Wong, 2309.16666]

Soft modes in stochastic signals

Remember, the presence of a long wavelength gravitational wave can be thought as a diff in the background geometry

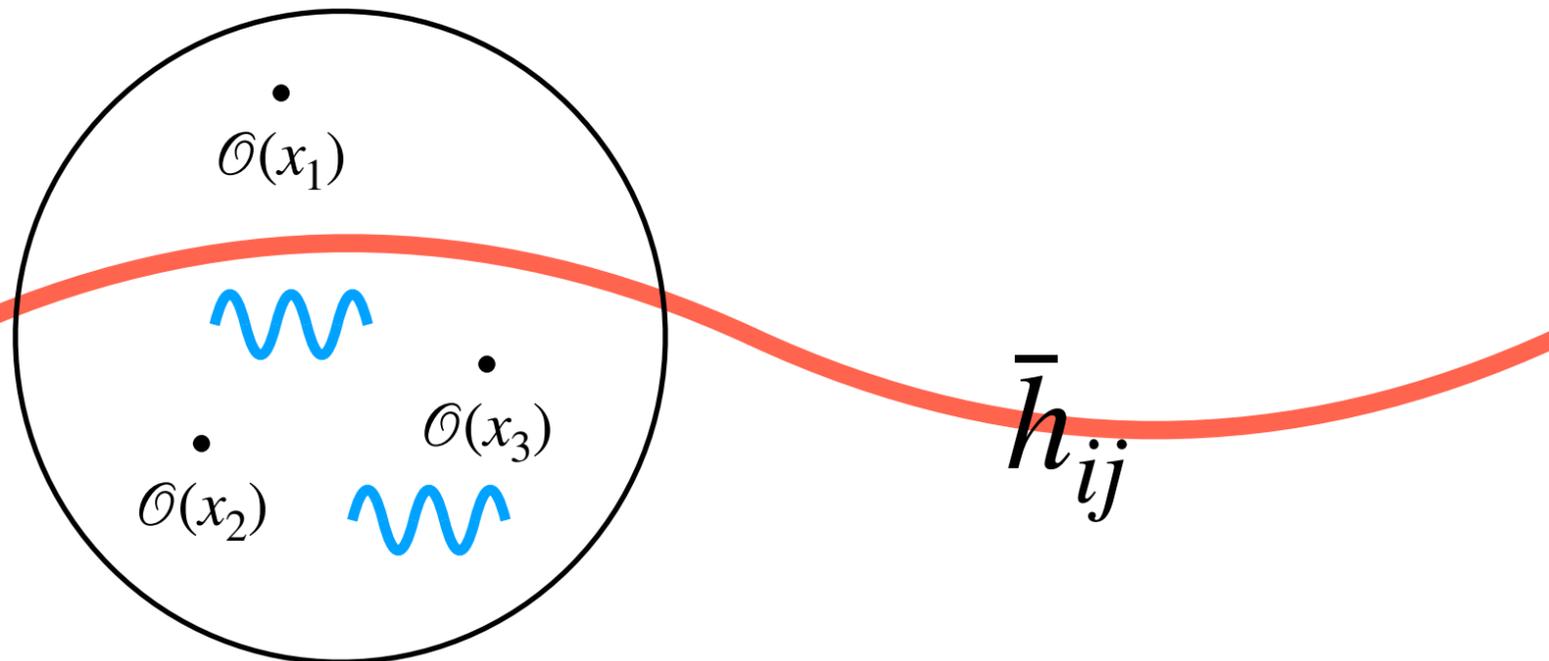
in the presence of a long mode

$$\langle \mathcal{O}(x_1, \dots, x_N) \rangle_{h \rightarrow \text{const.}} = \langle \mathcal{O}(\tilde{x}_1, \dots, \tilde{x}_N) \rangle$$

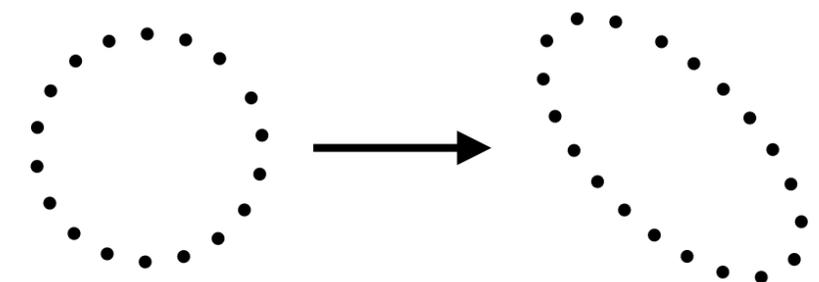
in a rescaled coordinate

expanding in the long mode

$$\Rightarrow \langle \mathcal{O}(\tilde{x}_1, \dots, \tilde{x}_N) \rangle = \left(1 + \frac{1}{2} \bar{h}^k{}_\ell \sum_{m=1}^N x_m^\ell \frac{\partial}{\partial x_m^k} + \mathcal{O}(\bar{h}^2) \right) \langle \mathcal{O}(x_1, \dots, x_N) \rangle$$



$$\tilde{x}^i = x^i + \frac{1}{2} \bar{h}^i{}_j x^j$$



Anisotropic rescaling

Soft modes in stochastic signals

in the presence of a long mode

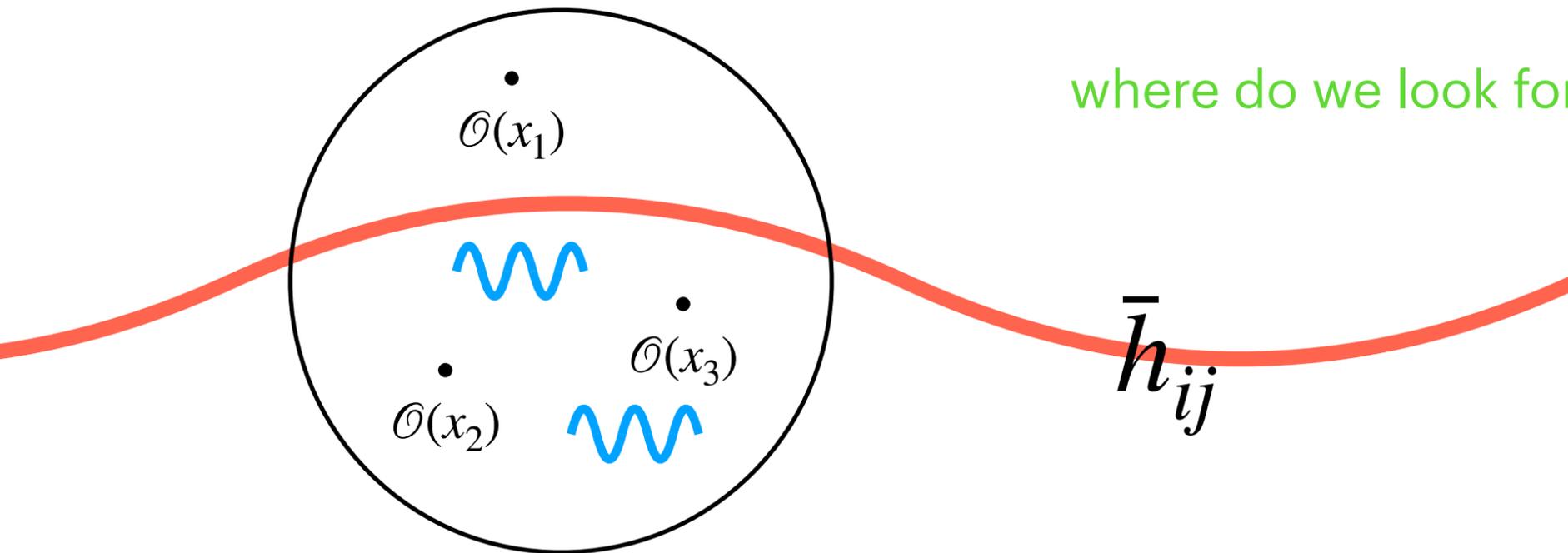
$$\langle \mathcal{O}(x_1, \dots, x_N) \rangle_{h \rightarrow \text{const.}} = \langle \mathcal{O}(\tilde{x}_1, \dots, \tilde{x}_N) \rangle$$

in a rescaled coordinate

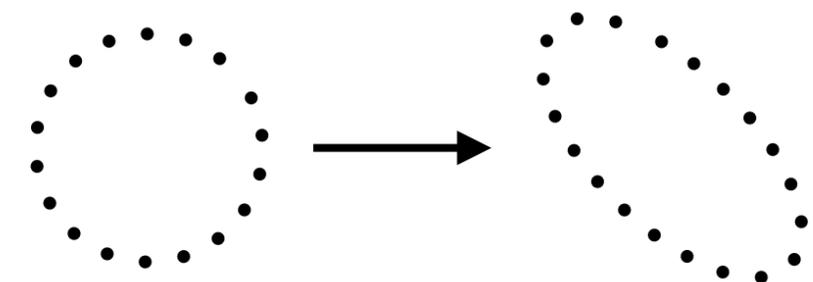
expanding in the long mode

$$\Rightarrow \langle \mathcal{O}(\tilde{x}_1, \dots, \tilde{x}_N) \rangle = \left(1 + \frac{1}{2} \bar{h}^k{}_\ell \sum_{m=1}^N x_m^\ell \frac{\partial}{\partial x_m^k} + \mathcal{O}(\bar{h}^2) \right) \langle \mathcal{O}(x_1, \dots, x_N) \rangle$$

$$\lim_{h \rightarrow \text{const.}} \langle h_{ij} \mathcal{O}(x_1, \dots, x_N) \rangle = \lim_{h \rightarrow \text{const.}} \frac{1}{2} \langle h_{ij} \bar{h}^k{}_\ell \rangle \sum_{m=1}^N x_m^\ell \frac{\partial}{\partial x_m^k} \langle \mathcal{O}(x_1, \dots, x_N) \rangle$$

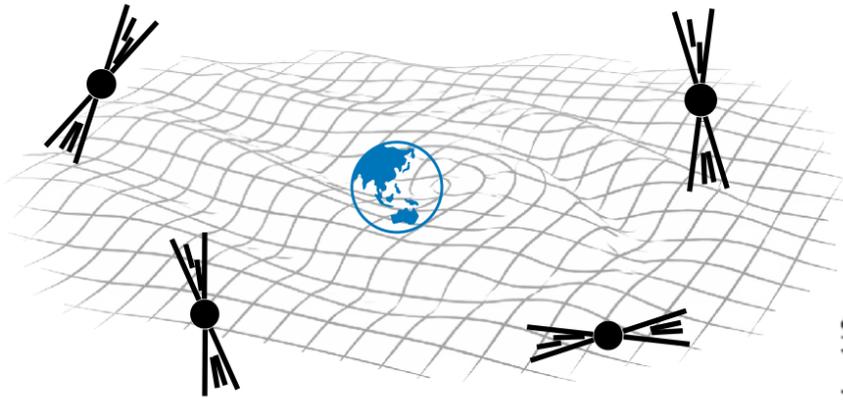


$$\tilde{x}^i = x^i + \frac{1}{2} \bar{h}^i{}_j x^j$$



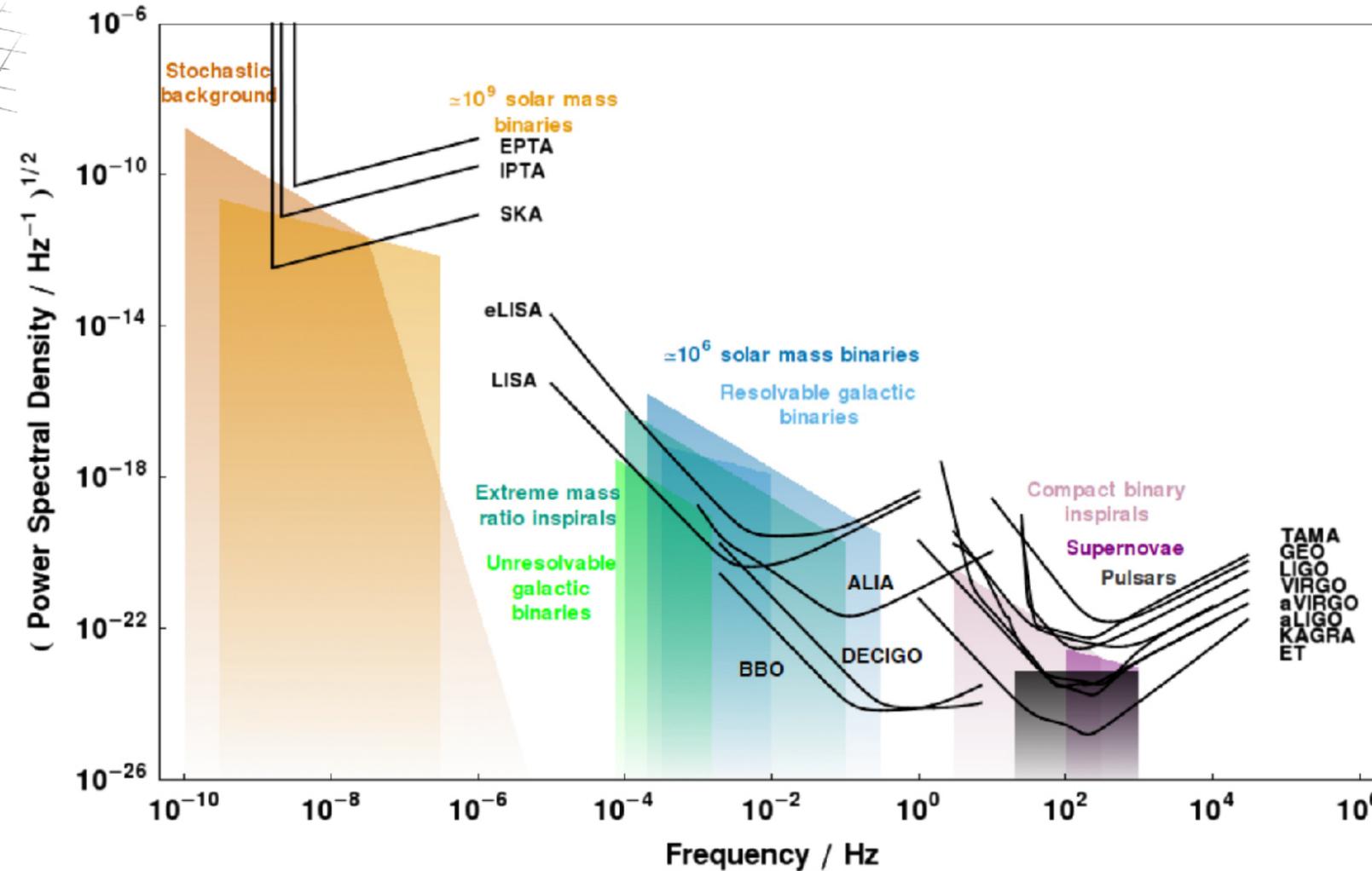
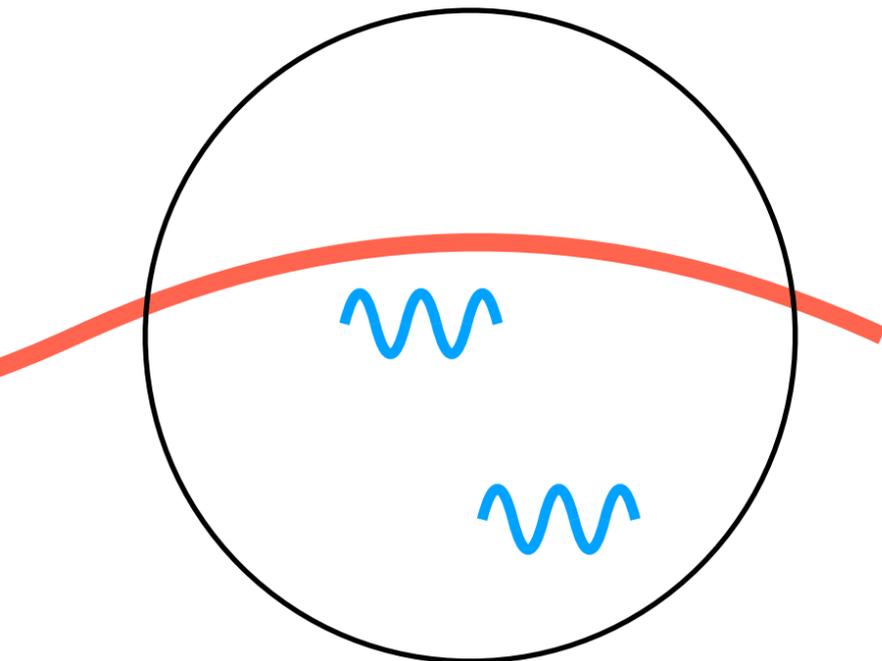
Anisotropic rescaling

PTA - Interferometer correlations



(PTA/astrometry)

Low frequency GWs,
 ~ 10 nHz, signal: $z_a(t)$

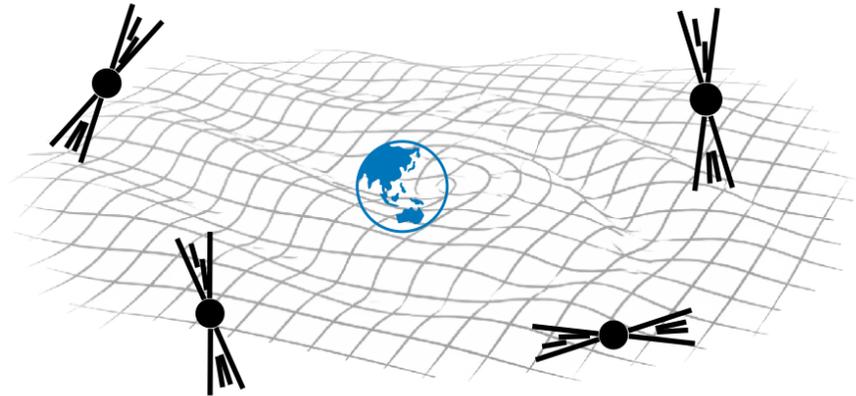


Higher frequency GWs,
 $10^{-4} \sim 1$ Hz, signal: $D(t)$

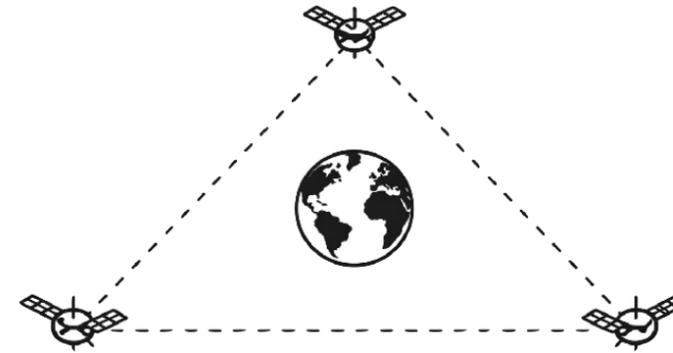
$$\text{Schematically: } \langle z_a(t) D_A(t) D_B(t) \rangle \sim \langle h_{f_1} h_{f_2} h_{f_3} \rangle_{f_1 \ll f_2, f_3}$$

PTA - Interferometer correlations

Relating three point functions to power spectrums at different frequency ranges

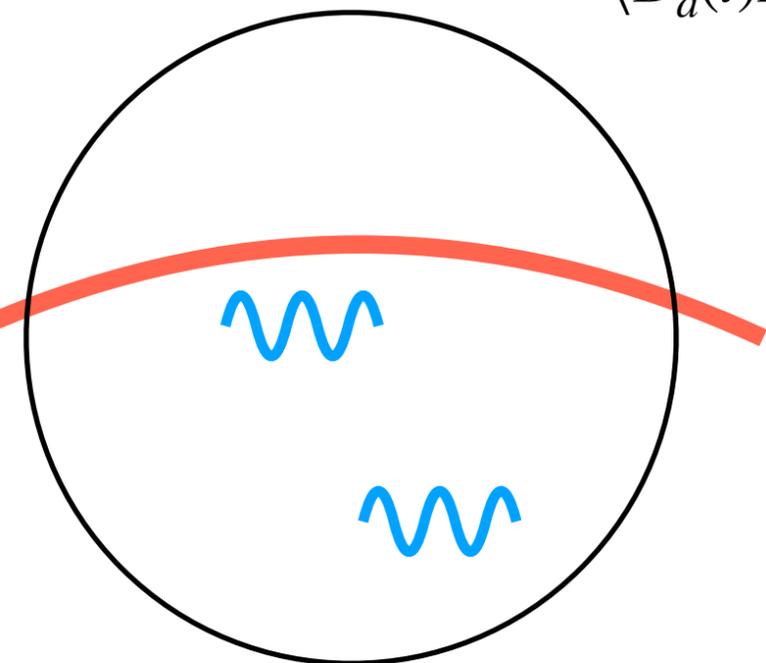


signal: $z_c(t)$, power spectrum at f_0



signal: $D_a(t)$

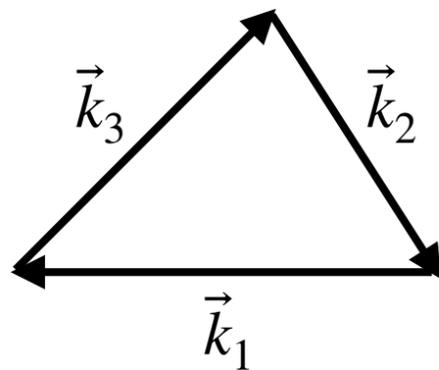
$$\begin{aligned}
 \langle D_a(t)D_b(t)z_c(t) \rangle = & \frac{1}{8} \int df_0 S_h(f_0) \int df S_h(f) \left\{ \int \frac{d^2\Omega}{4\pi} [1 - 3(n_c \cdot \Omega)^2] \sum_A F_a^A(\Omega) F_b^A(\Omega) \right. \\
 & + \int \frac{d^2\Omega}{4\pi} \left[\frac{1}{3} \Omega^j l_b^i (l_b \cdot \Omega) - \frac{1}{3} \Omega^j m_b^i (m_b \cdot \Omega) - l_b^i l_b^j (n_c \cdot l_b)^2 - m_b^i m_b^j (n_c \cdot m_b)^2 \right. \\
 & \left. \left. + n_c^i l_b^j ((n_c \cdot l_b) - (n_c \cdot \Omega)(l_b \cdot \Omega)) + n_c^i m_b^j ((n_c \cdot m_b) - (n_c \cdot \Omega)(m_b \cdot \Omega)) \right] \sum_A F_a^A(\Omega) e_{ij}^A(\Omega) \right. \\
 & \left. + (a \leftrightarrow b) \right\}
 \end{aligned}$$



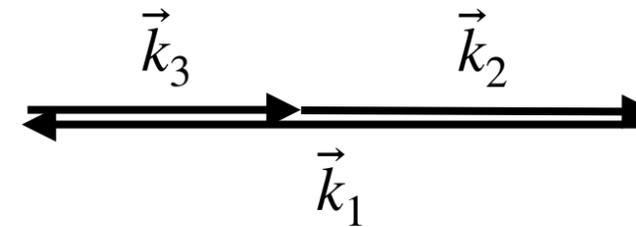
Non-gaussianities in the bispectrum

- The squeezed limit obtained from $\langle z_a(t)D_b(t)D_c(t) \rangle$ is a test of gravity (diff. inv.) and a signal of local non-gaussianity
- Correlations of different detectors, $\langle D_a(t)D_b(t)D_c(t) \rangle$, contain interesting information such as “time translational invariance” of the signal

wave vectors of the modes:



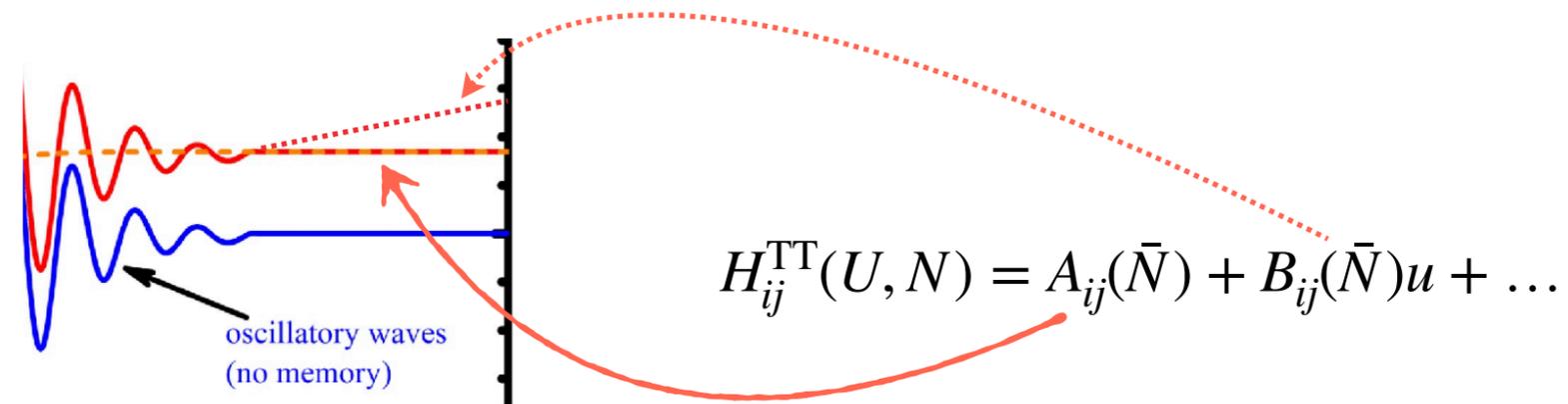
early universe signals



time translational invariance enforces $\delta(f_1 + f_2 + f_3)$

Brief summary

Single event signal: memories characterized by symmetry



Stochastic signal: symmetry dictates the squeezed limit, fixing local non-gaussianity

$$\langle z_a(t) D_b(t) D_c(t) \rangle \sim S(f_{\text{low}}) S(f_{\text{high}})$$

