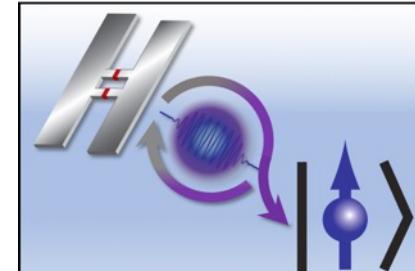


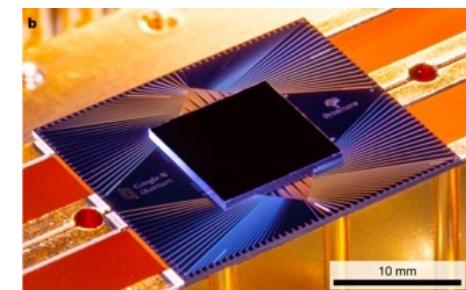
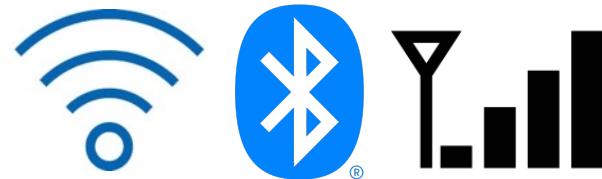
Quantum Information Technologies with Hybrid Systems

Yuimaru Kubo



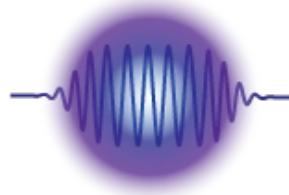
Advantages of microwave

- Well established (4G, 5G) -> operable
 - easy to stabilize phase
 - economically reasonable below ~ 10 GHz
- Easy mode matching (impedance matching) -> handy!
- Solid devices -> integrable



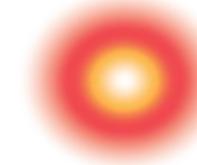
Quantum microwave at millikelvin

$5 \text{ GHz} \times h (= 250 \text{ mK} \times k_B) \gg 10 \text{ mK} \times k_B$



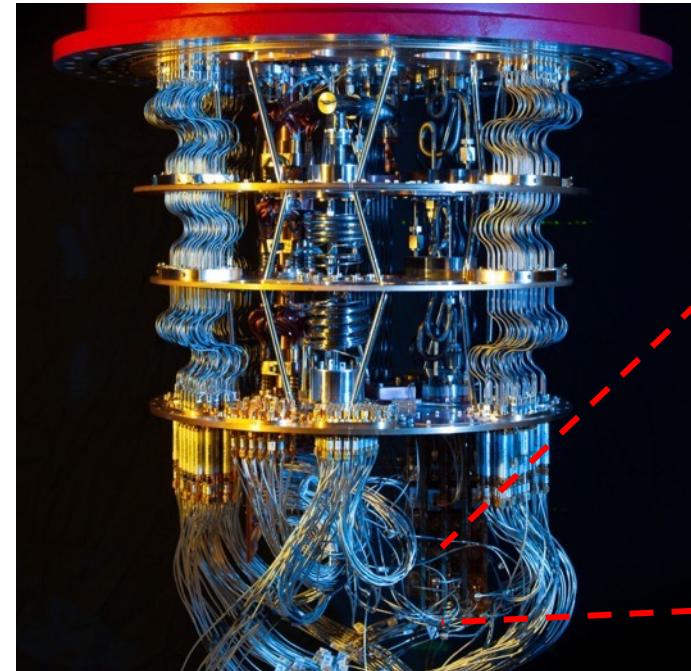
Quantum microwave
(microwave photon)

$400 \text{ THz} (= 10^4 \text{ K}) \gg 300 \text{ K}$

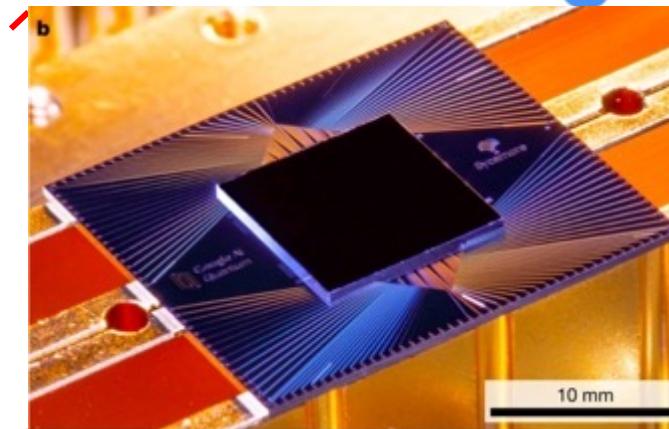


Optical photon

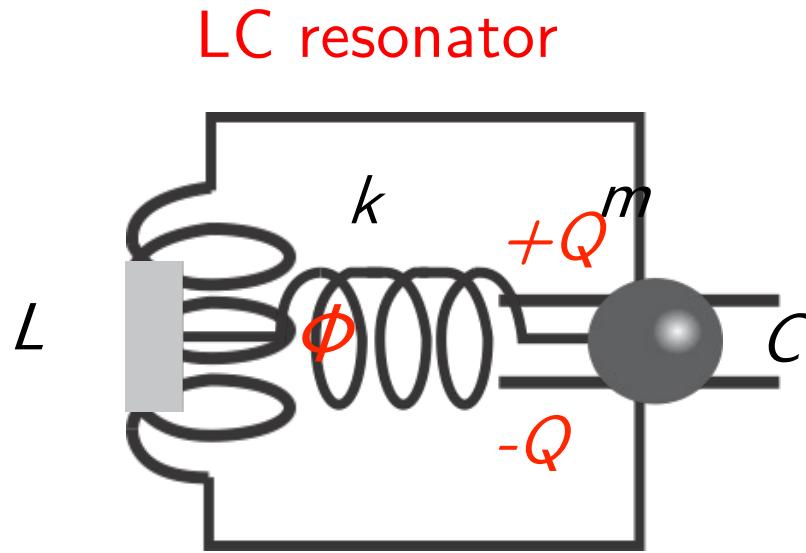
Dilution refrigerator



Google

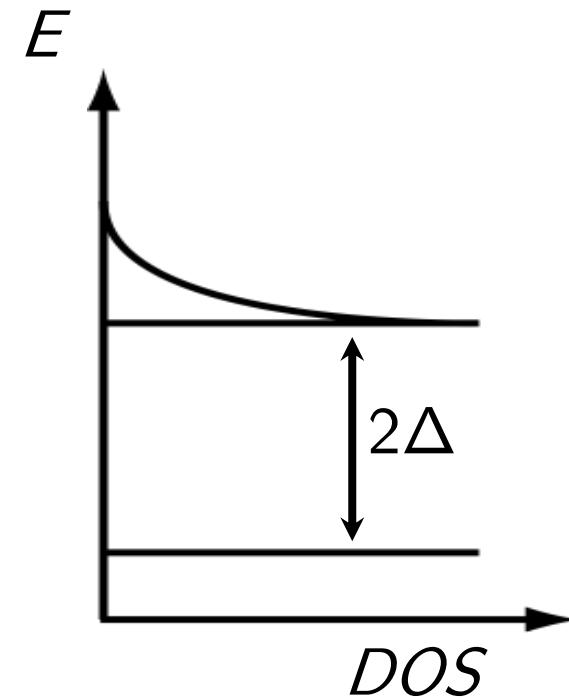


Quantum integrated circuit



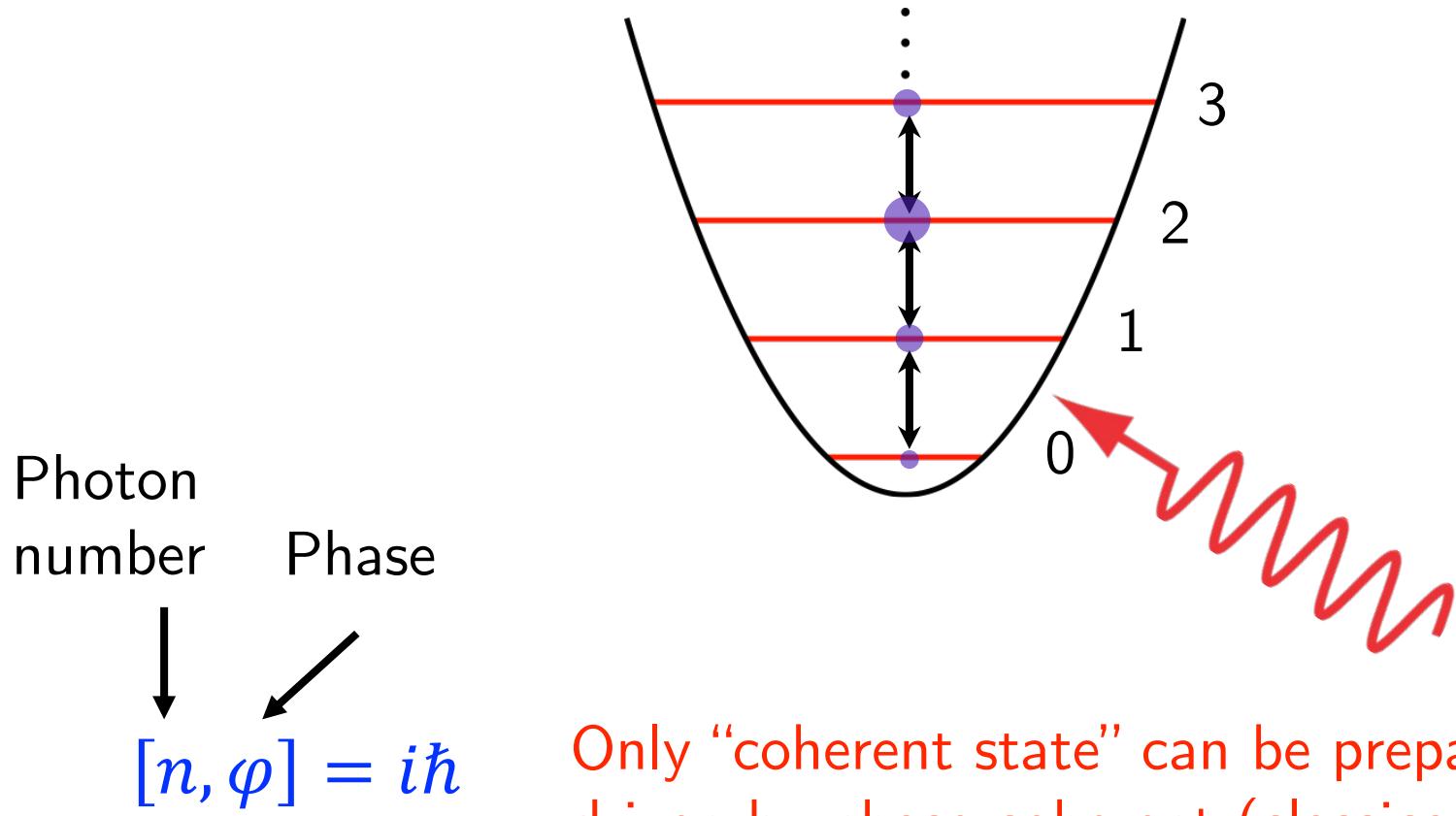
$$\omega = \frac{1}{\sqrt{LC}} \quad [\phi, Q] = i\hbar$$

For details, see e.g., Devoret and Martinis,
Quantum Inform. Processing 3, 163 (2004)



- Low loss -> high Q
- Macroscopic phase coherence
- Superconducting gap ($> \sim 100$ GHz)

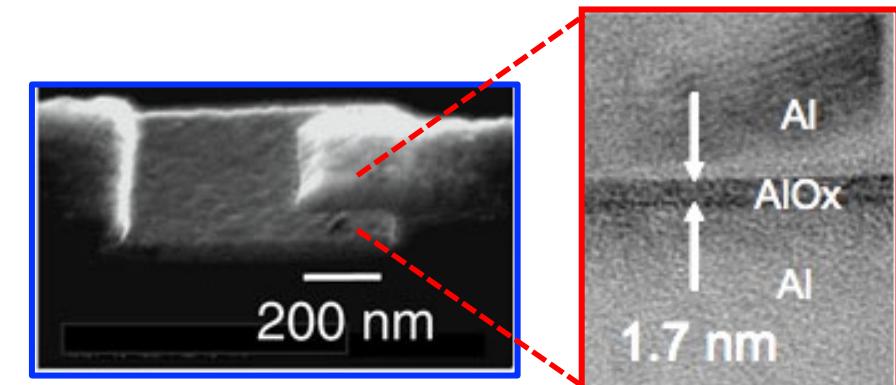
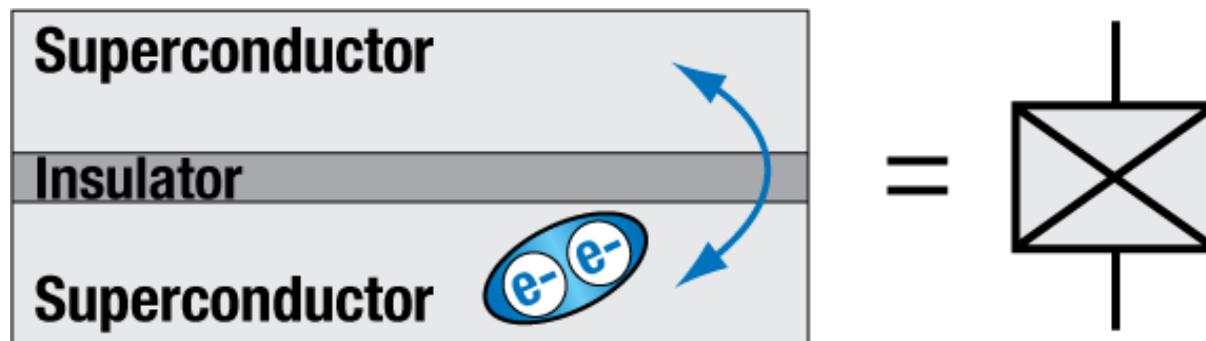
Harmonic oscillator



Only “coherent state” can be prepared when driven by phase-coherent (classical) light

No quantum “bit” possible

Superconducting qubit: Josephson junction



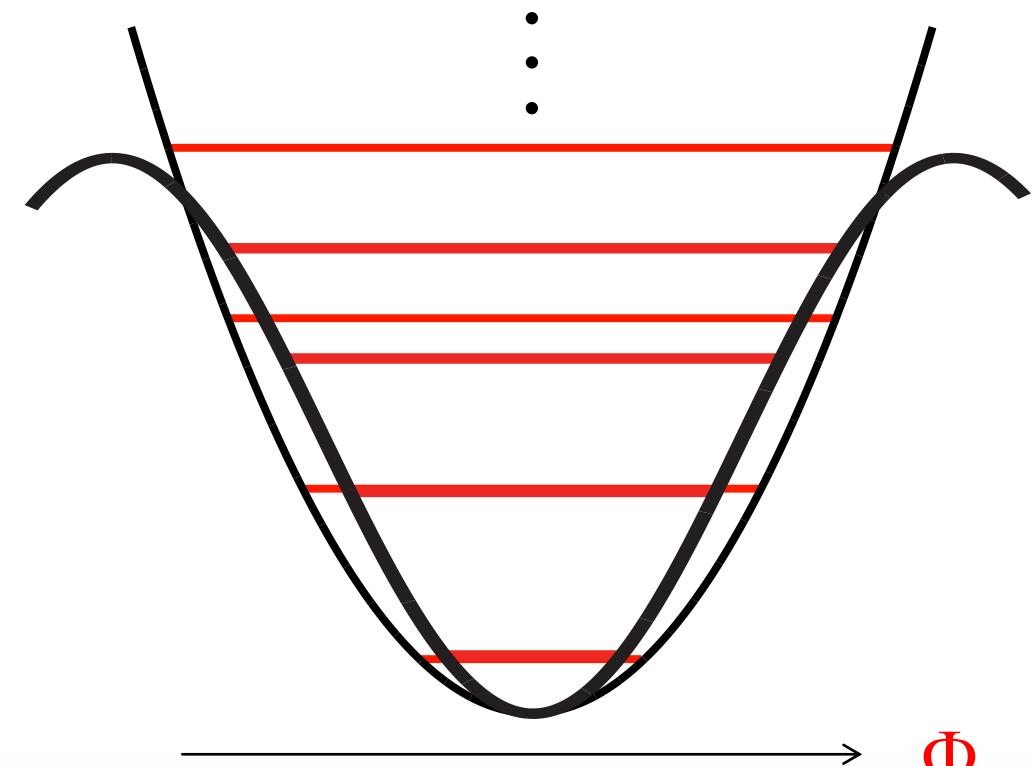
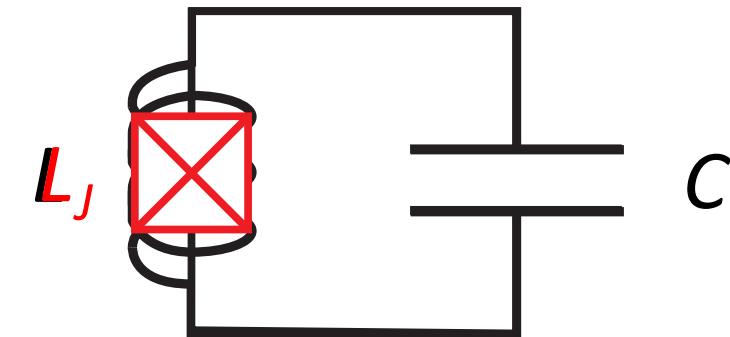
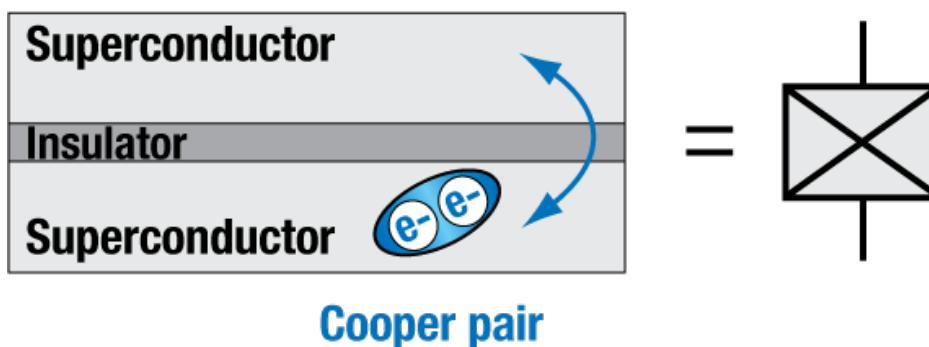
Cooper pair

$$H_q = \frac{Q^2}{2C} - \frac{\hbar I_0}{2e} \cos\left(2\pi \frac{\Phi}{\Phi_0}\right)$$

I_0 : critical current of Junction
 $\Phi_0 = h/2e$: Flux quantum

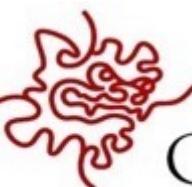
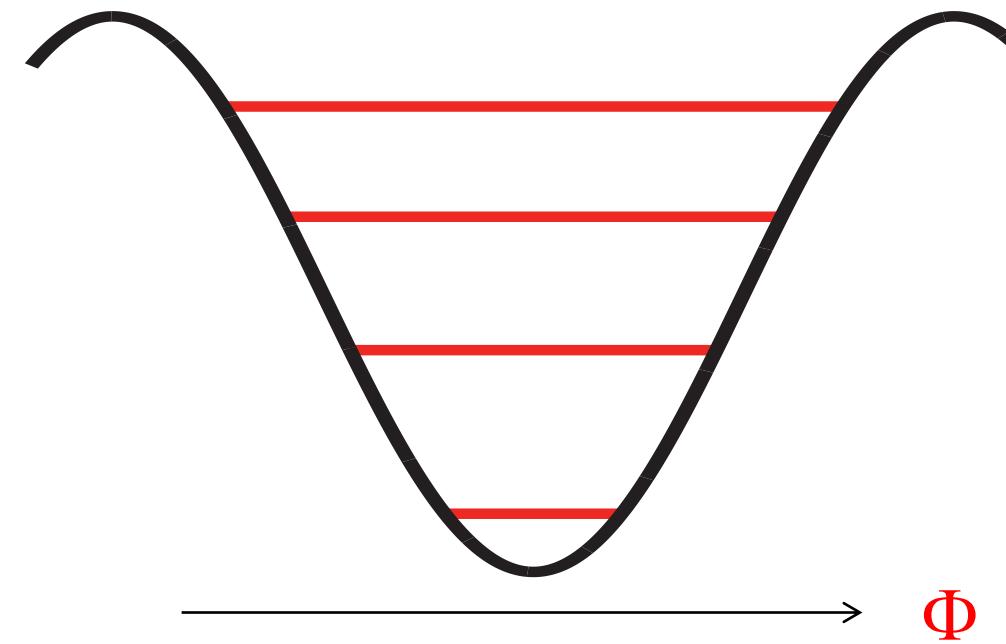
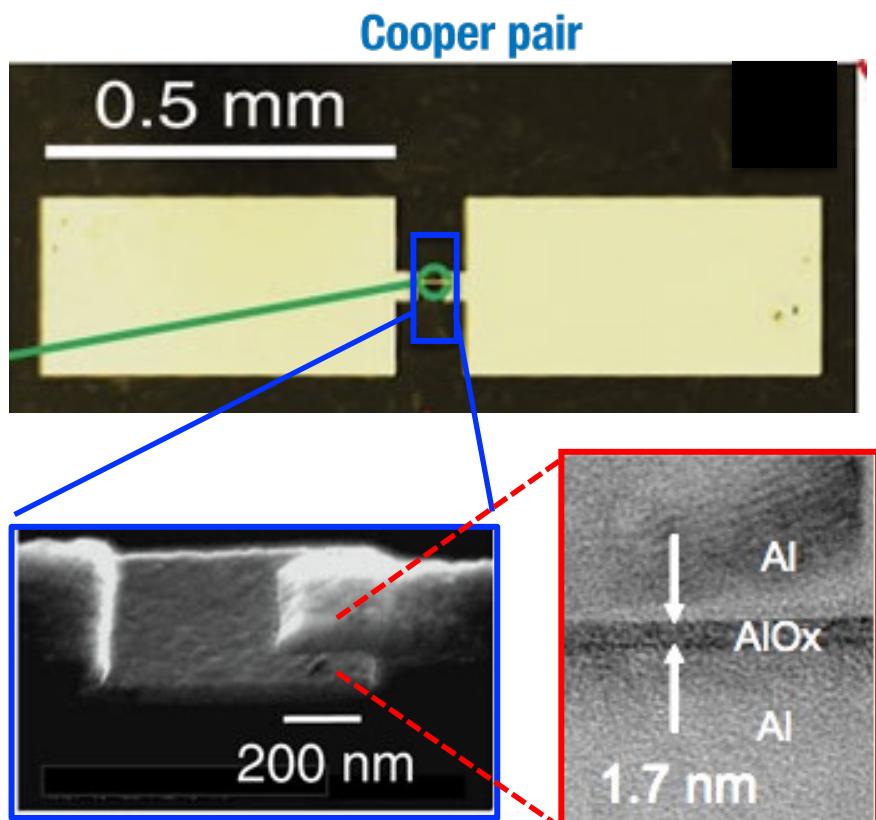
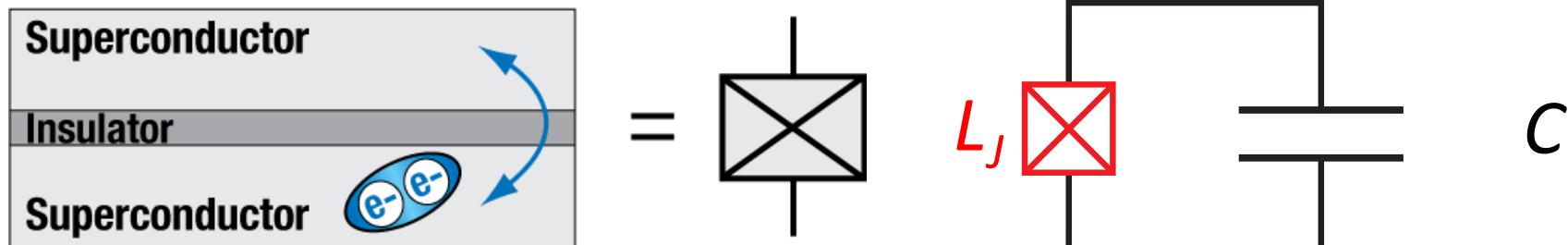
- Point-like lumped element inductance
- Non-linear & Non dissipative inductance at a single quantum of energy (single “microwave photon”)

Superconducting qubit: Josephson junction

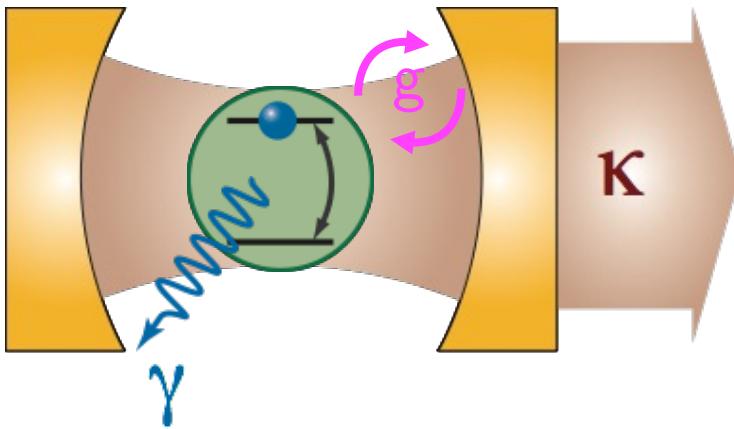


- Point-like lumped element inductance
- Non-linear & **Non dissipative** at **single quantum of energy (single microwave photon)**

Superconducting qubit: Josephson junction



Cavity-QED

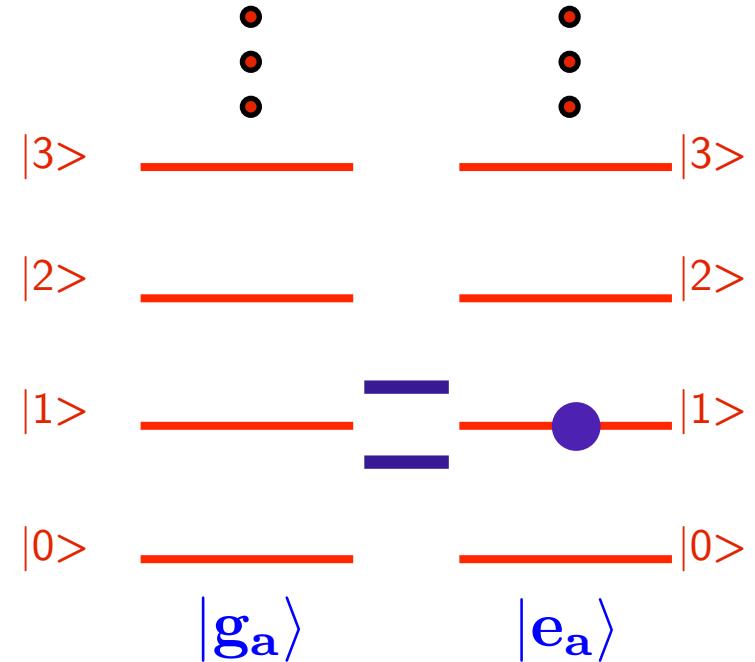
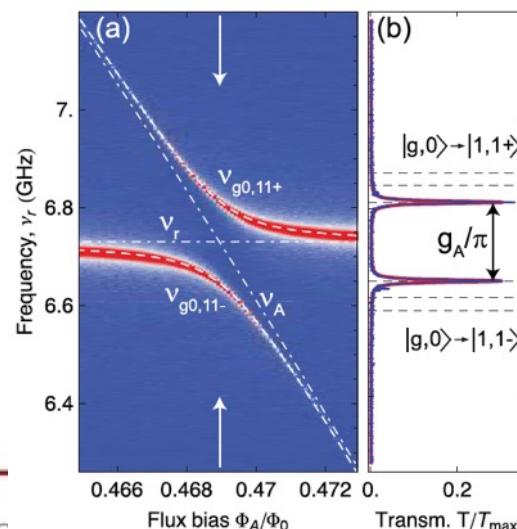


Vacuum Rabi Oscillation

$$H_{JC}/\hbar = \omega_r a^\dagger a + \frac{1}{2} \omega_a \sigma_z + g(a^\dagger \sigma^- + a \sigma^+)$$

$$g \gg \kappa, \gamma$$

Strong coupling

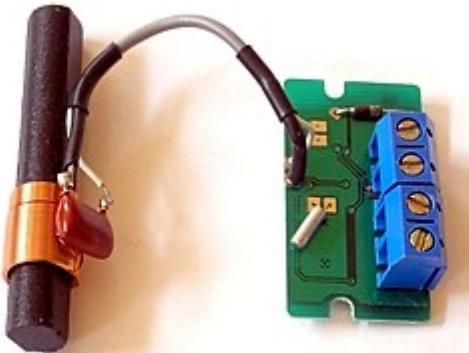


OIST

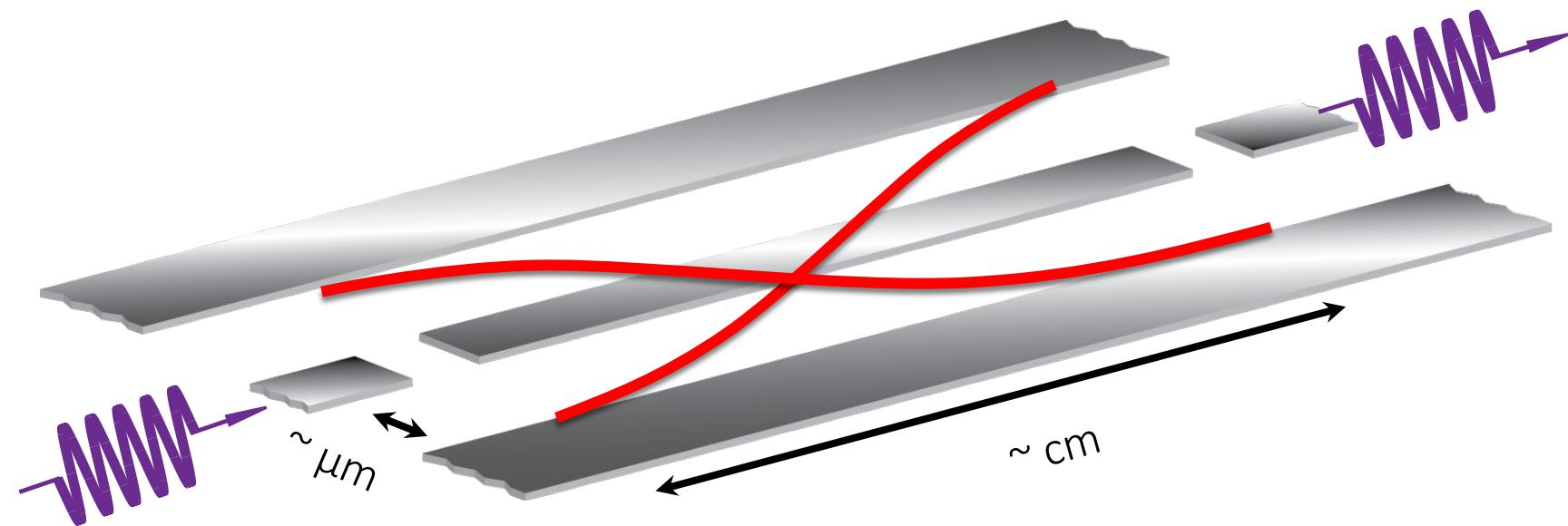
OKINAWA INSTITUTE OF SCIENCE AND TECHNOLOGY

Yuimaru Kubo |

Microwave LC resonator

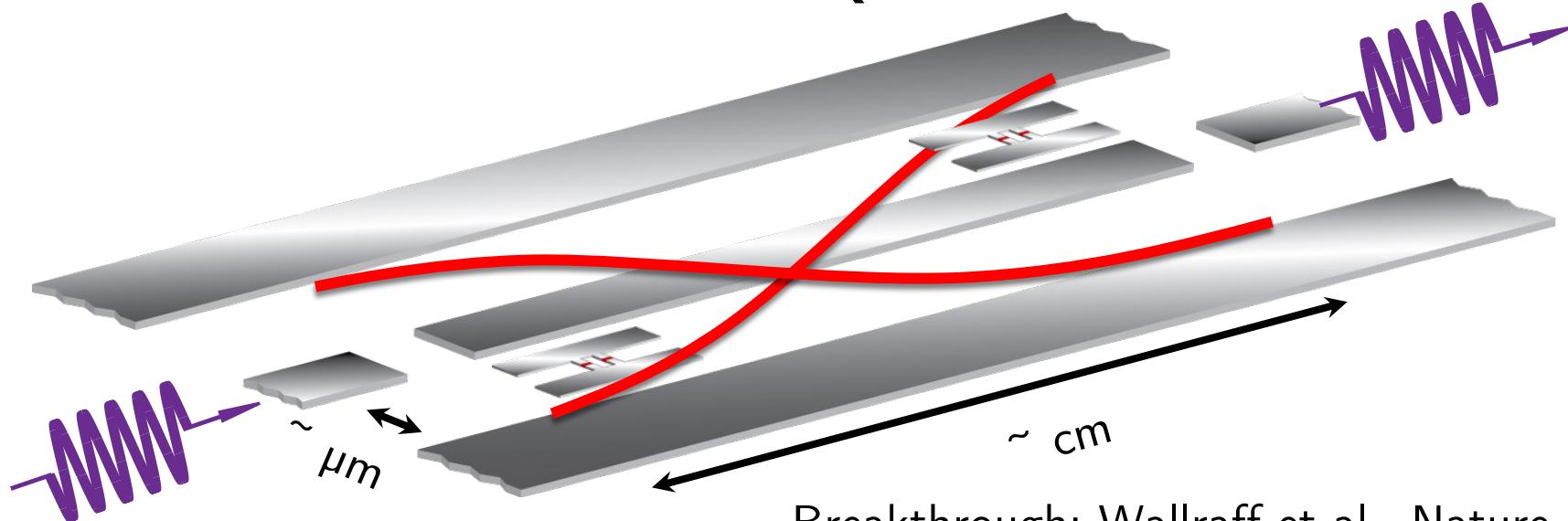


(from wikipedia)



- Losses and parasitic effect in “bulky lumped element”
- Low-loss and reliable designability using planar superconductor
- 2D version of “coaxial line” (e.g. cable from wall to TV)

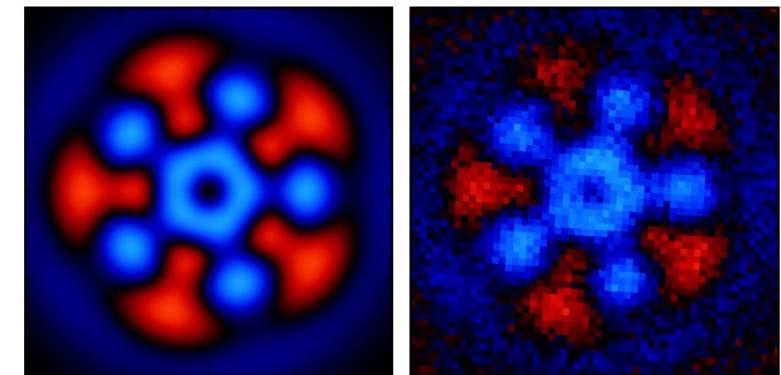
Circuit-QED



Breakthrough: Wallraff et al., Nature 431 162 (2004)

- Strong coupling of “macroscopic” qubit to microwave photons
- Gigantic dipole -> unprecedented parameter regime
- Manipulating and detecting **single artificial atoms** and **single microwave photons**

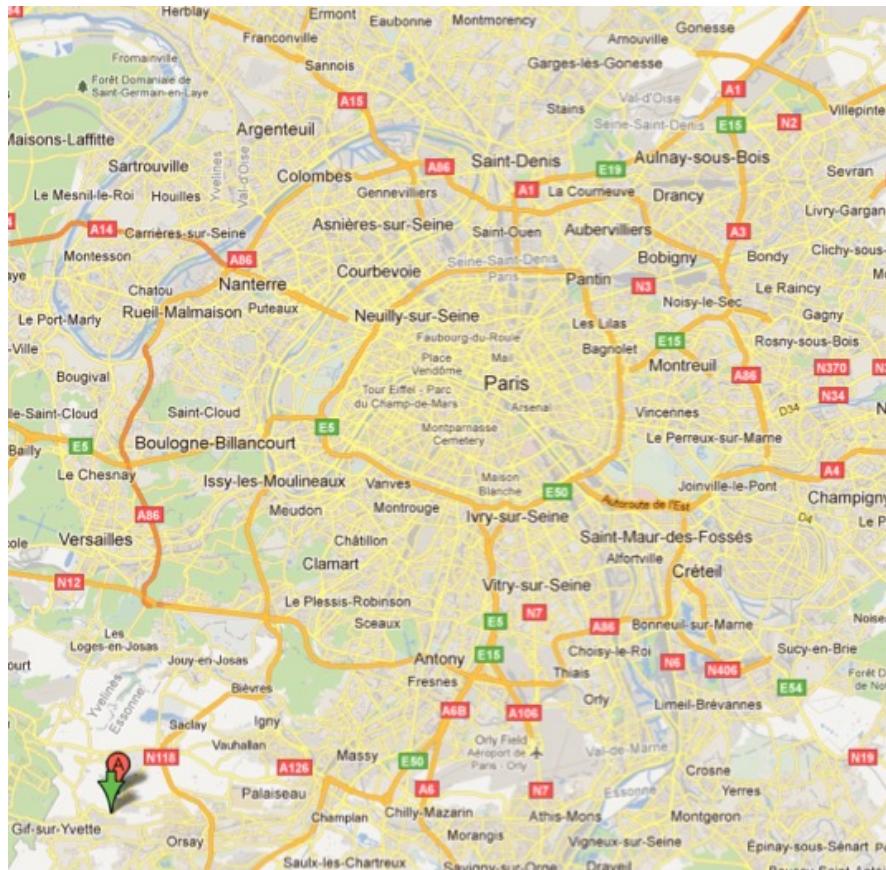
Wigner function of $|0\rangle + |5\rangle$



M. Hofheinz et al. Nature 459, 546 (2009)

CEA-Saclay

CEA = Atomic Energy Comission (and Alternative Energy)

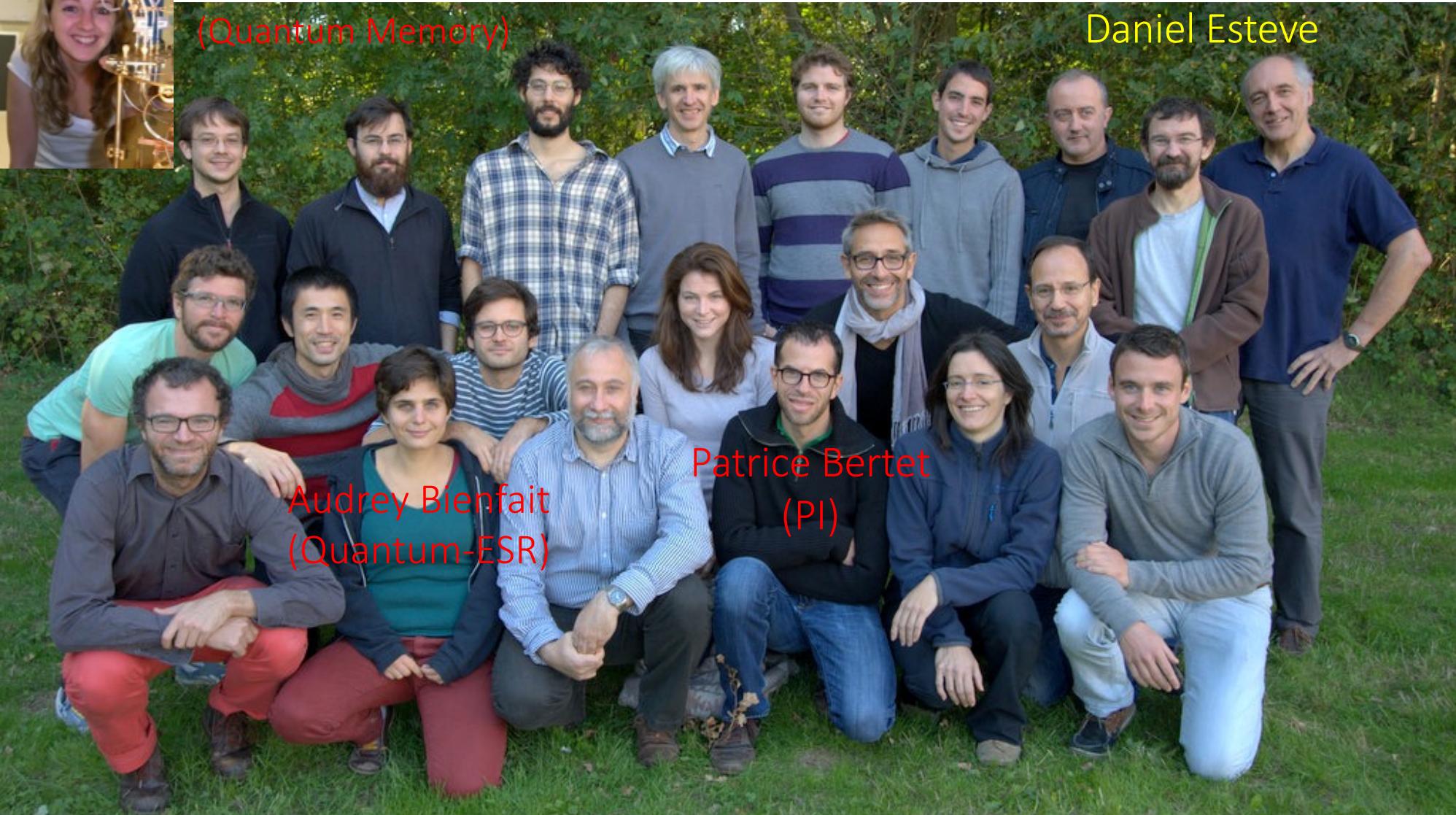


20 km from Paris

QUANTum electRONICS Group at CEA-Saclay



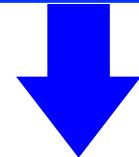
Cecile Grezes
(Quantum Memory)



Spins in solids

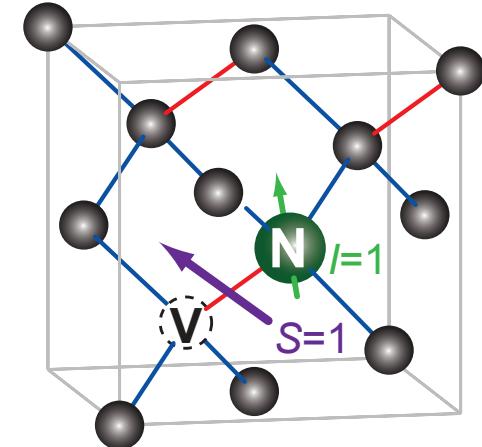
- Very good coherence in “silent” environment (e.g., ^{12}C , ^{28}Si)
- Both microwave (spin) and optical (orbital) transitions
- Scalability, designability ?
- Control on demand ?

Complementary for superconducting circuits

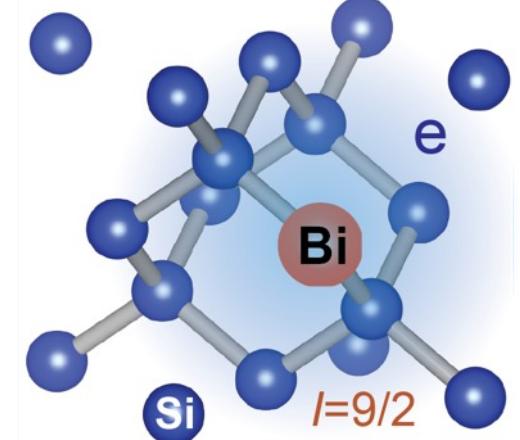


Hybrid quantum systems

NV centers in diamond

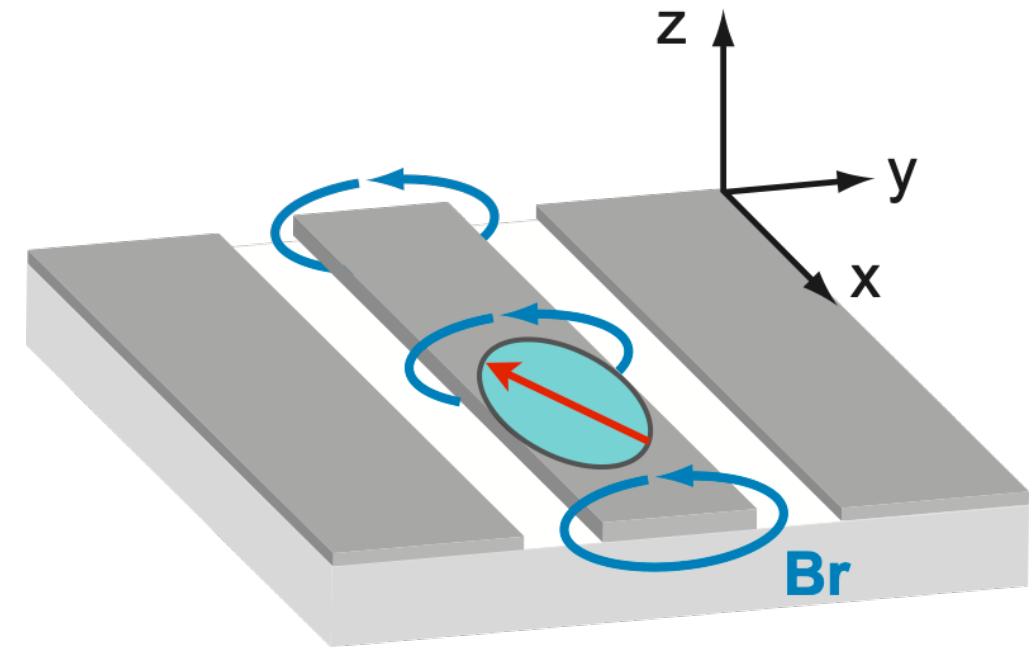


P (Bi) donors in Silicon



Coupling a single spin to a resonator

- magnetic dipole coupling constant:
$$g/2\pi = \frac{g_{NV}\mu_B\delta B_0}{h}$$
- Assuming ‘conventional’ coplanar waveguide geometry,
- Single spin coupling:
 $\sim 1 - 10 \text{ Hz}$



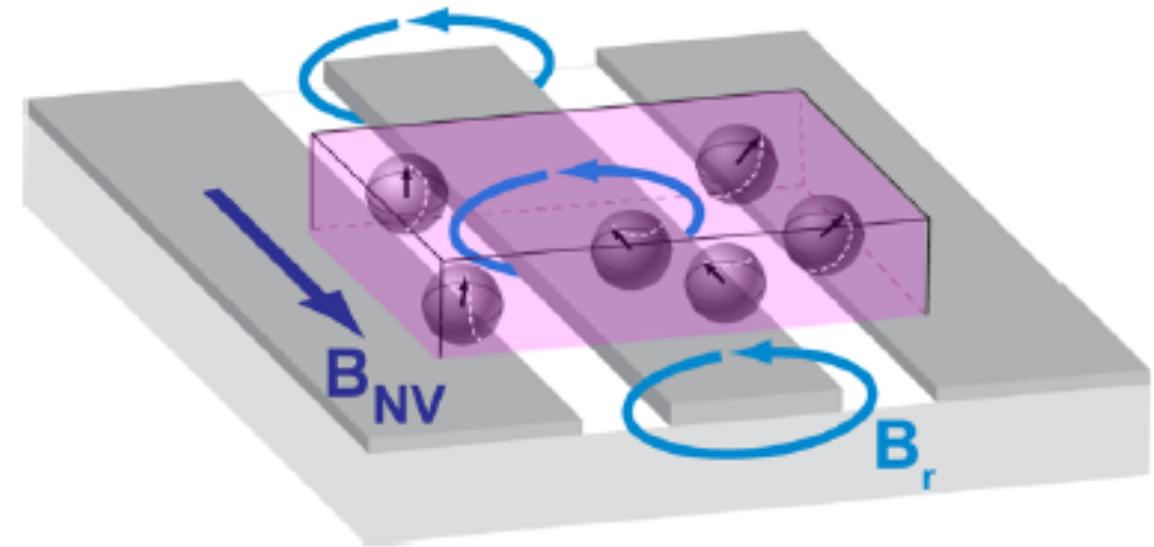
Coupling many spins to a resonator

- Coupling with N spins:

$$g_{ens} \sim \sqrt{N} g$$

- Coupling $N = 10^{12}$ ($\sim 10^{18} \text{ cm}^{-3}$) spins

- Collective coupling of
 $\sim 1 - 10 \text{ MHz}$



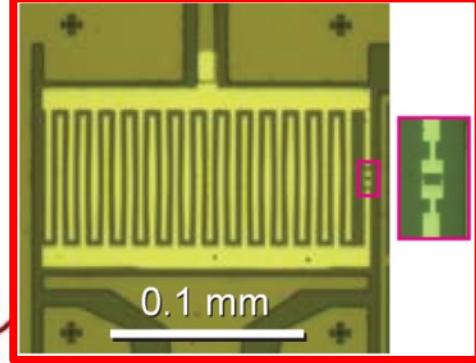
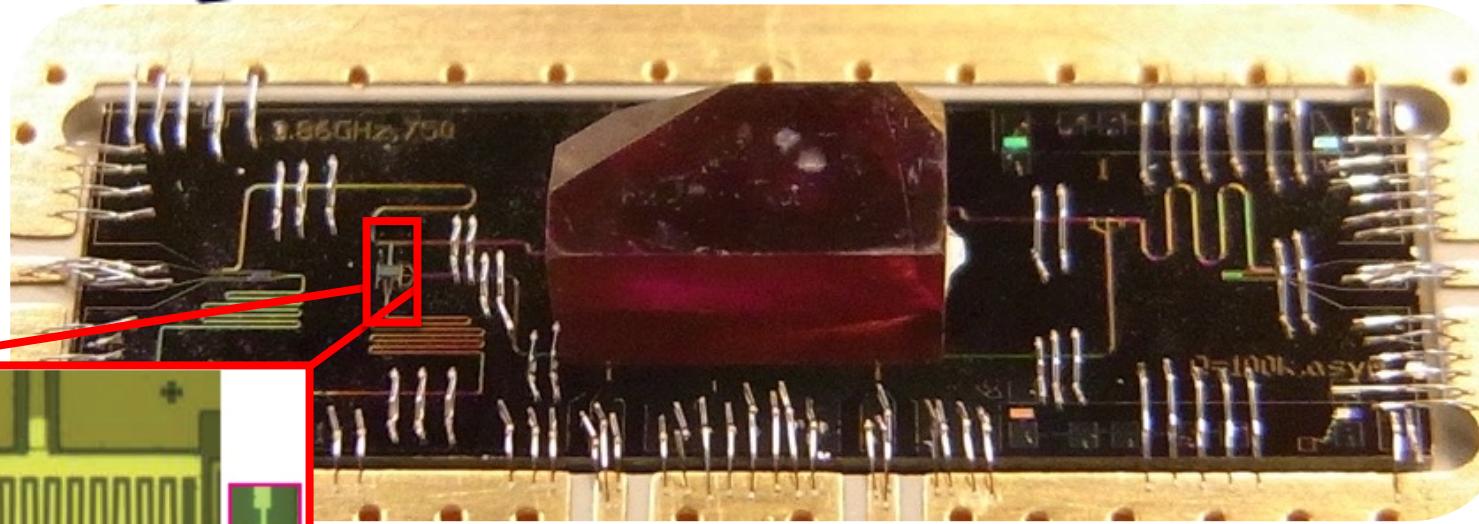
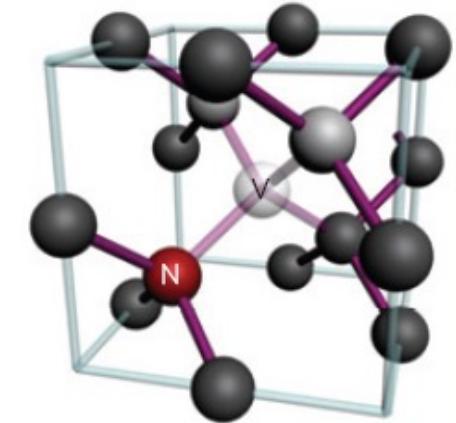
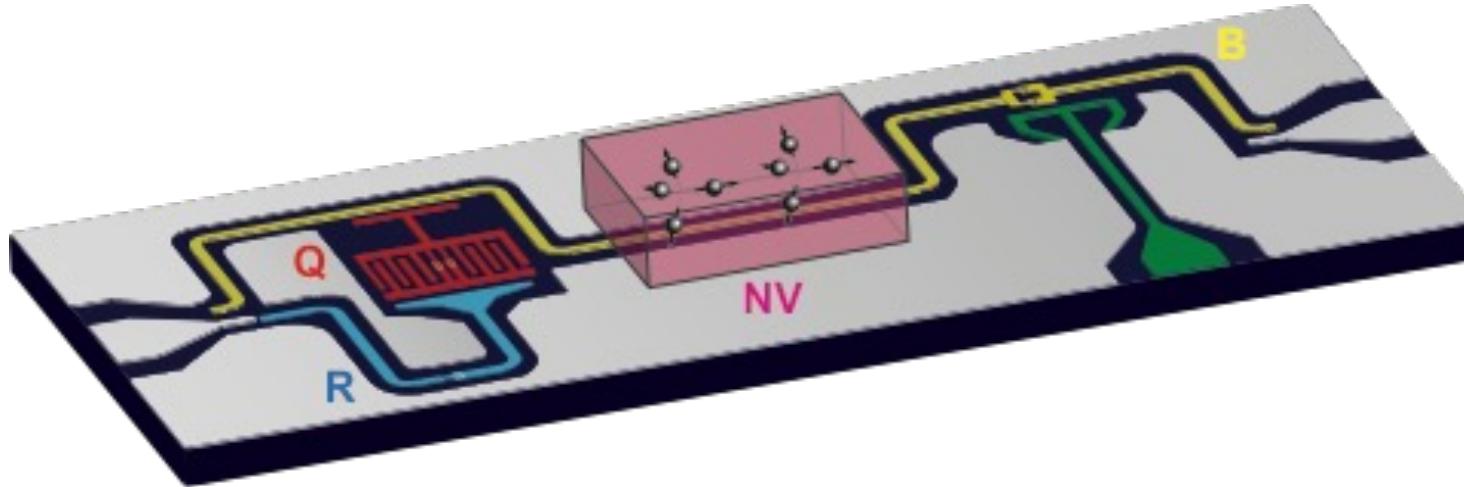
$$H = \sum g_k (a\sigma_{k,+} + a^+\sigma_{k,-})$$

$$\rightarrow H = g_{ens} (aS^+ + a^+S^-)$$

$$\text{with } S = \sum (g_k(r_k)/g_{ens})\sigma_{k,-}$$

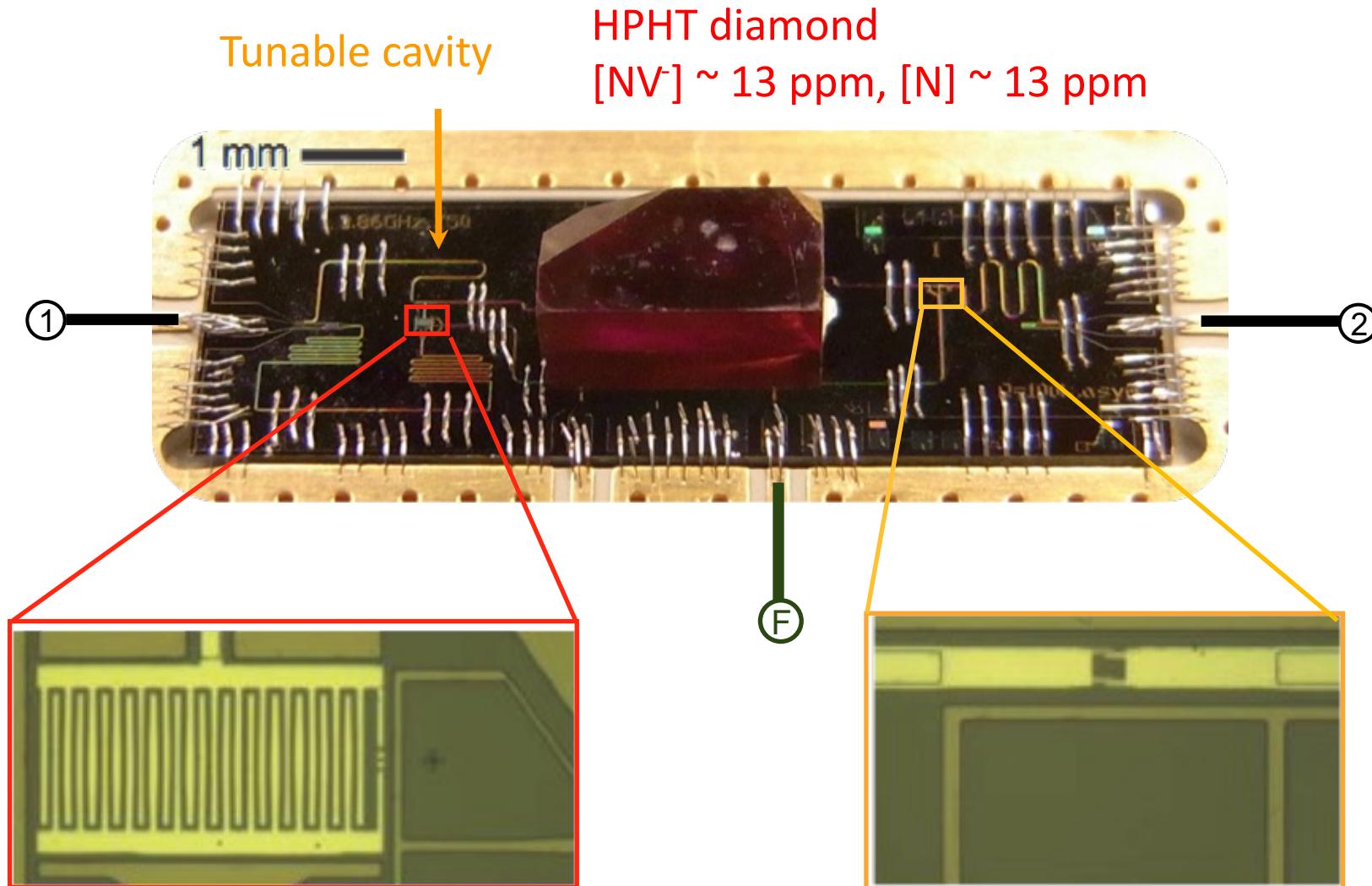
$$|1\rangle \otimes |0_1 0_2 \dots 0_n\rangle \rightarrow |0\rangle \otimes (g_1|1_1 0_2 \dots 0_n\rangle + \dots + g_n|0_1 0_2 \dots 1_n\rangle)/g_{ens}$$

The hybrid quantum circuit

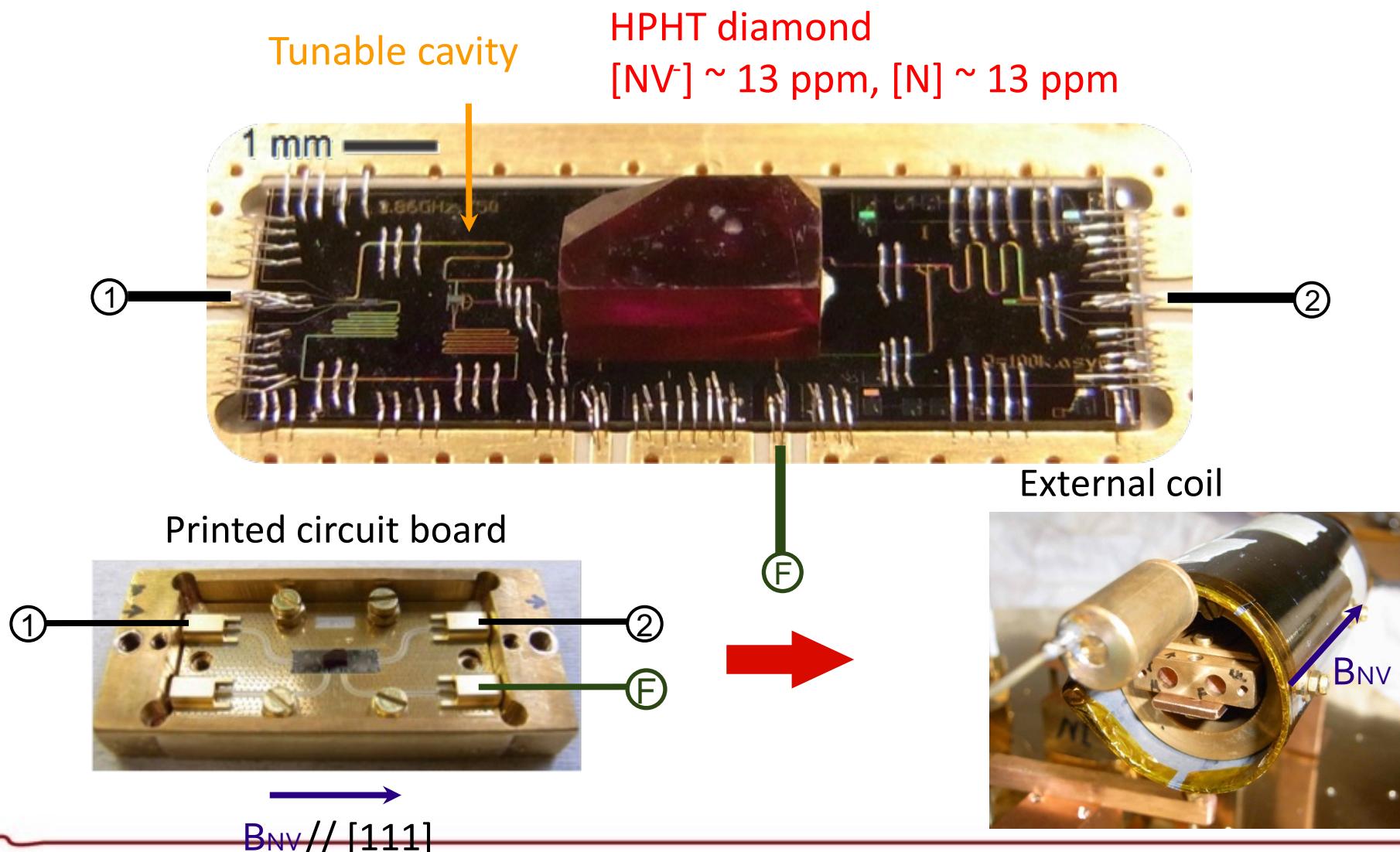


Kubo et al., PRL 105, 140502 (2010) (photon & spins)
Kubo et al., PRL 107, 220501 (2011) (qubit, photon, spins)

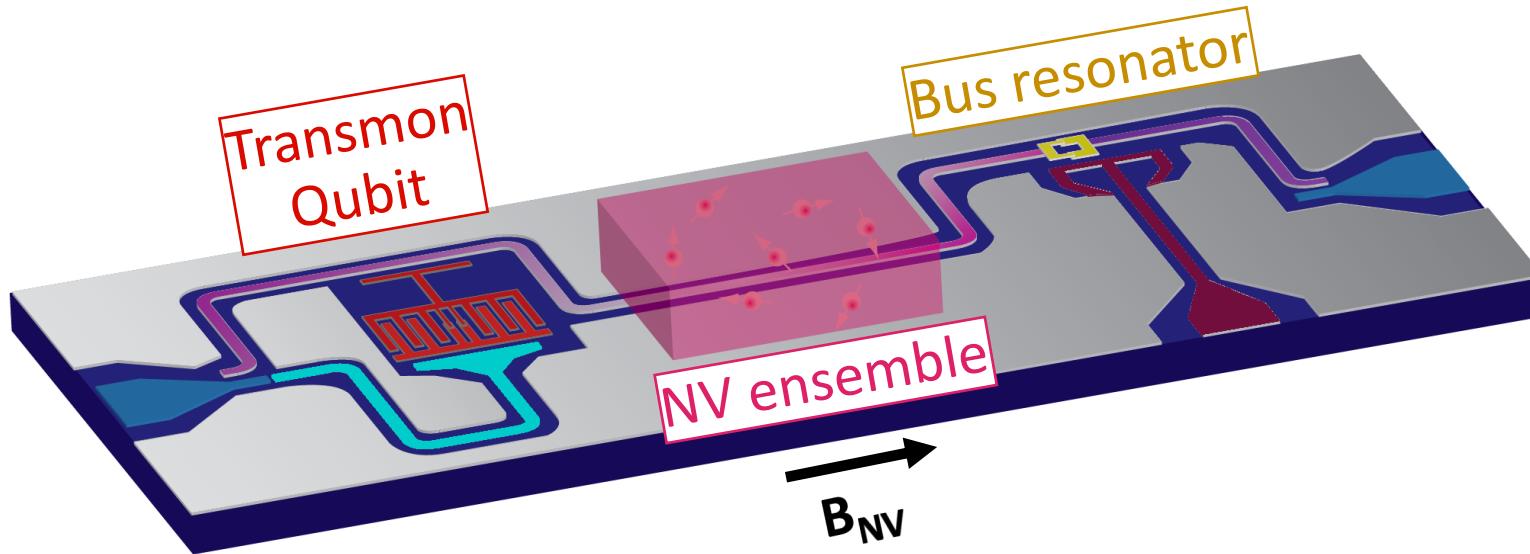
The hybrid quantum circuit



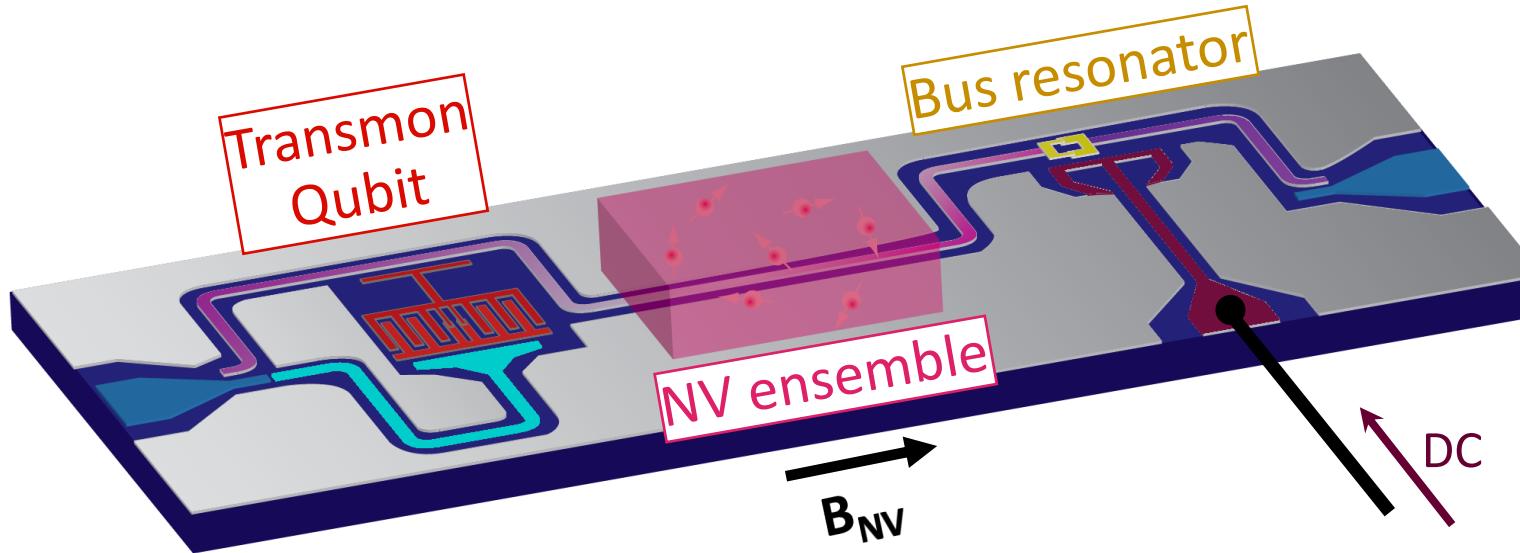
STEP 1 (WRITE): the device



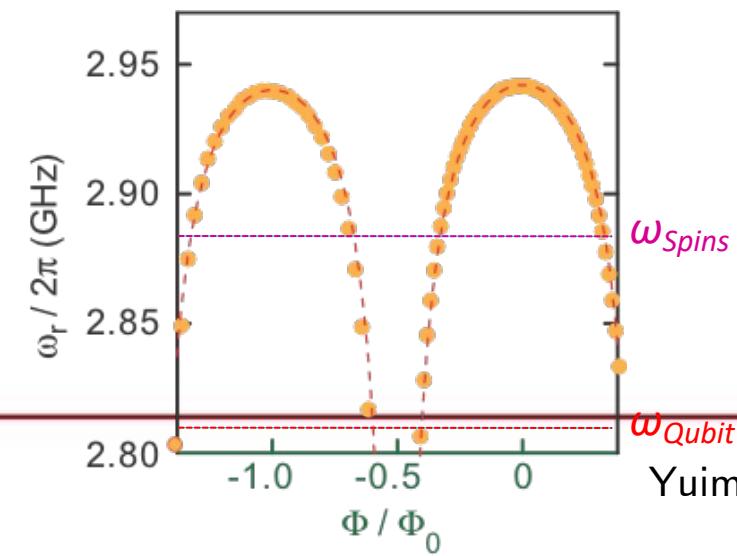
The hybrid quantum circuit



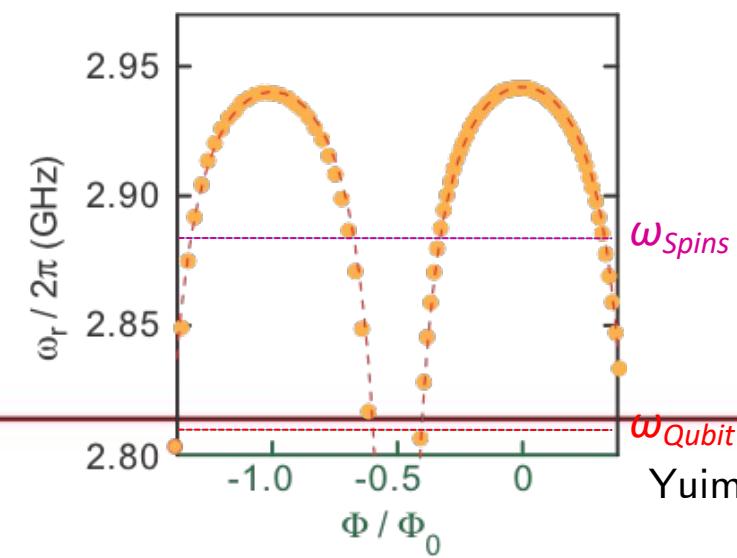
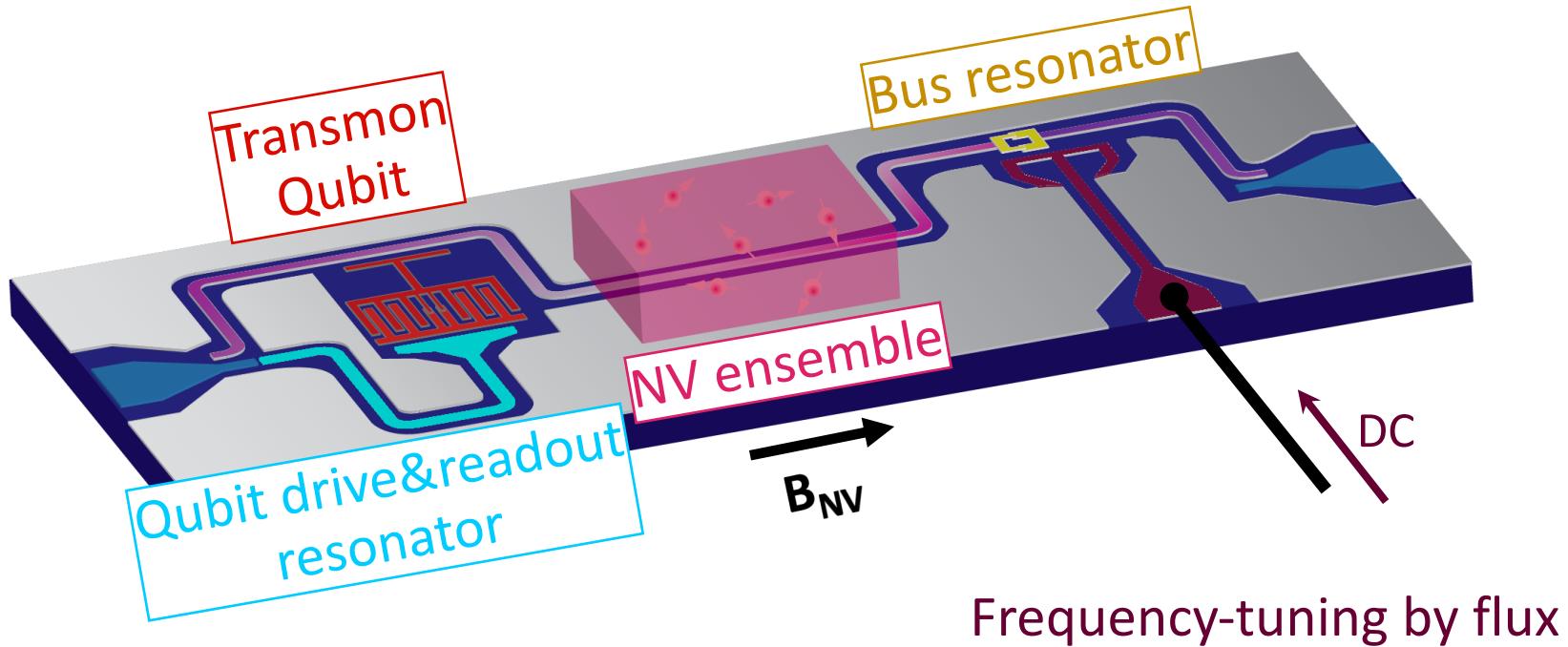
The hybrid quantum circuit



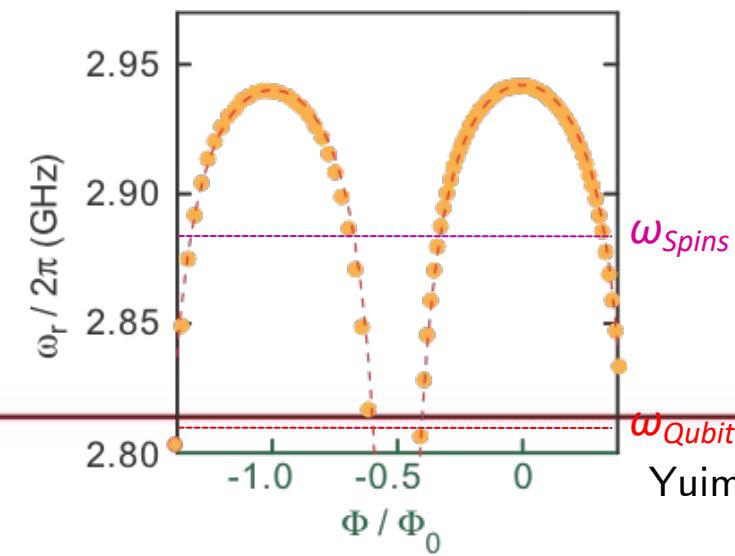
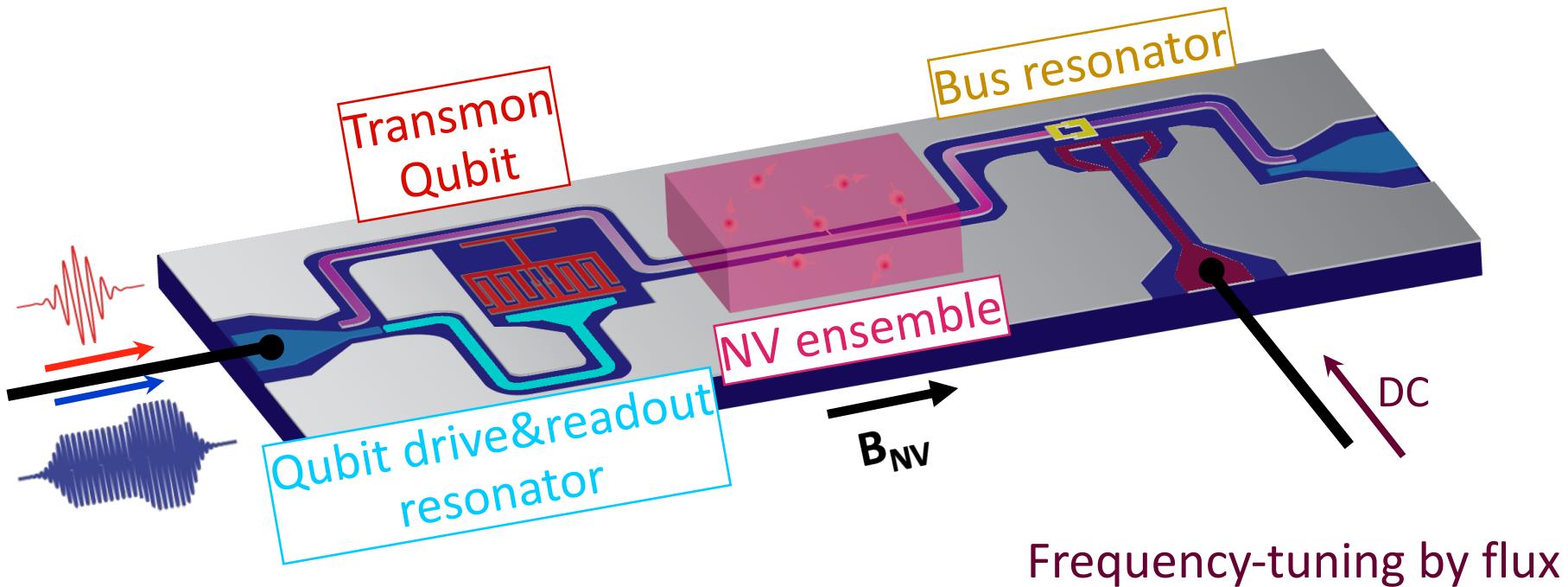
Frequency-tuning by flux



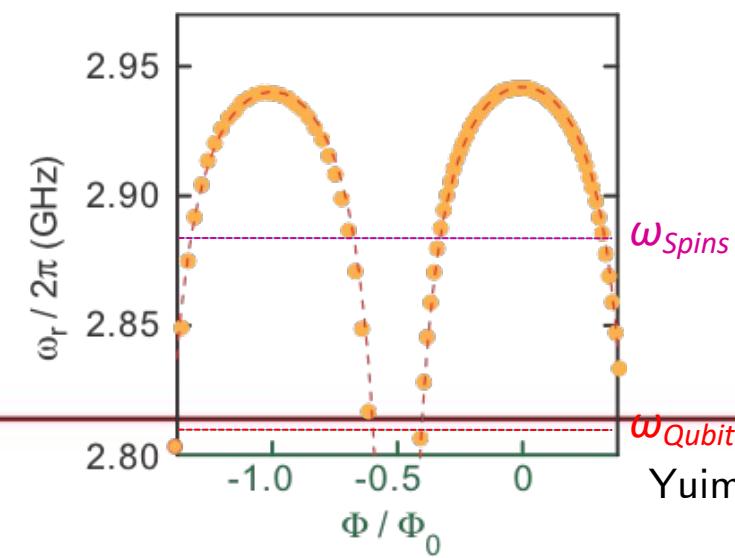
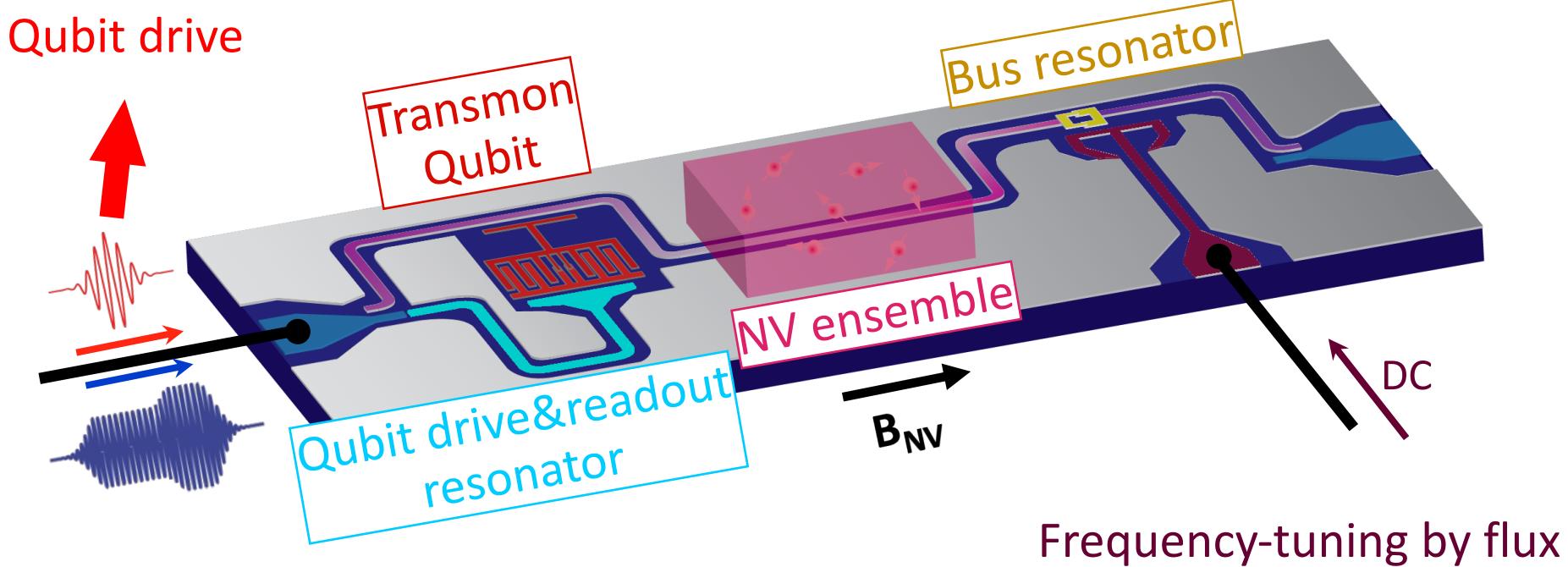
The hybrid quantum circuit



The hybrid quantum circuit

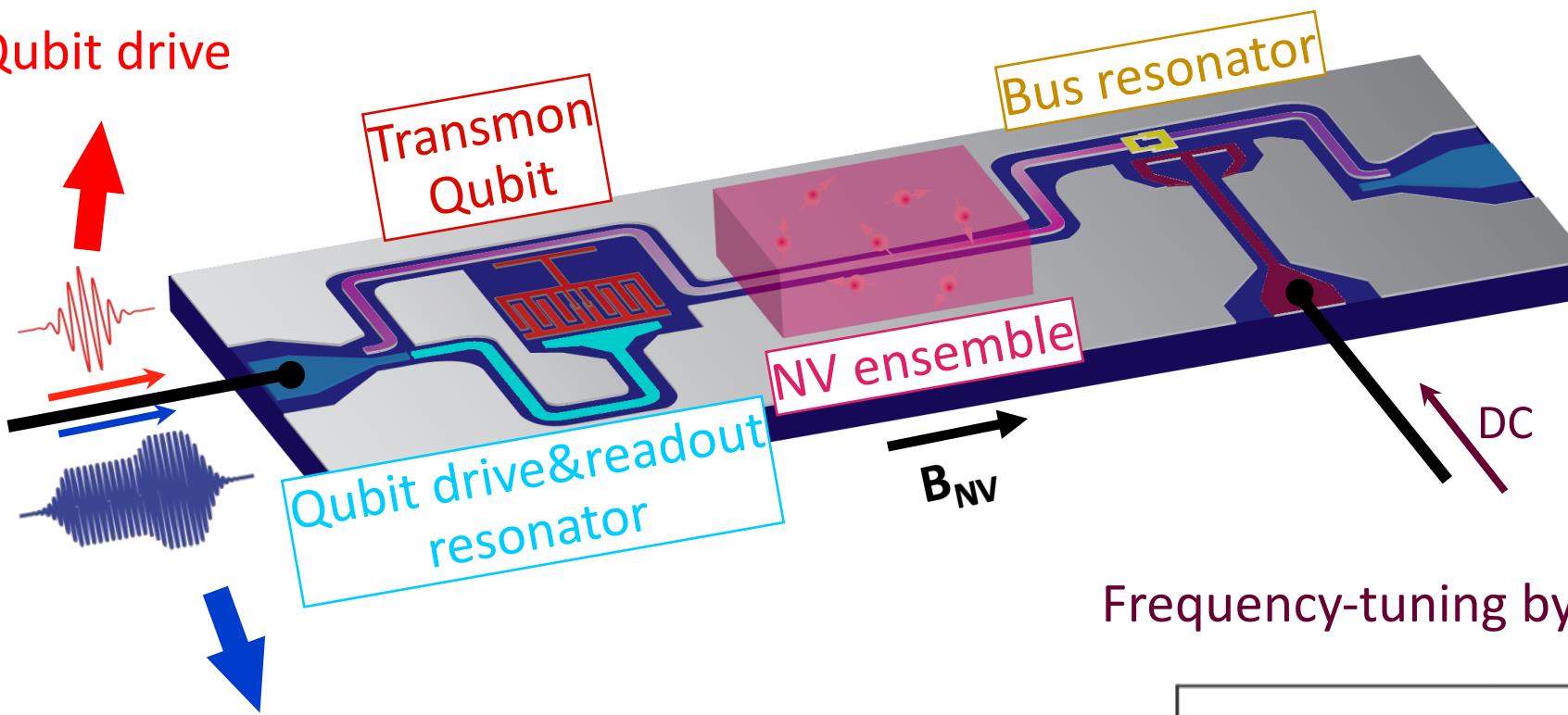


The hybrid quantum circuit

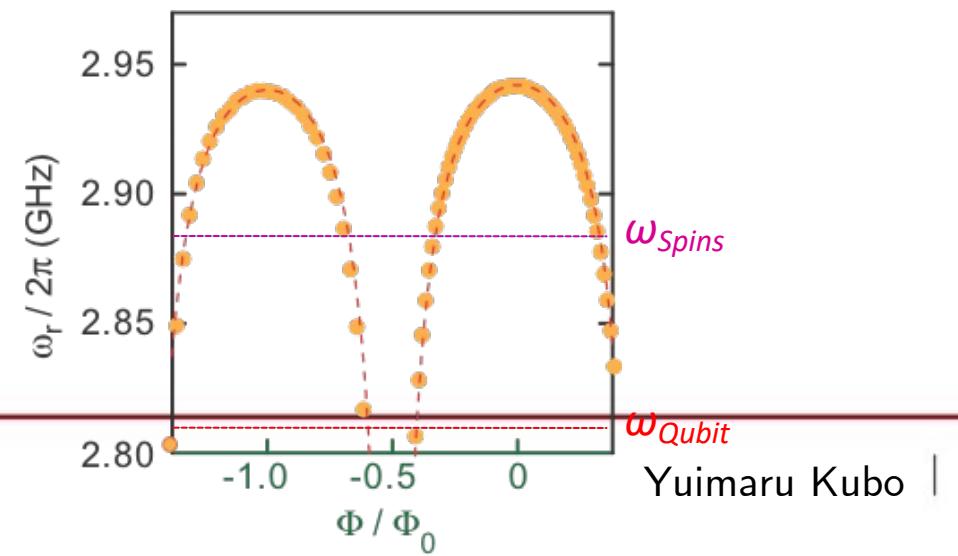
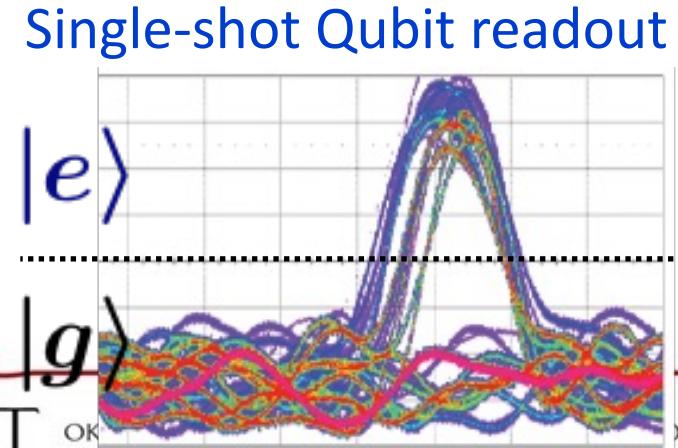


The hybrid quantum circuit

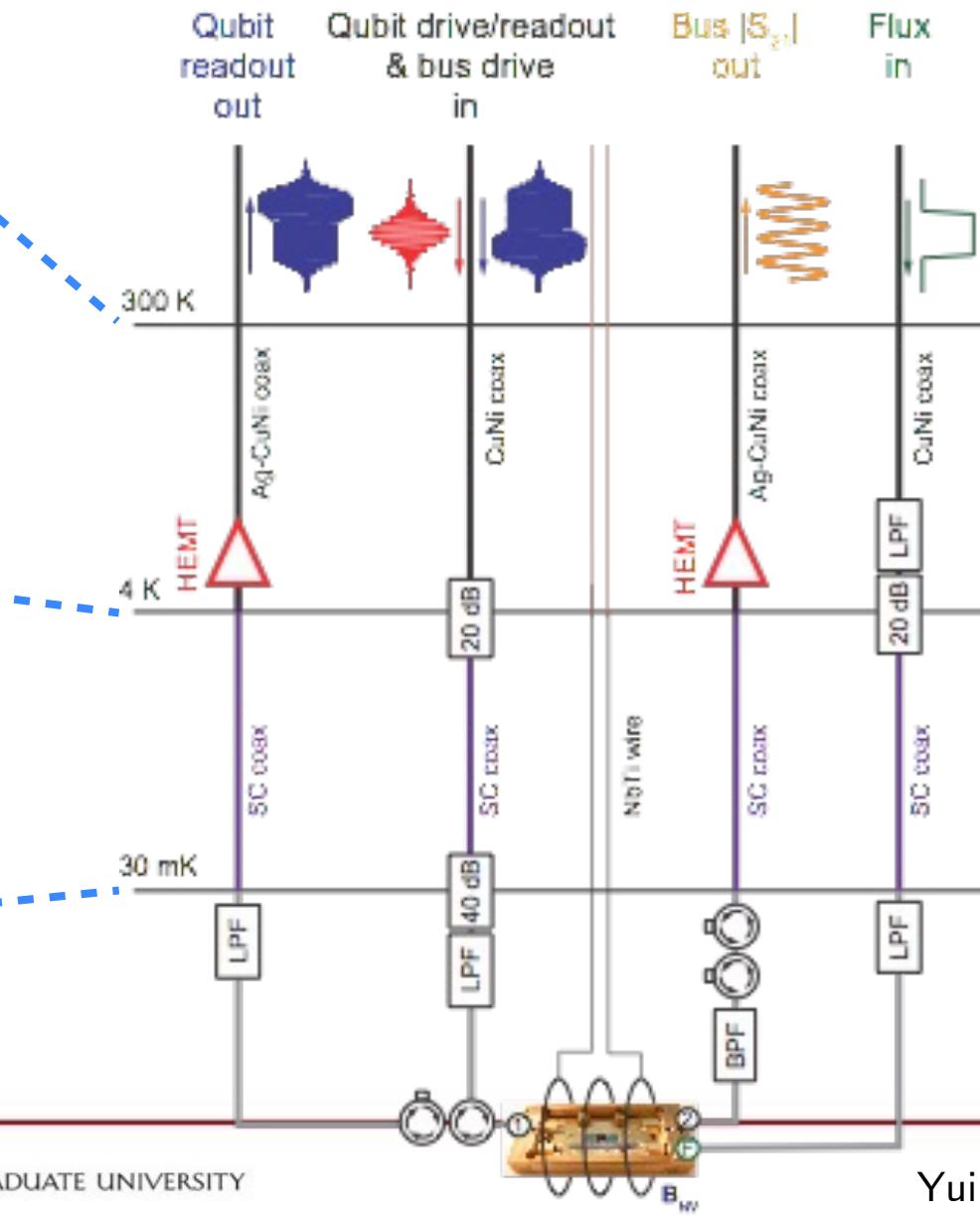
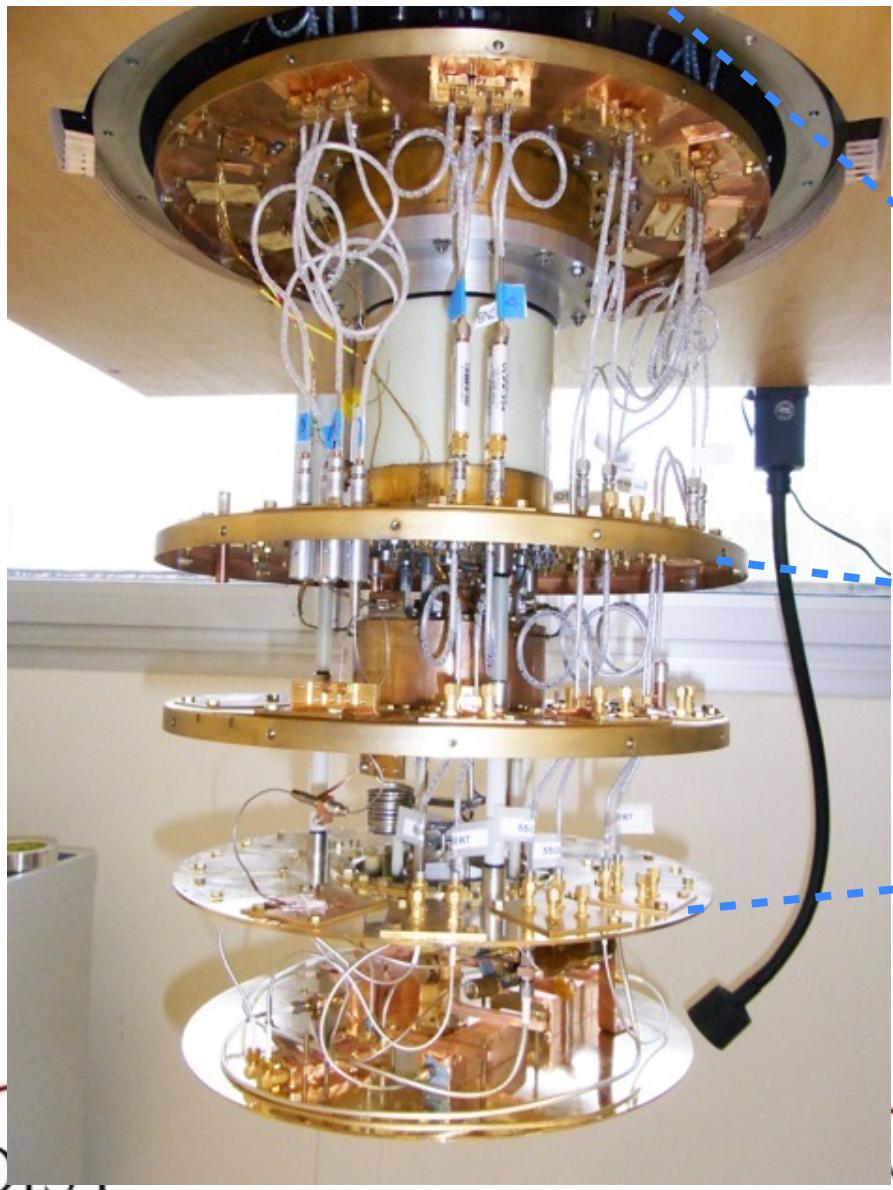
Qubit drive



Frequency-tuning by flux



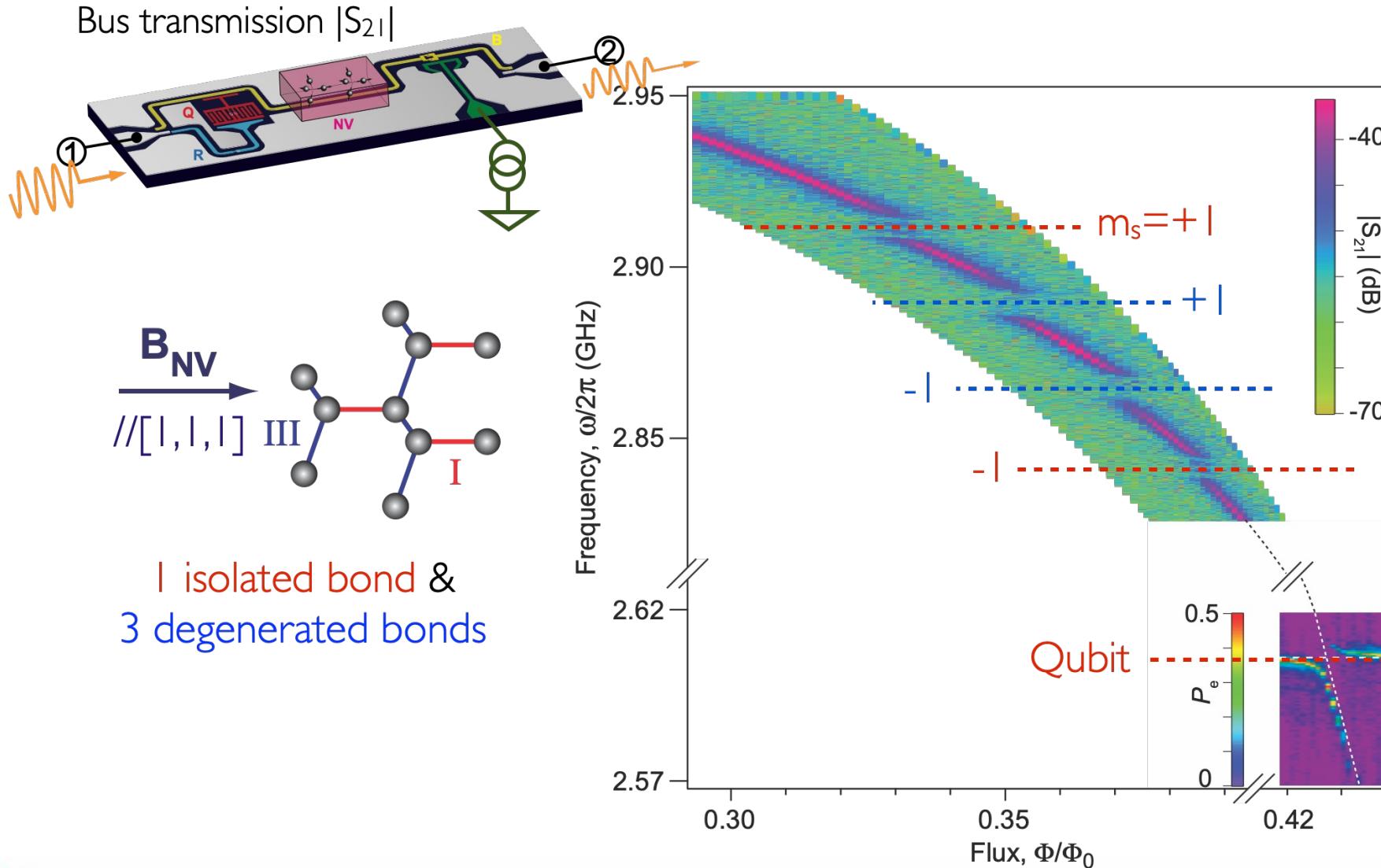
Setup in a dilution fridge



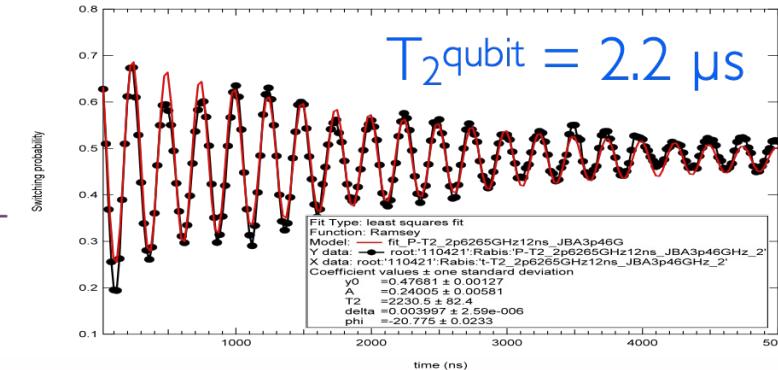
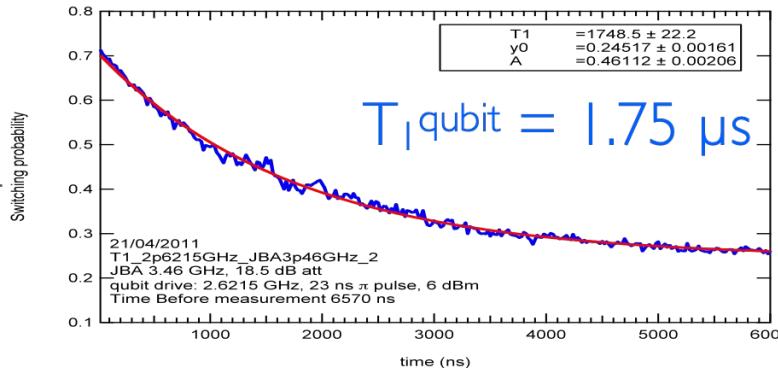
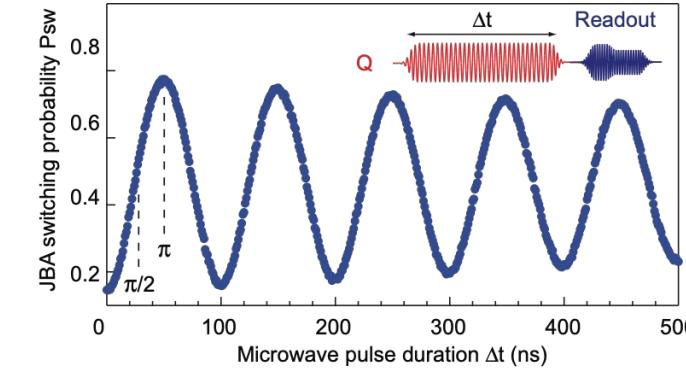
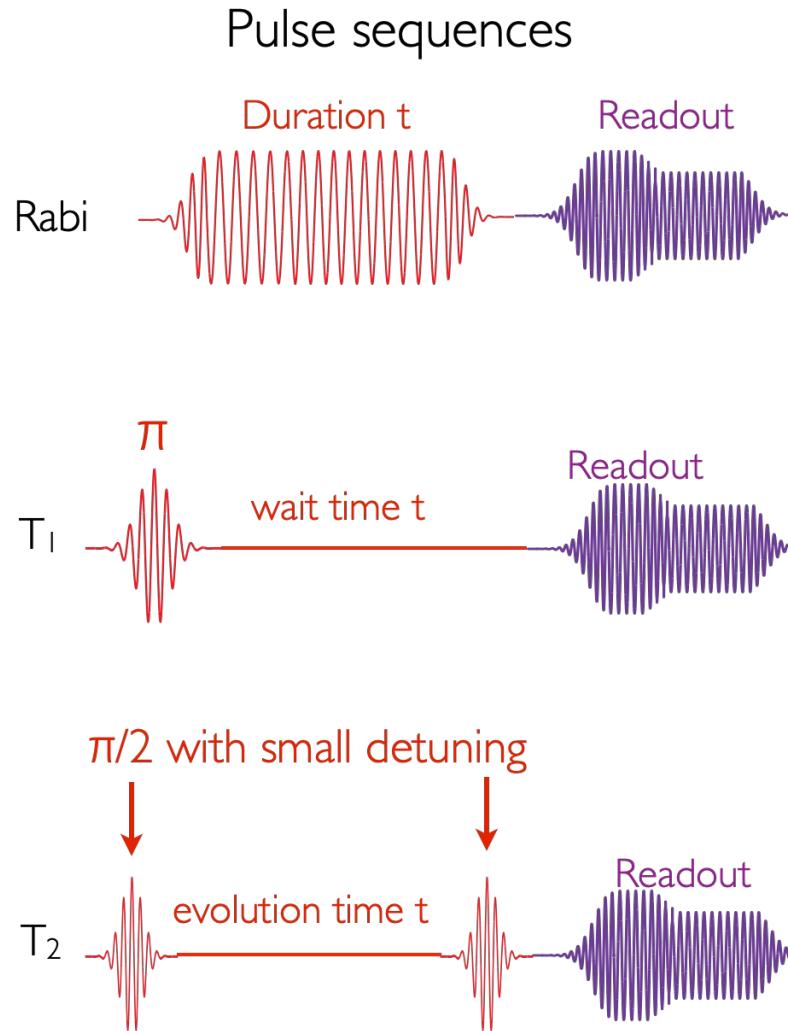
GRADUATE UNIVERSITY

Yuimaru Kubo |

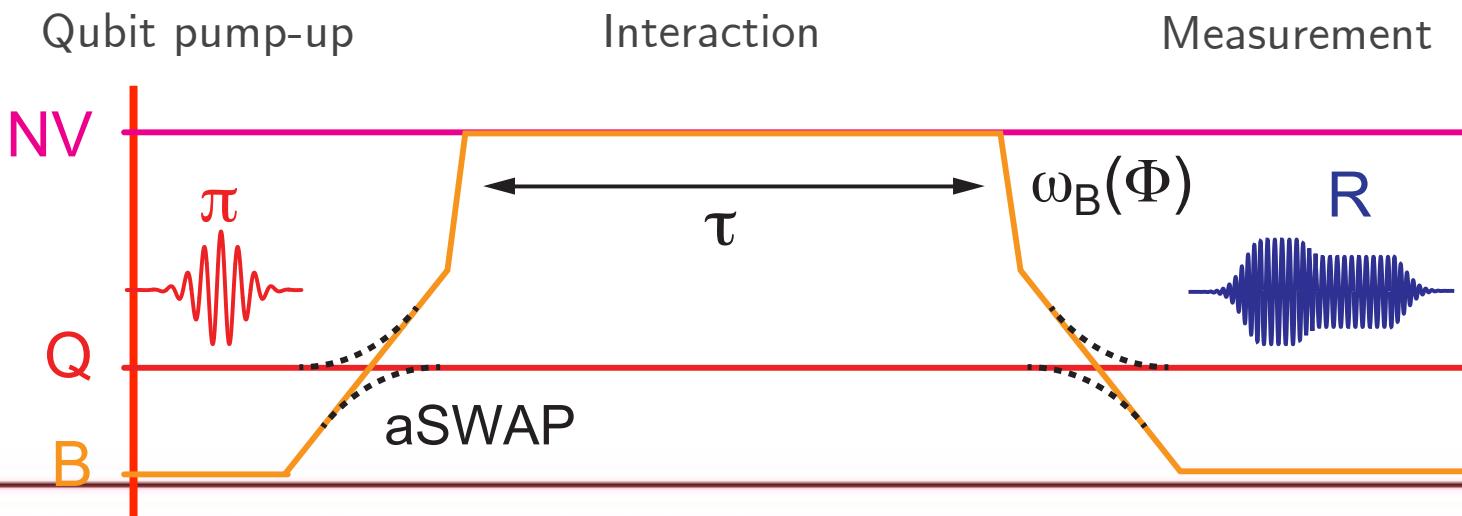
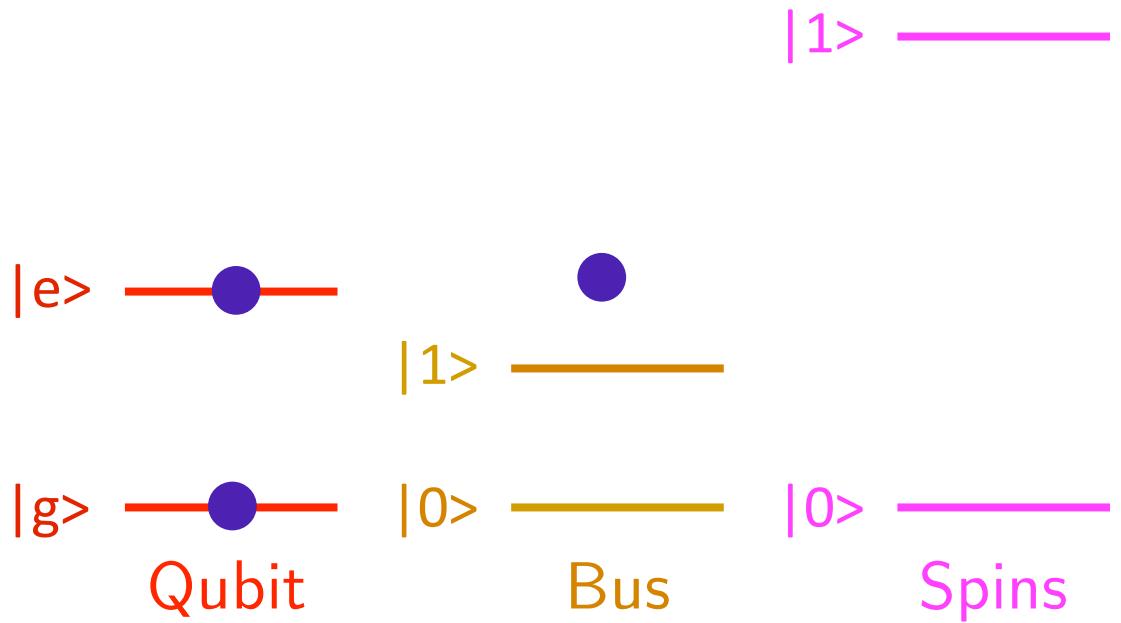
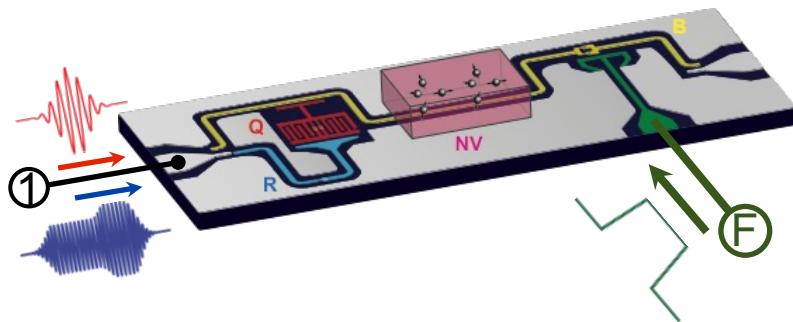
Device characterization: spectroscopy



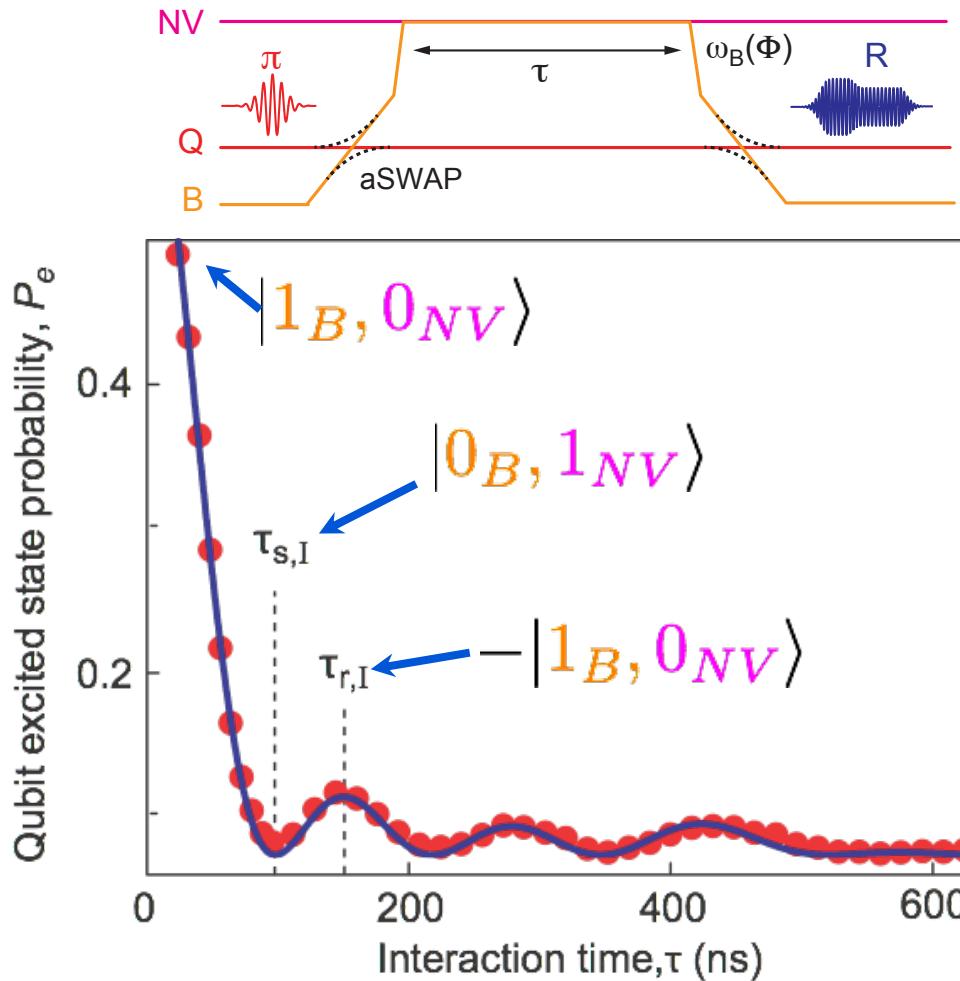
The ‘bad’ qubit (so-so at the time...)



Quantum memory: single μ -wave photon swap



Coupling qubit to spin ensemble: single photon swap



Kubo et al., PRL 107, 220501 (2011)

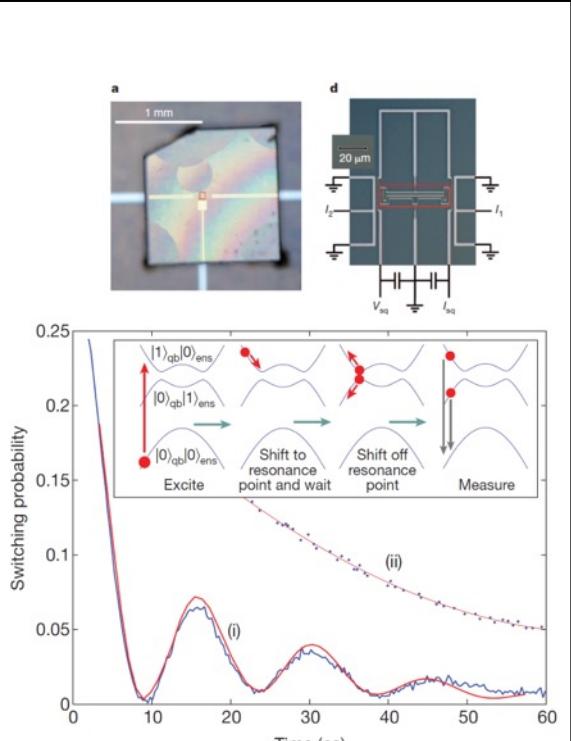
See also: Zhu et al. Nature 478 221 (2011)

Storage/retrieval of a **SINGLE** microwave photon into/from a spin ensemble

Related works

Storage of $|1\rangle$

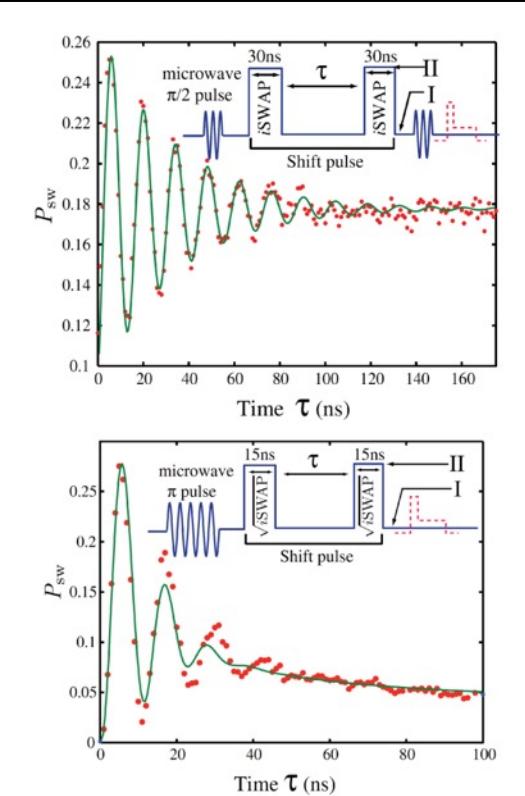
NTT, NV centers



Zhu et al.,
Nature, **478**, 7368 (2011)

Storage of $|0\rangle + |1\rangle$

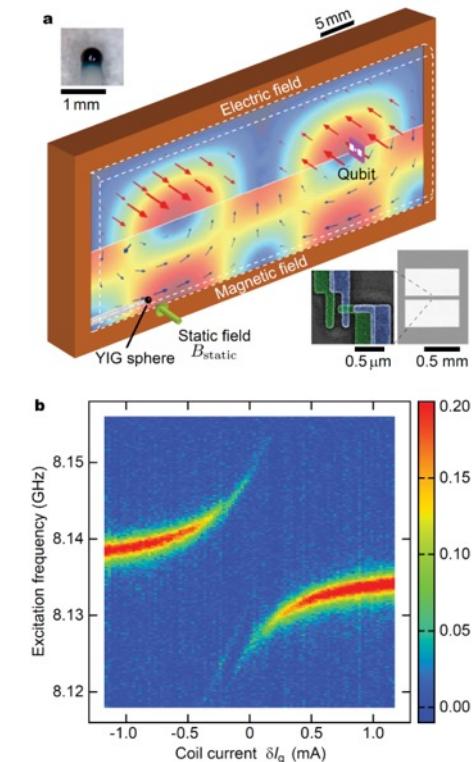
NTT, NV centers



Saito et al.,
PRL, **111** 107008 (2013)

Coherent coupling with a qubit

Tokyo University, YIG



Tabuchi et al.,
Science **349** 405 (2015)

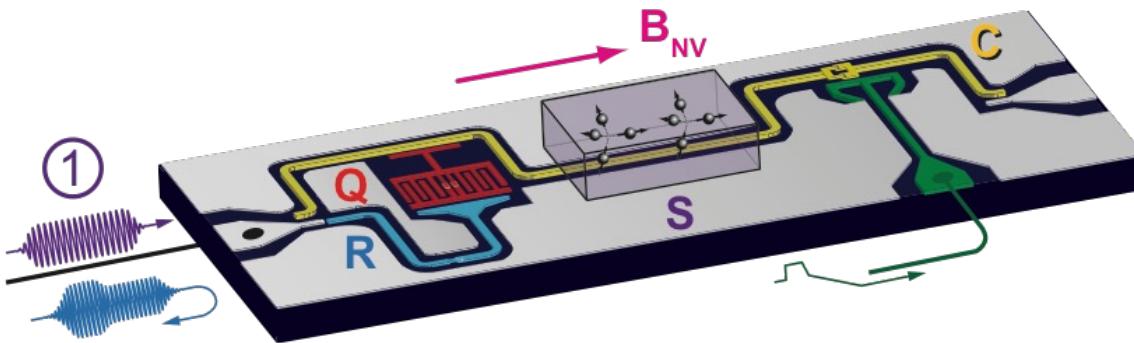


OIST

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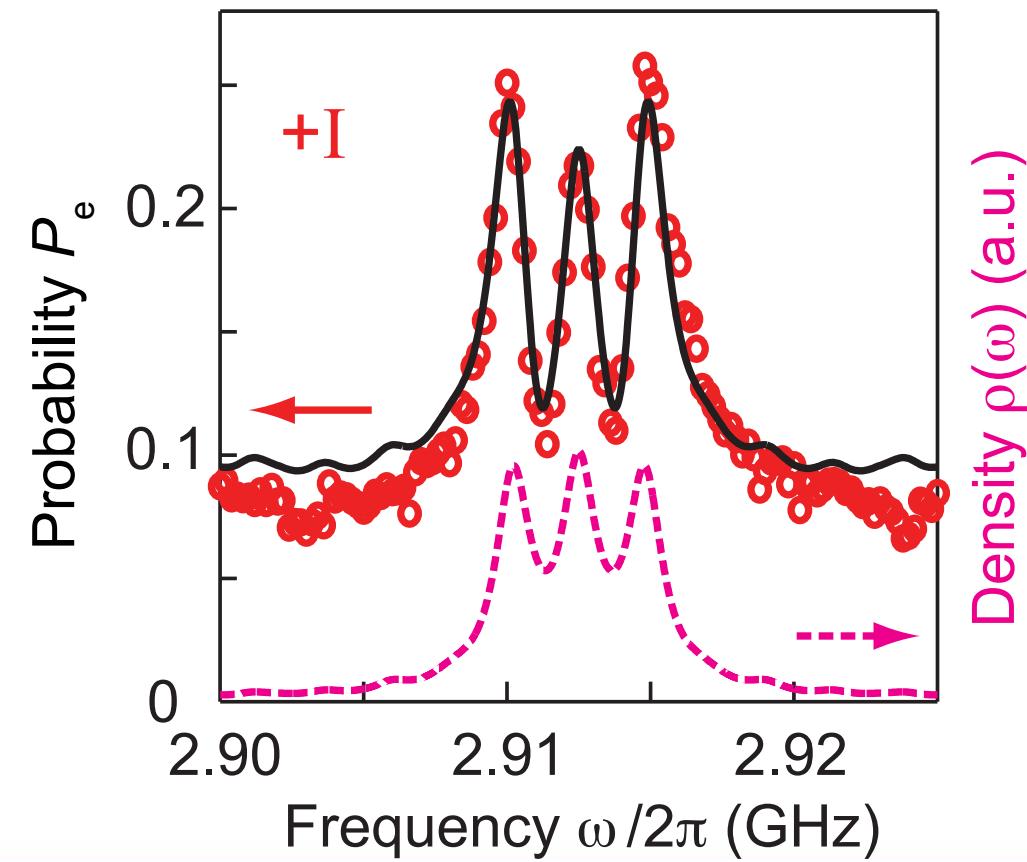
Yuimaru Kubo |

Qubit detected ESR

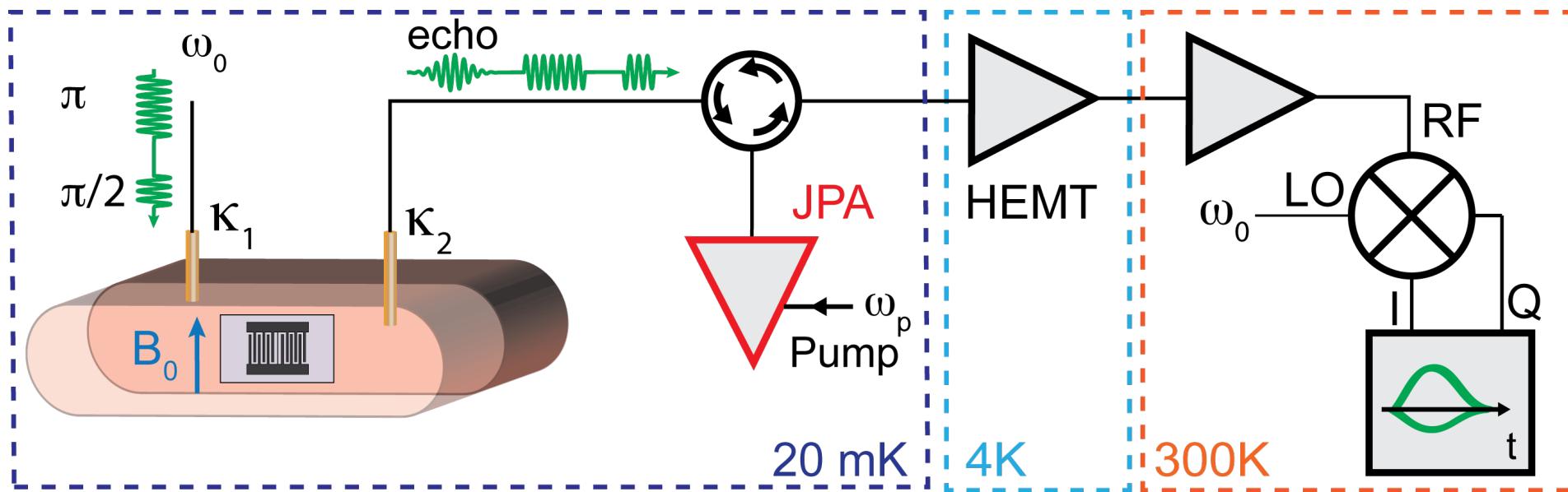


Kubo et al., PRB 86, 064514 (2012)

- Clear ^{14}N hyperfine
- Ultra-low excitation
(≈ 15 spins) was detected
-> 10^5 spins $\text{Hz}^{-1/2}$



Quantum-limited ESR Spectrometer

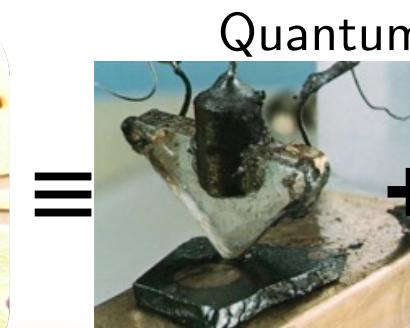
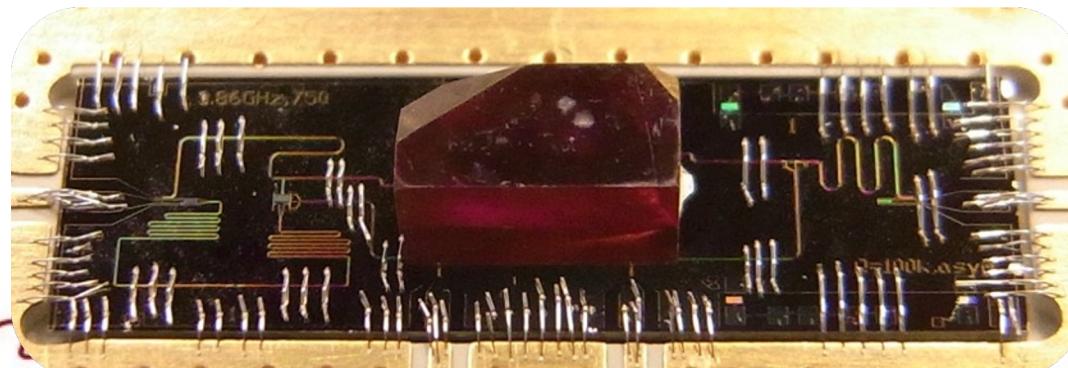


- Full spin polarization @ 20 mK
- High Q (3×10^5) resonator and large coupling (~ 50 Hz)
- Quantum-limited Josephson parametric Amplifier ($T_N \sim < 200$ mK)

Bienfait et al., Nature Nanotech 11, 253 (2016)

Summary & Perspective

- Hybrid quantum circuit with a spin ensemble and superconducting circuit
 - PoC for a superconducting quantum CPU + spin ensemble quantum RAM
 - Quantum RAM operation
- Hybrid circuit as a sensitive magnetization detector
 - Single-spin microwave detection [Wang et al., Nature (2023)]



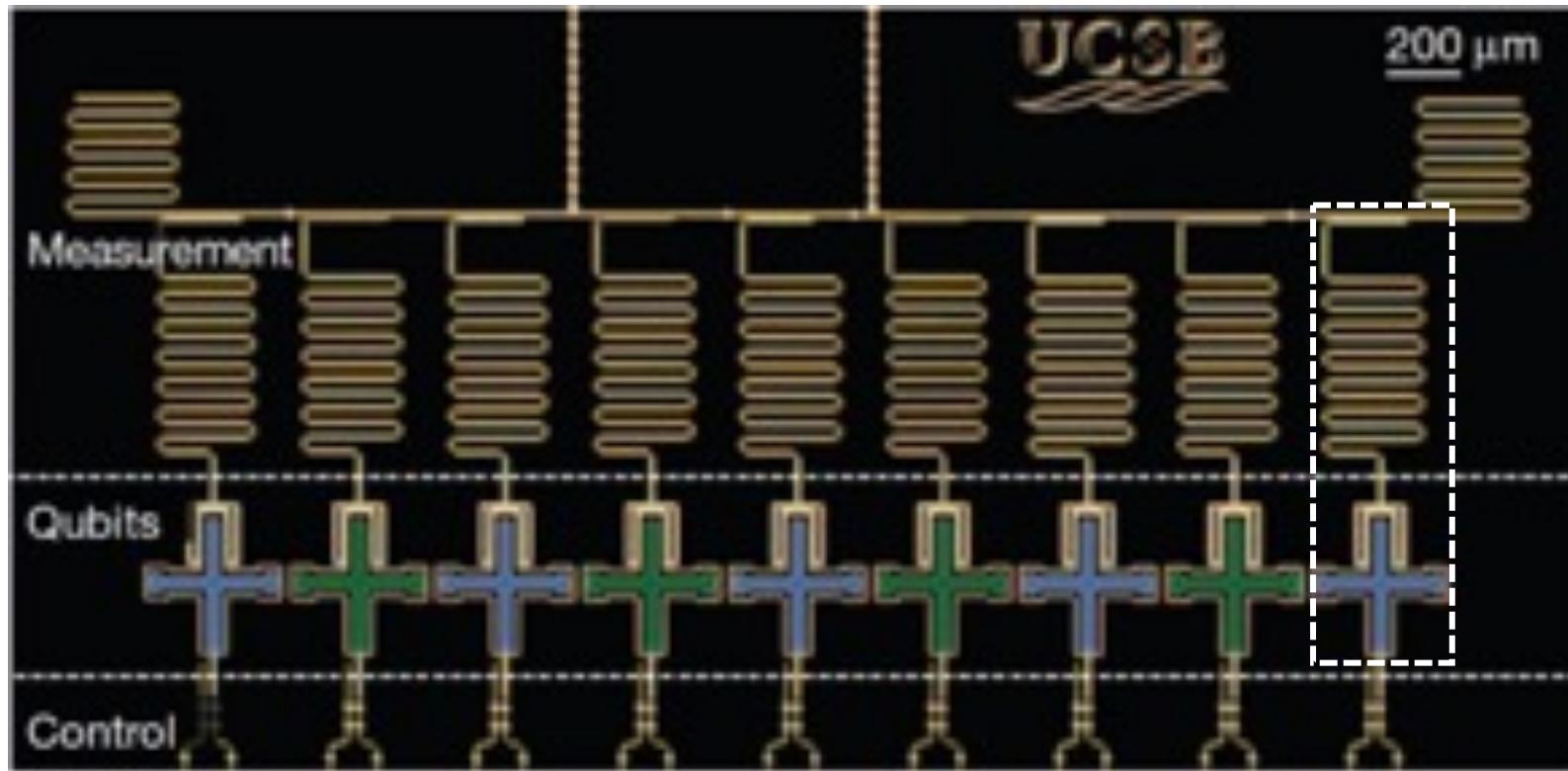
=



Quantum version of...

+

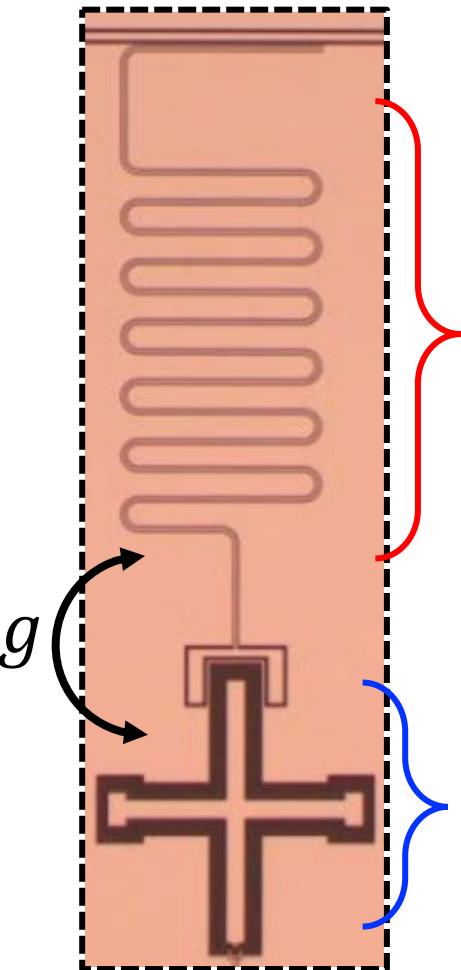
Measuring qubit states with a resonator



Kelly et al., Nature 519, 66 (2015)

$$\mathcal{H}_{JC}/\hbar = \left(\omega_r + \frac{g^2}{\omega_r - \omega_q} \sigma_z \right) (\hat{a}^\dagger \hat{a} + 1/2) + \frac{\omega_q \sigma_z}{2}$$

(Dispersive regime: $|\omega_r - \omega_q| \gg g$)

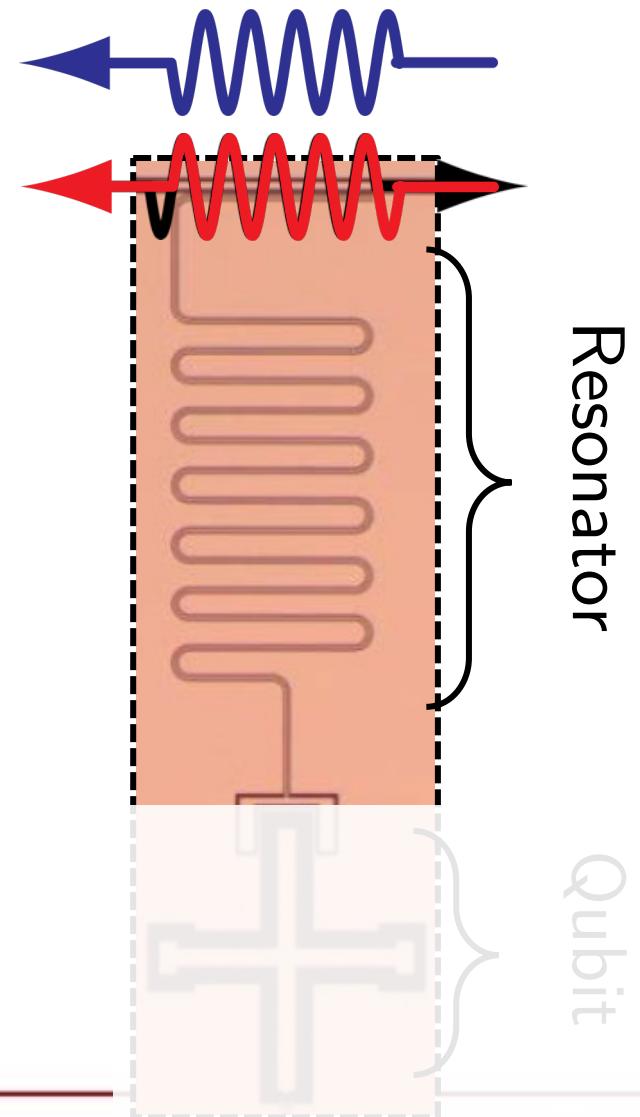
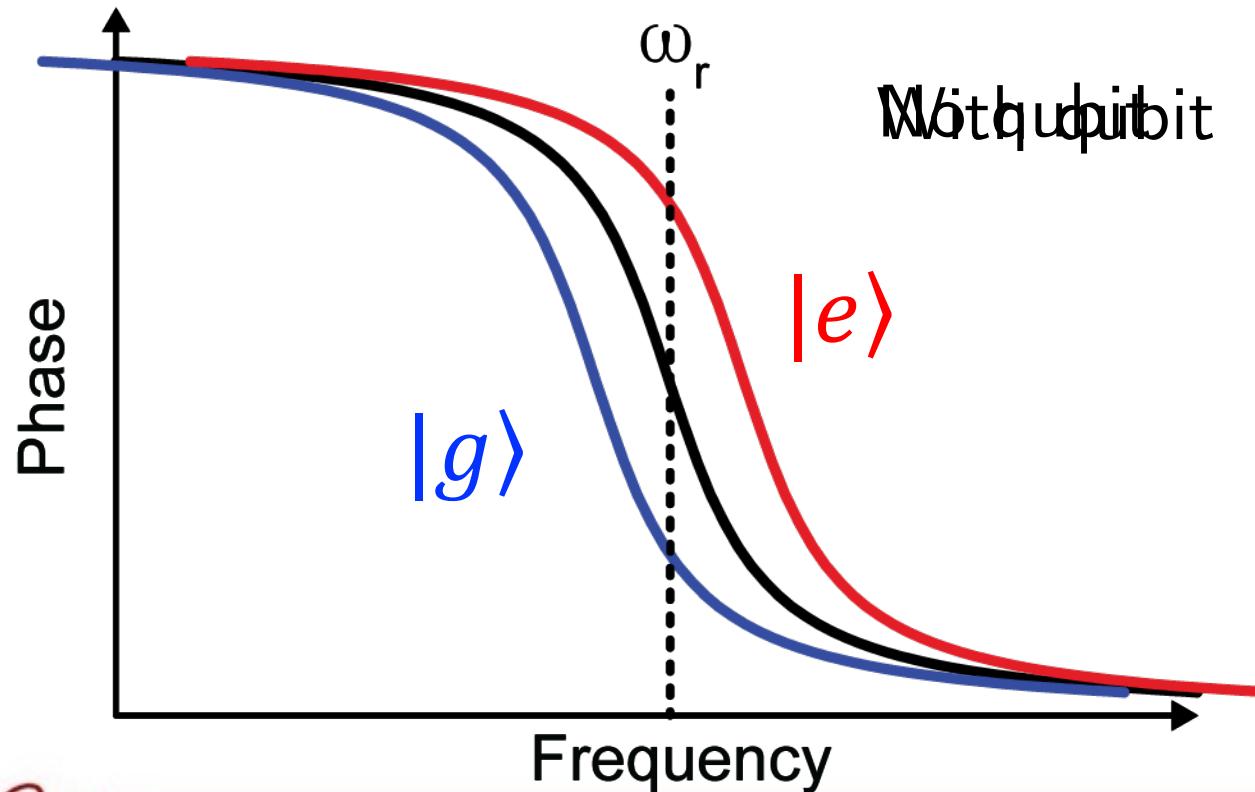


Resonator
Qubit

Measuring qubit states with a resonator

$$\mathcal{H}_{JC}/\hbar = \left(\omega_r + \frac{g^2}{\omega_r - \omega_q} \sigma_z \right) (\hat{a}^\dagger \hat{a} + 1/2) + \frac{\omega_q \sigma_z}{2}$$

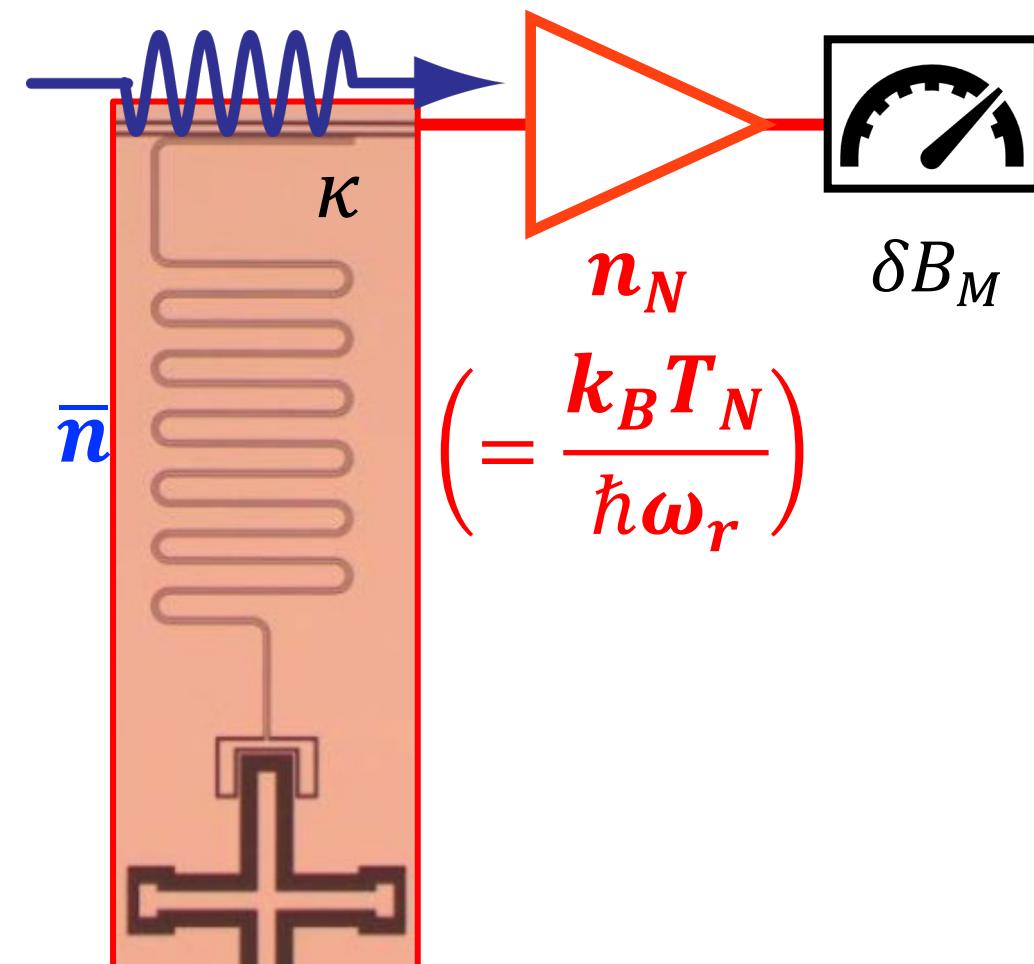
"Dispersive" shift depending on qubit state



Requirements for single-shot readout

$$\text{SNR} \sim \sqrt{\frac{\bar{n}\kappa}{n_N\delta B_M}}$$

- Measurement bandwidth $\delta B_M \gtrsim 10 \text{ MHz}$
- Resonator leak: $\kappa \lesssim 10 \text{ MHz}$
- **Low power of the readout pulse:**
 $\bar{n} \sim 10 \text{ photons} \equiv \lesssim \text{fW}$
- **Small noise from the amplifier:**
 $n_N \sim 1 (\equiv \lesssim 250 \text{ mK})$



Noiseless (quantum-limited) microwave amplifier needed!!

Commercial cryogenic amplifier

Low Noise Factory
(<http://www.lownoisefactory.com>)



- Cryogenic HEMT (high electron mobility transistor) amplifier
- $T_N \approx 3 \text{ K} \gg \hbar\omega/k_B = 250 \text{ mK}$
- 1-10 mW dissipation $\gg \sim 100 \mu\text{W}$
- How to implement a **quantum-limited amplifier** at microwave frequency **at millikelvin?**
(hopefully without much dissipation)

State-of-the-art: Josephson parametric amplifier

- **Josephson junction:**
Non-linear & non-dissipative inductance
- Quantum limited noise
- Wideband ($\sim < 1$ GHz)
- Versatile superconducting quantum technology

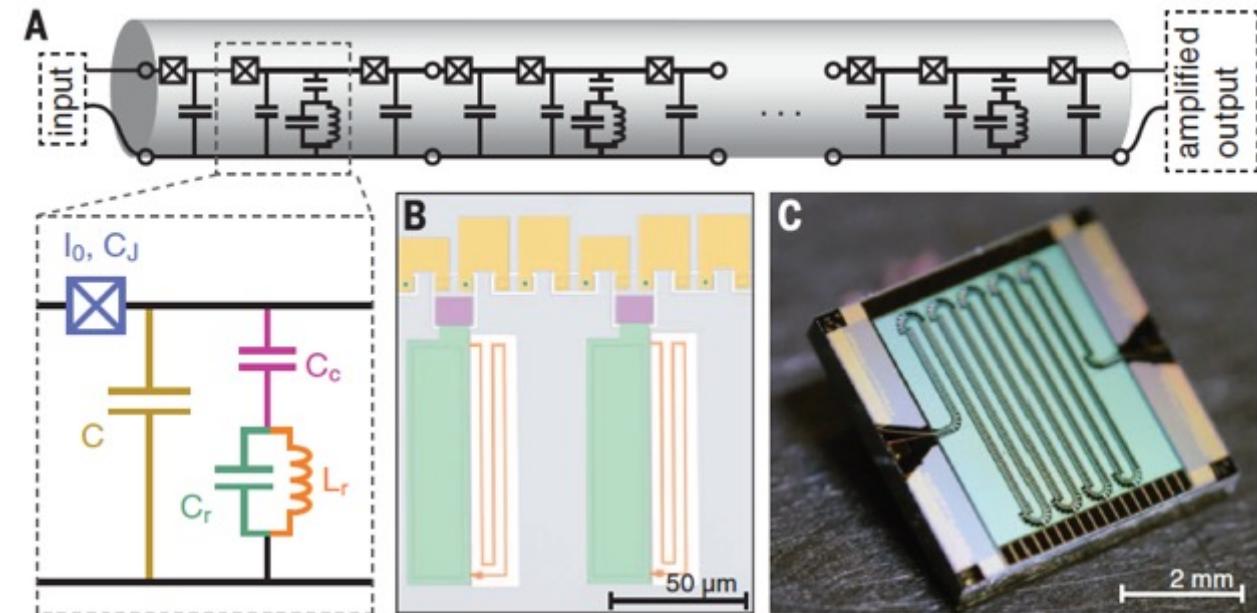
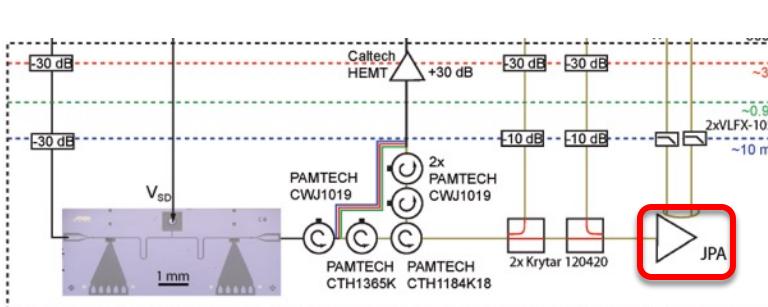


Fig. 1. Josephson traveling-wave parametric amplifier. (A) Circuit diagram. The JTWPA is implemented as a nonlinear lumped-element transmission line; one unit cell consists of a Josephson junction with critical current $I_0 = 4.6 \mu\text{A}$ and intrinsic capacitance $C_J = 55 \text{ fF}$ with a capacitive shunt to ground $C = 45 \text{ fF}$. Every third unit cell includes a lumped-element resonator designed with capacitance $C_r = 6 \text{ pF}$ and inductance $L_r = 120 \text{ pH}$, with coupling strength set by a capacitor $C_c = 20 \text{ fF}$. The value of C in the resonator-loaded cell is reduced to compensate for the addition of C_c . (B) False-color optical micrograph. The coloring corresponds to the inset in (A), with the lower metal layer shown in gray. (C) Photograph of a 2037 junction JTWPA. The line is meandered several times on the 5 mm by 5 mm chip to achieve the desired amplifier gain.

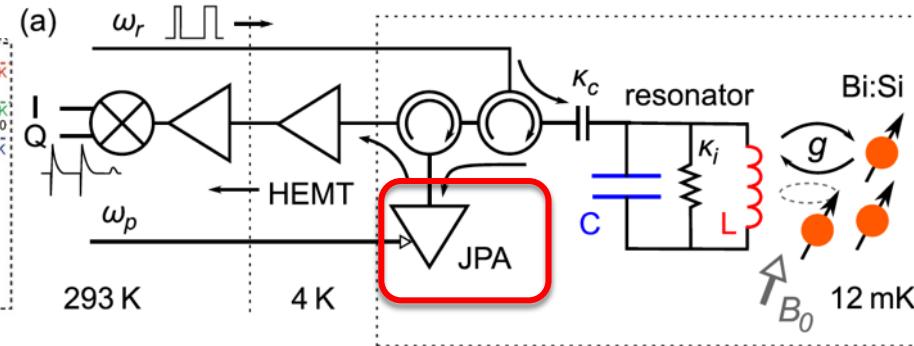
Macklin et al., Science 350, 307 (2015).

Paramp as microwave quantum technology

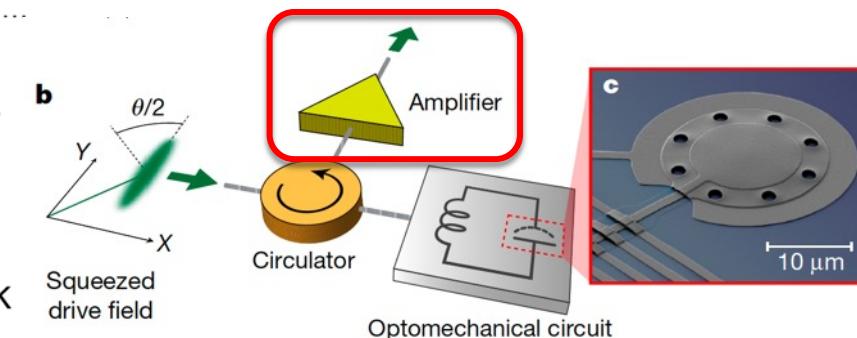
Quantum-dot qubits



Spin resonance



Mechanical motion



Liu et al., Science **347**, 285 (2015)
Mi et al., Nature **555**, 599 (2018)

Eichler et al., PRL **118**, 037701 (2017)
Bienfait et al., Nature Nano **11**, 253 (2016)
Probst et al., APL **111**, 202604 (2017)

Clark et al., Nature **541**, 191 (2017)
Satzinger et al., Nature **563**, 661 (2018)
Noguchi et al., PRL **119**, 180505 (2017)

Drawbacks of Josephson parametric amplifier

- Limited dynamic range
(currently reported:
 $-100 \text{ dBm} = 0.1 \text{ pW}$)
-> readout **ONLY 20 qubits simultaneously**
(5 qubits so far)

- Inoperable under
(moderate) magnetic fields

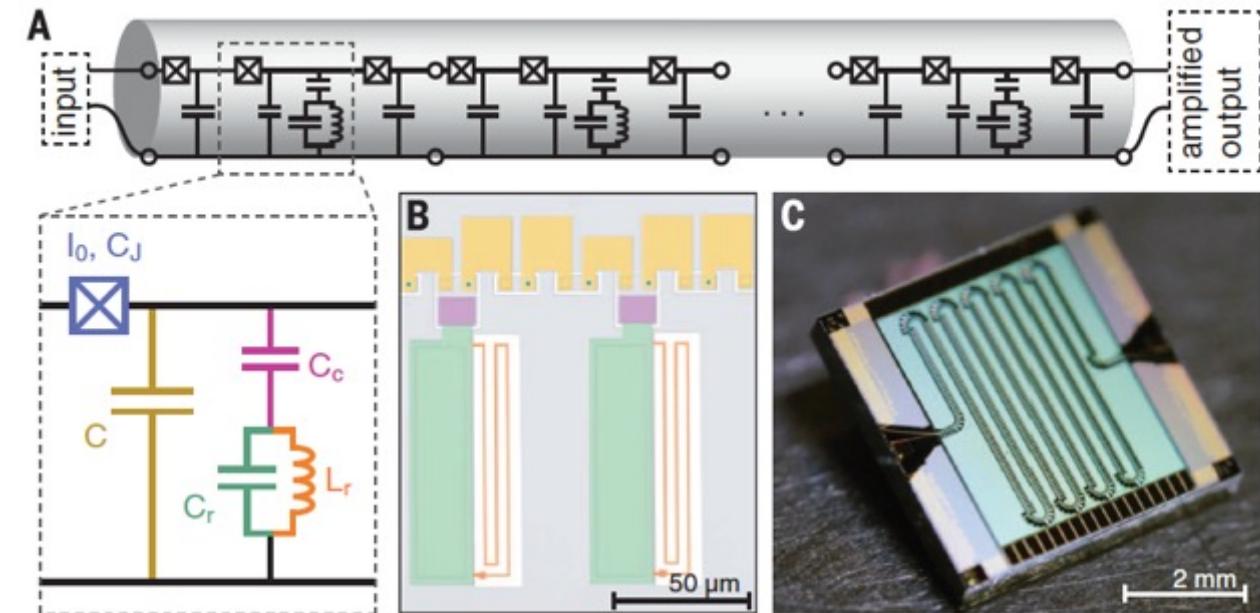


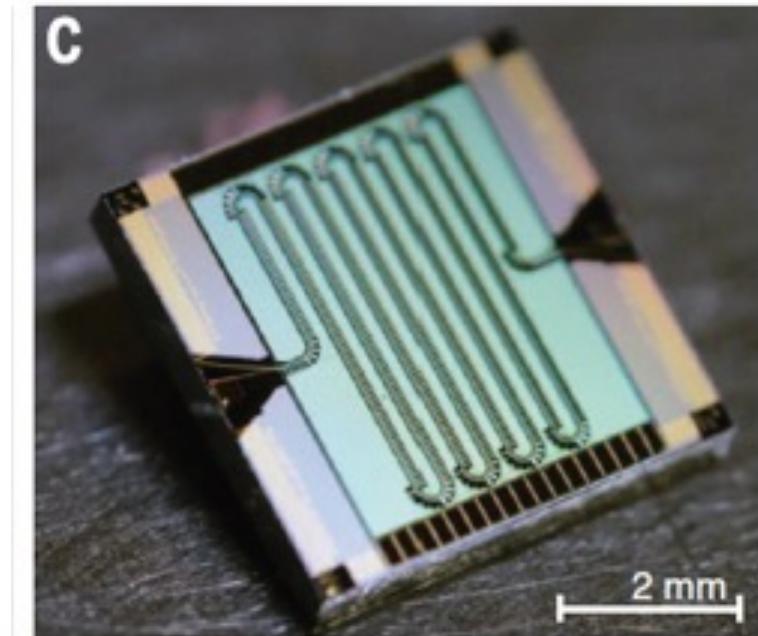
Fig. 1. Josephson traveling-wave parametric amplifier. (A) Circuit diagram. The JTWPA is implemented as a nonlinear lumped-element transmission line; one unit cell consists of a Josephson junction with critical current $I_0 = 4.6 \mu\text{A}$ and intrinsic capacitance $C_J = 55 \text{ fF}$ with a capacitive shunt to ground $C = 45 \text{ fF}$. Every third unit cell includes a lumped-element resonator designed with capacitance $C_r = 6 \text{ pF}$ and inductance $L_r = 120 \text{ pH}$, with coupling strength set by a capacitor $C_c = 20 \text{ fF}$. The value of C in the resonator-loaded cell is reduced to compensate for the addition of C_c . (B) False-color optical micrograph. The coloring corresponds to the inset in (A), with the lower metal layer shown in gray. (C) Photograph of a 2037 junction JTWPA. The line is meandered several times on the 5 mm by 5 mm chip to achieve the desired amplifier gain.

Macklin et al., Science 350, 307 (2015).

Summary of microwave amplifier for qu-technologies

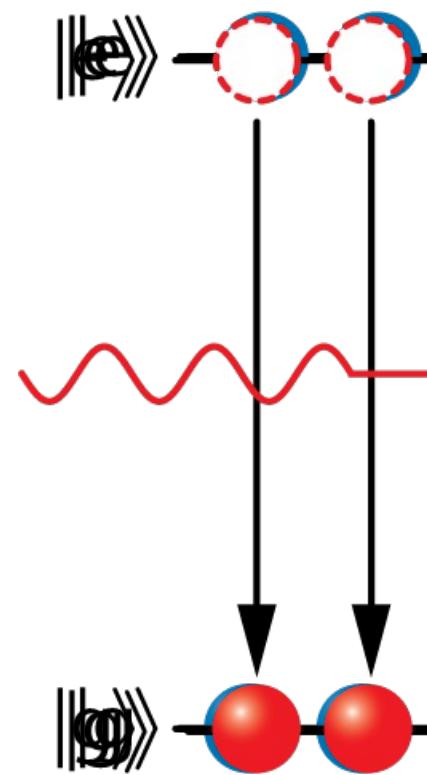
- Necessary to operate at millikelvin without adding noise without dissipating too much power for the fridge
- **The only solution (to date):**
Josephson parametric amplifier (JPA)

Is it really true?

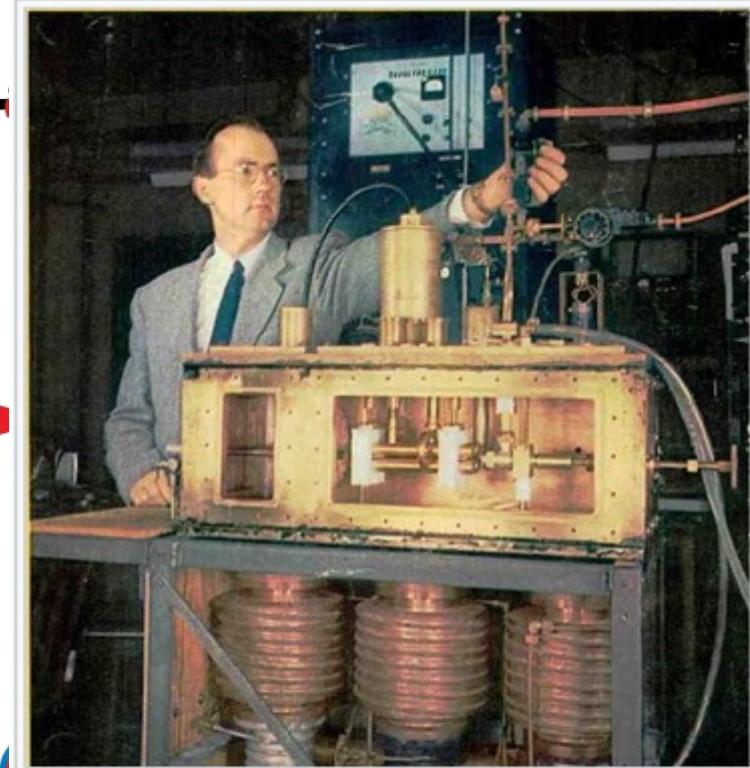


Stimulated emission

- Maser (**Microwave amplification by stimulated emission of radiation**)
- Population inversion (negative temperature)
- Intensively studied in 1950's - 60's but **left behind** due to
- birth of laser
- low temperature operation (**useless for CLASSICAL information technologies**)

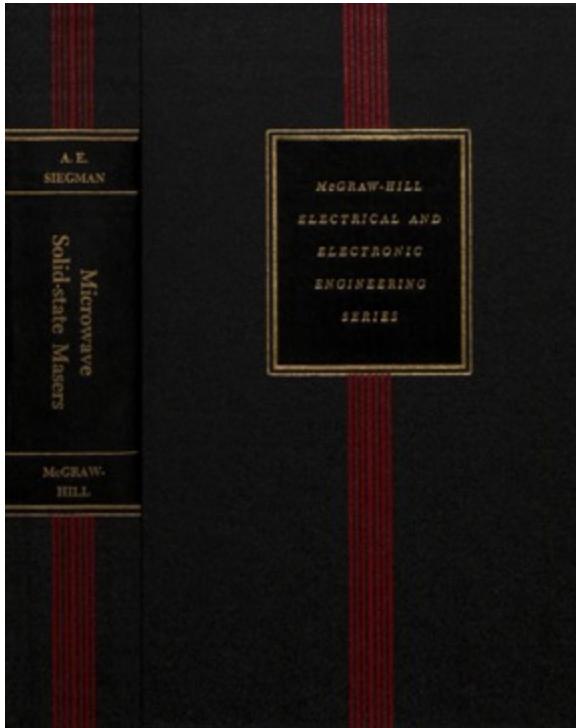


Charles Townes (Nobel prize in 1964)



First prototype ammonia maser and inventor Charles H. Townes. The ammonia nozzle is at left in the box, the four brass rods at center is the **quadrupole** state selector, and the resonant cavity is at right. The 24 GHz microwaves exit through the vertical **waveguide** Townes is adjusting. At bottom are the vacuum pumps.

How maser can be good ?



Maser noise figures

The low-temperature operation of the maser is in part a virtue in that it is at least partly responsible for the major virtue of the maser amplifier—its extraordinarily low noise figure. In discussing ultralow noise figures, it is most convenient to discuss not the noise figure F as conventionally defined, but rather the equivalent excess input noise temperature T_a of the amplifier. This noise temperature is related to the conventional noise figure F by the equation

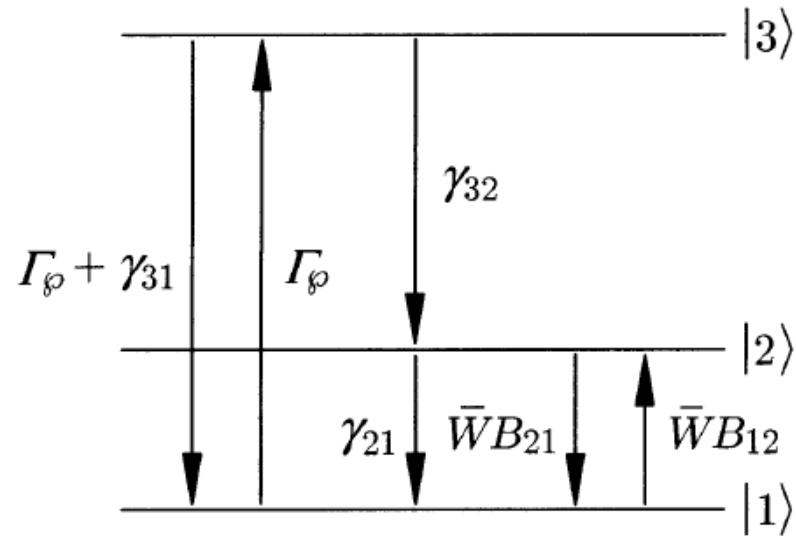
$$F = 1 + T_a/290^{\circ}\text{K} \quad (1-2-1)$$

Using this notation, typical masers have noise temperatures of the same order as their operating temperature, or about 4°K . The noise figure is then $F \sim 1.01$, or a few one-hundredths of a decibel.

Two reasons can be given for this extremely low noise figure. In the first place, the maser contains none of the usual noise-generating elements found in conventional amplifiers—no d-c currents, hot cathodes, or shot noise. The only noise source in the maser is thermal noise. Secondly,

Quantum-limited microwave amplifier possible by maser!?

Laser (Maser) principle



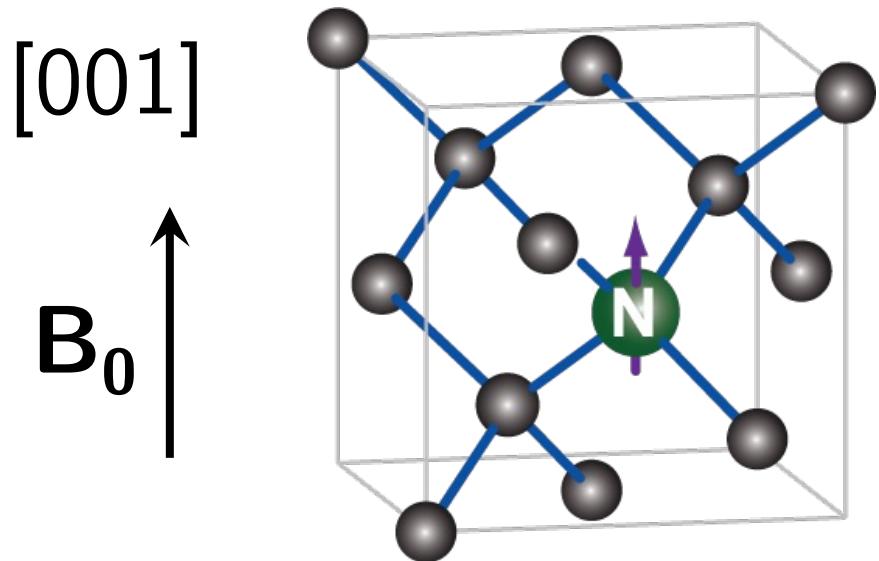
$$\dot{N}_1 = -\Gamma_p N_1 + \gamma_{21} N_2 + (\Gamma_p + \gamma_{31}) N_3 + \bar{W} B_{21} (N_2 - N_1),$$
$$\dot{N}_2 = -\gamma_{21} N_2 + \gamma_{32} N_3 - \bar{W} B_{21} (N_2 - N_1),$$
$$\dot{N}_3 = -(\Gamma_p + \gamma_{31} + \gamma_{32}) N_3 + \Gamma_p N_1.$$

[Carmichael, "Statistical methods in quantum optics 1" (1999)]

- N_i : atom (spin) population on level $|i\rangle$
- Γ_p : pump rate
- γ_{ij} : decay rate from $|i\rangle$ to $|j\rangle$
- $\bar{W}B_{21}$: stimulated emission (absorption) rate (\bar{W} : mode energy density, B_{21} : Einstein's coefficient)
- n : photons in the cavity, κ : cavity leak rate

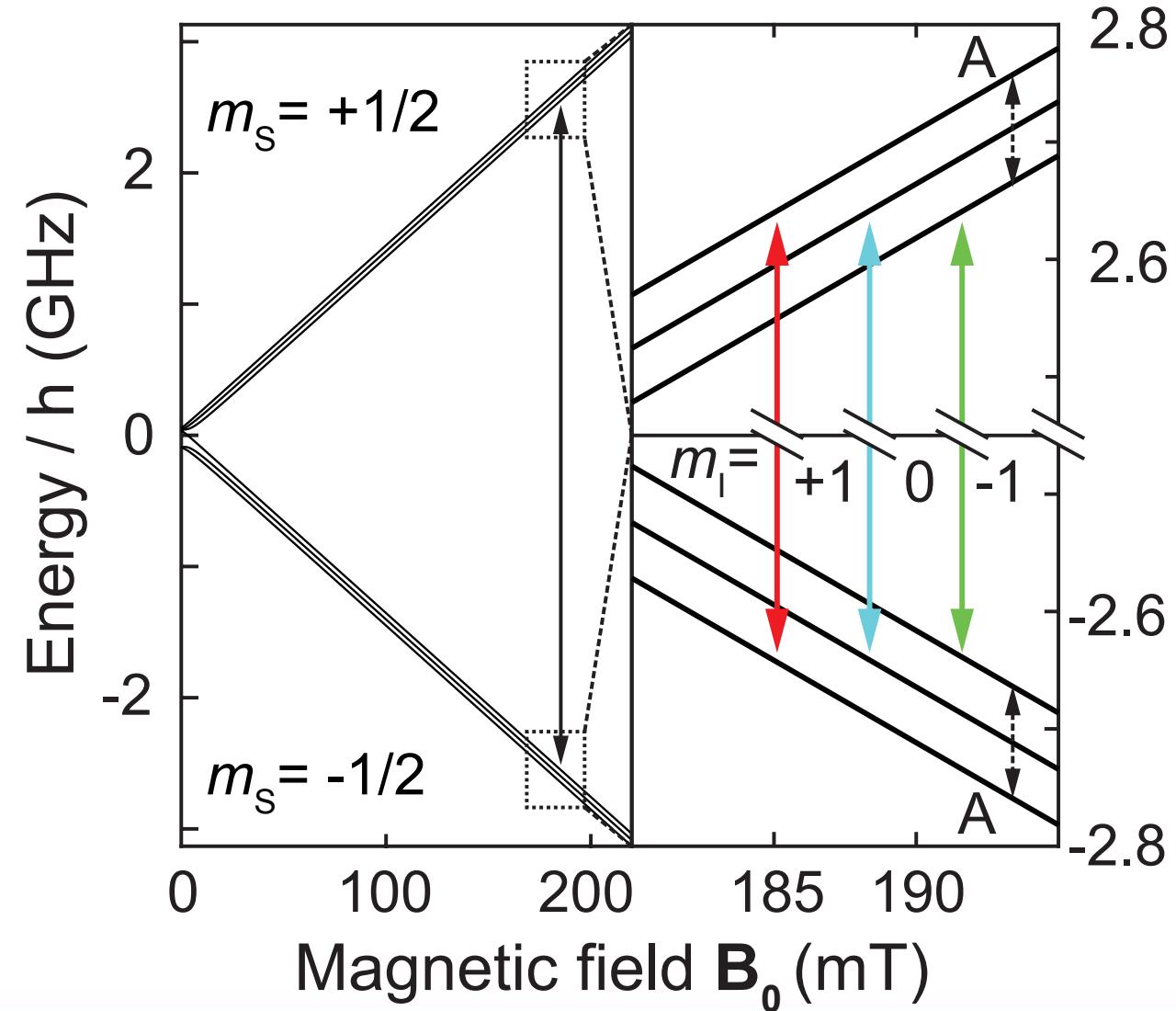
The spins: defect in diamond

Nitrogen impurity (**P1 center**)



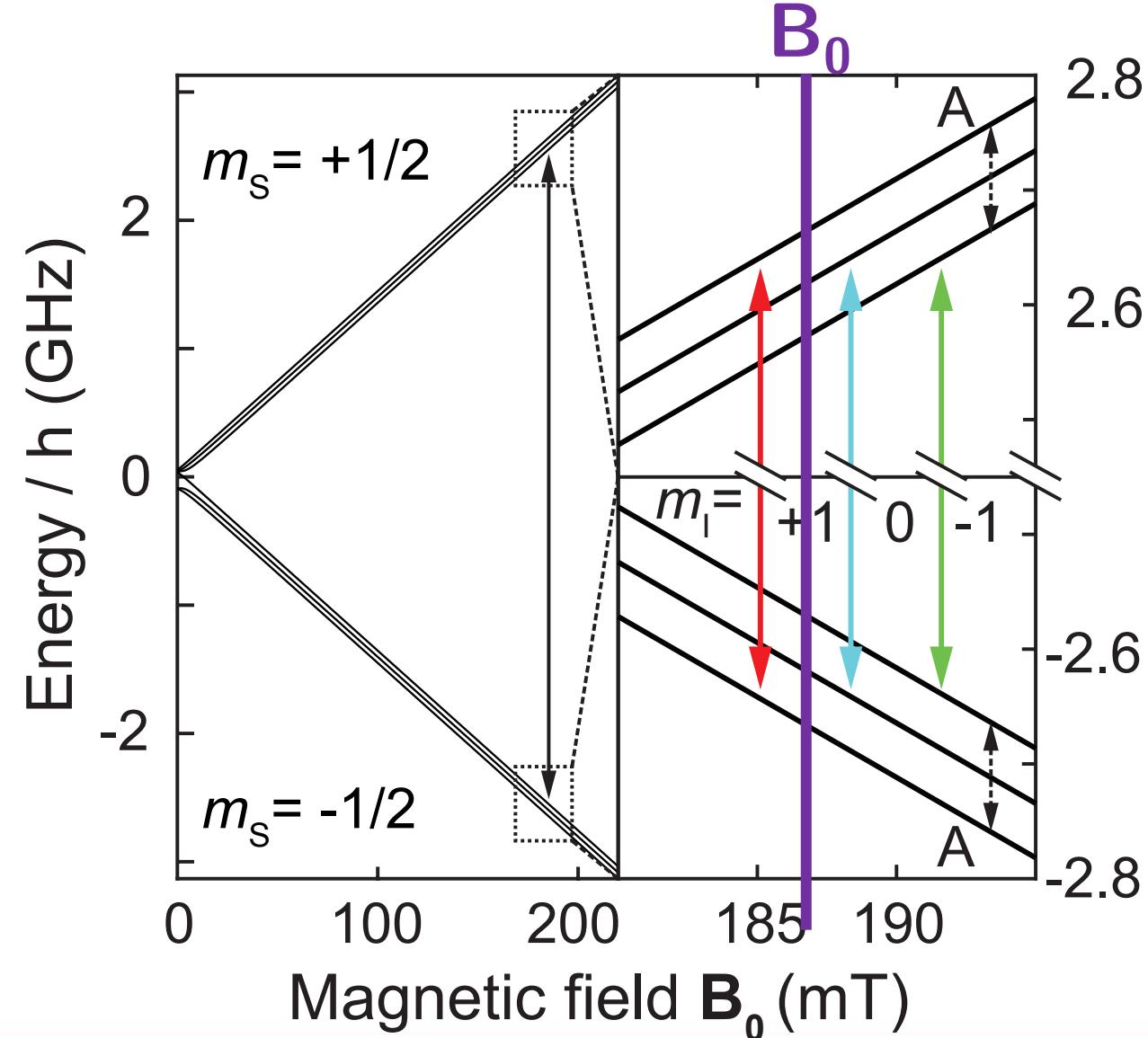
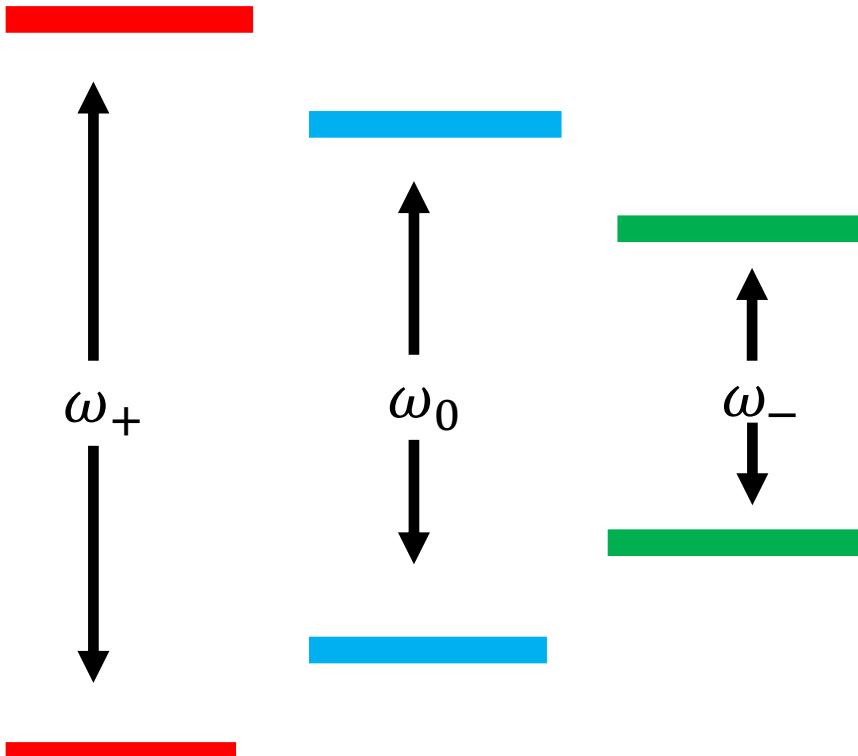
$$\mathbf{H}_{P1} = \gamma \vec{B} \cdot \vec{S} + \vec{S} \cdot \mathbf{A} \cdot \vec{I}$$

$$S = 1/2 \quad I = 1 \quad A_{\perp} = 114 \text{ MHz}$$
$$A_{||} = 81 \text{ MHz}$$



Three “evenly-spaced” 2-levels in P1 center

Nitrogen impurity (**P1 center**)



Four-spin flip-flop process in nitrogen centers

- Four-spin flip-flop process (cross-relaxation): two central spins' flips "up (down)"
= one spin on each satellite flips "down (up)"
- Measured to be ~ 10 ms
- Claimed could be used to generate inversion by **pumping central line** (reported, but data not shown)

Cross Relaxation Studies in Diamond*

P. P. SOROKIN, G. J. LASHER, AND I. L. GELLES

Research Laboratory, International Business Machines Corporation, Poughkeepsie, New York

(Received December 18, 1959)

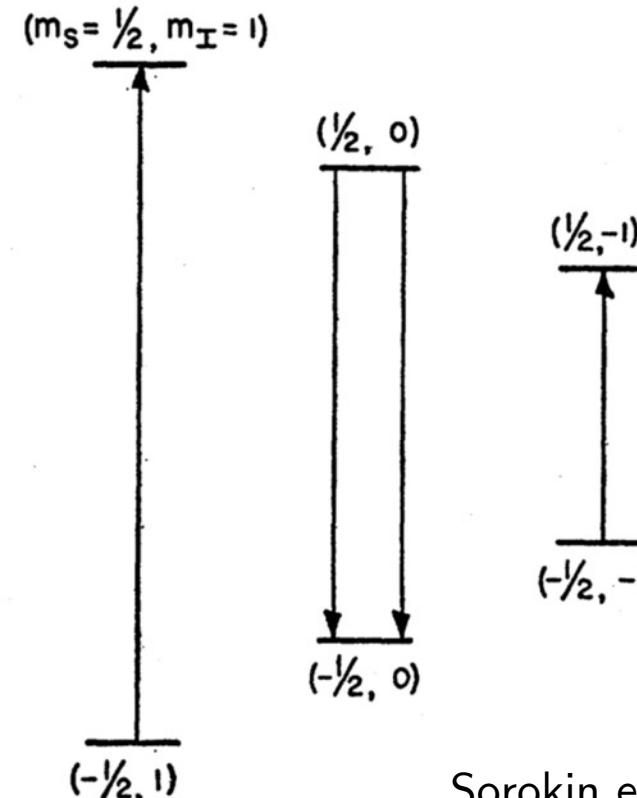
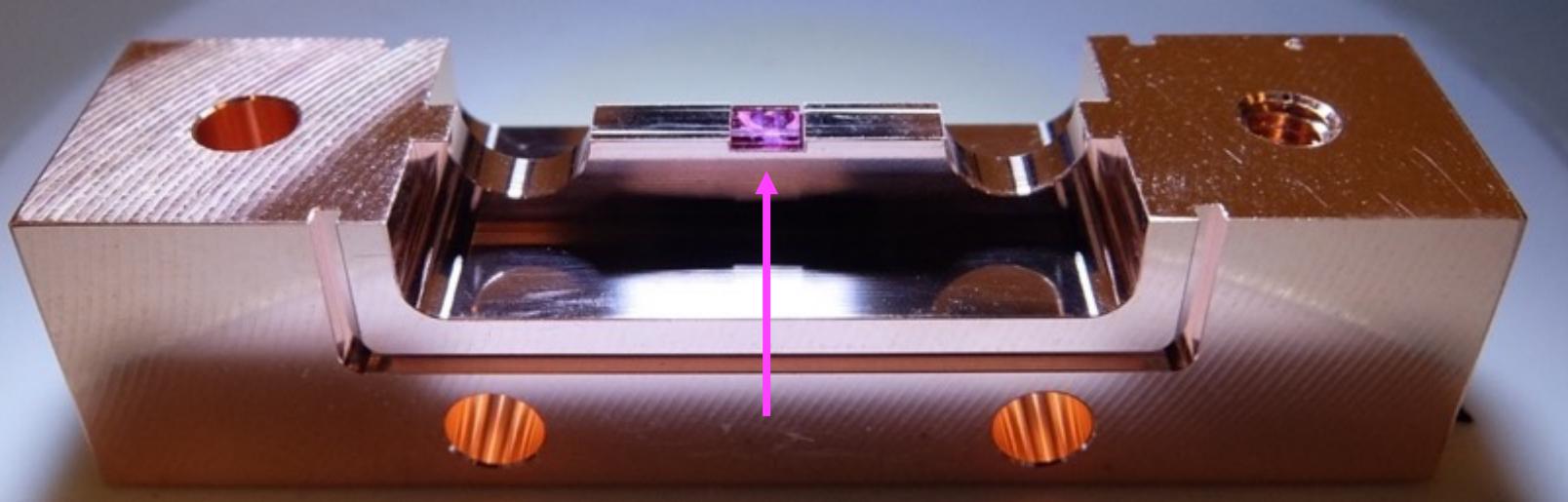
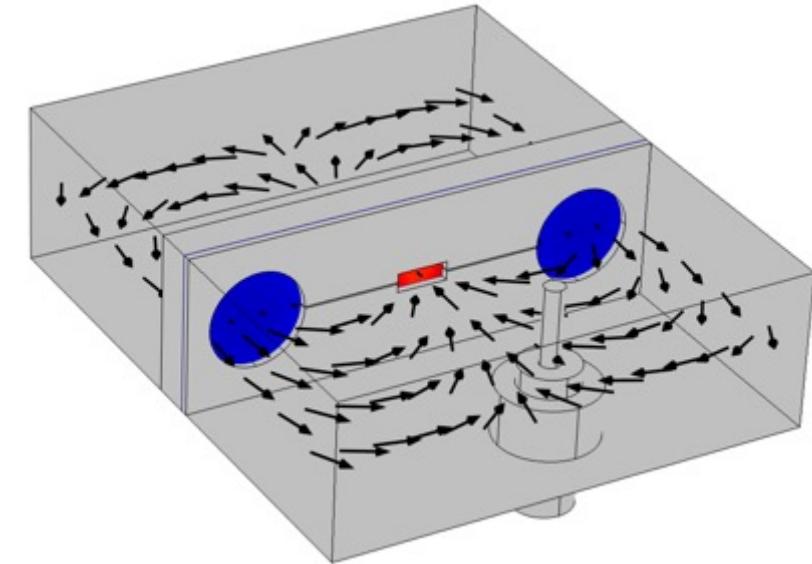
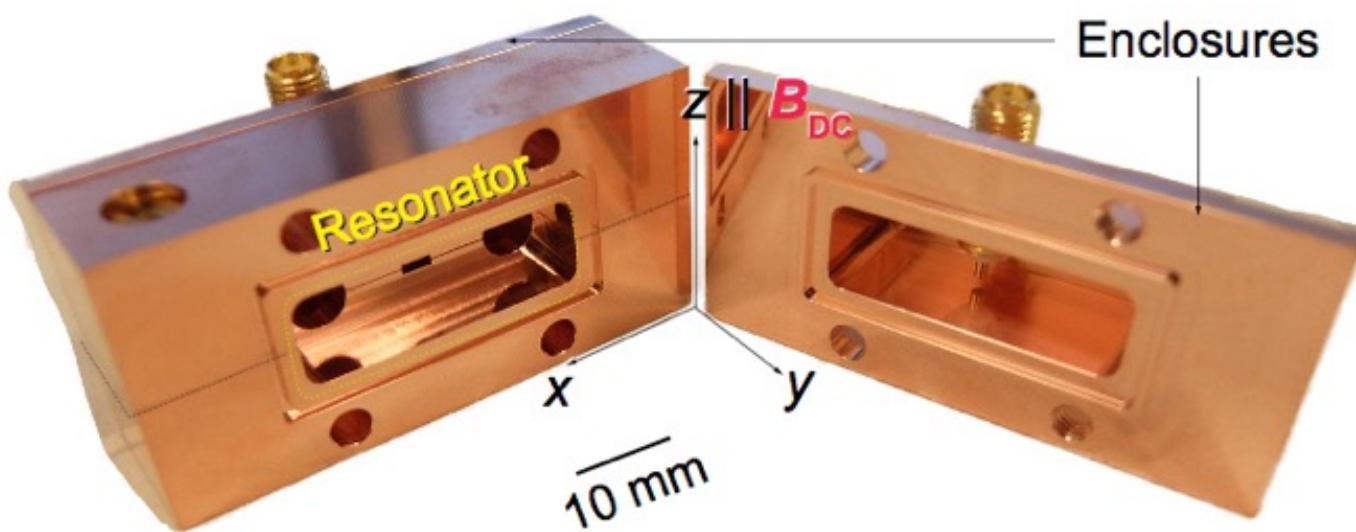


FIG. 1. Electron spin resonance energy level diagram for nitrogen centers in diamond ($H_0 \parallel [100]$) showing the four spin flip mechanism through which cross-relaxation proceeds. The hyperfine splitting Δ is 94.2 Mc/sec.

Sorokin et al., Phys. Rev. **118** (1960).

Loop-gap microwave resonator



Diamond (irrad. @QST)

Ball et al., Appl. Phys. Lett. **112** 204102 (2018)

$$\begin{aligned}f_r &= 6.25 \text{ GHz} \\Q_{\text{ext}} &= 750 \\Q_{\text{int}} &= 2500\end{aligned}$$



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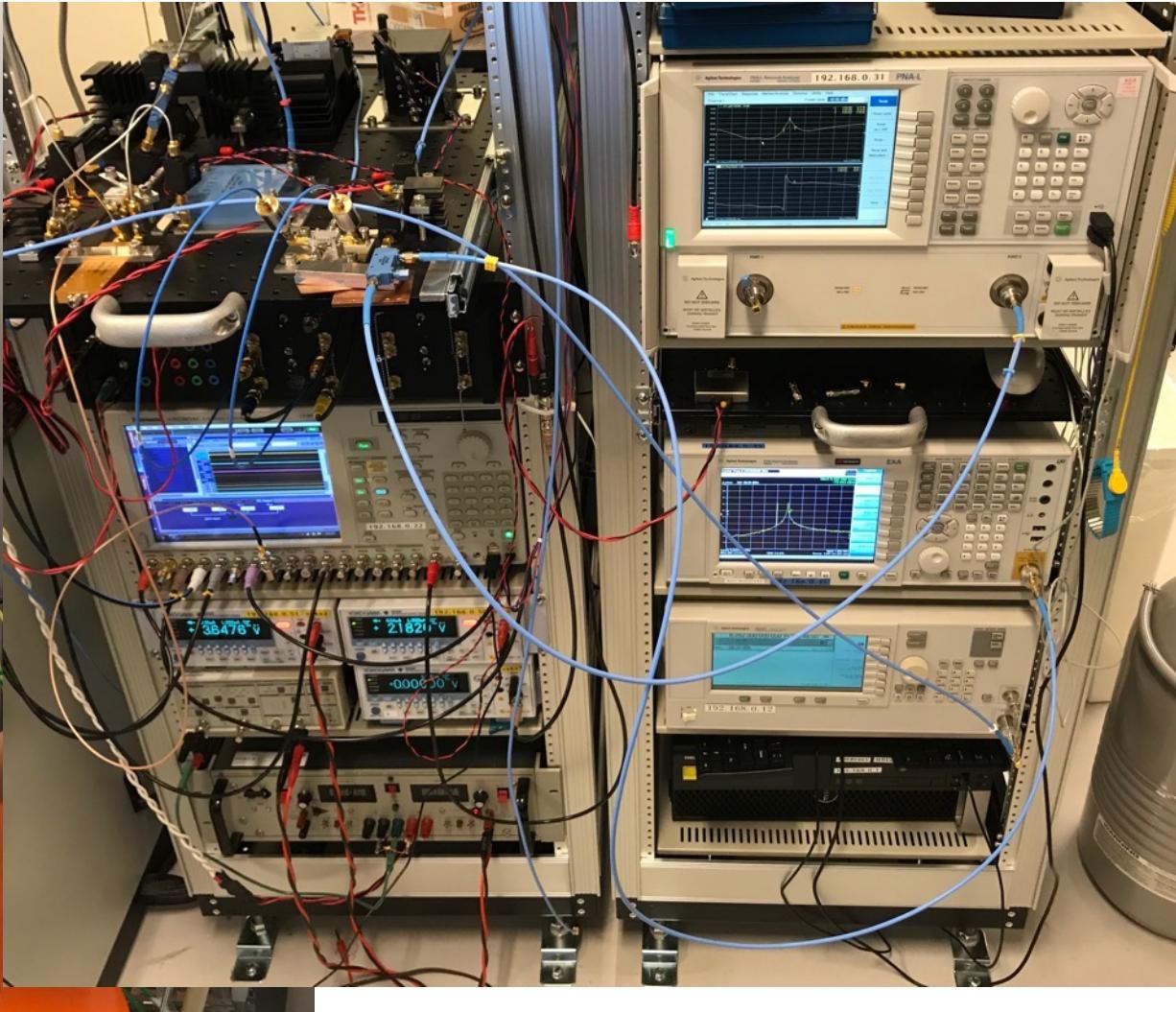
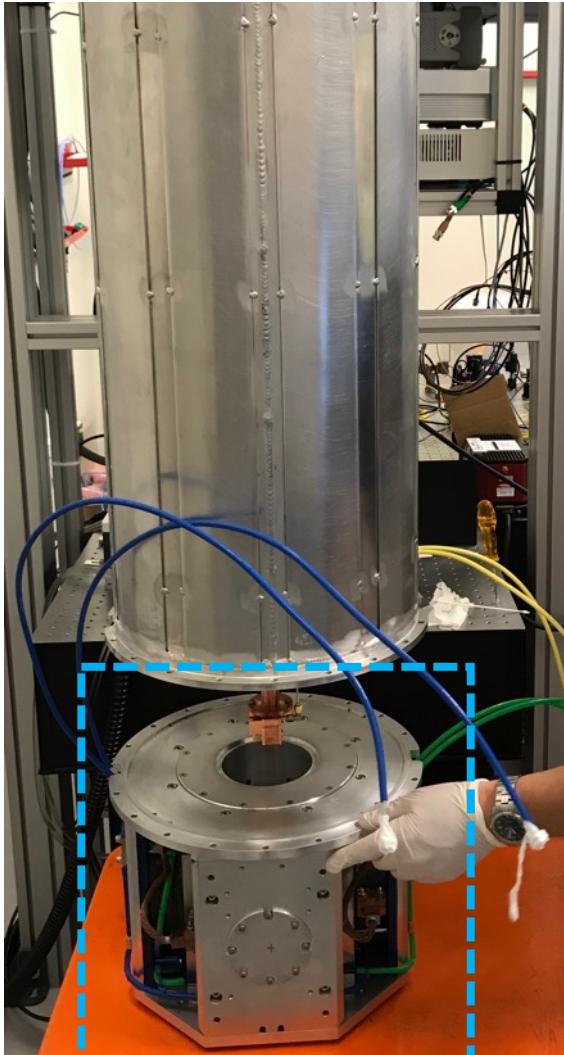
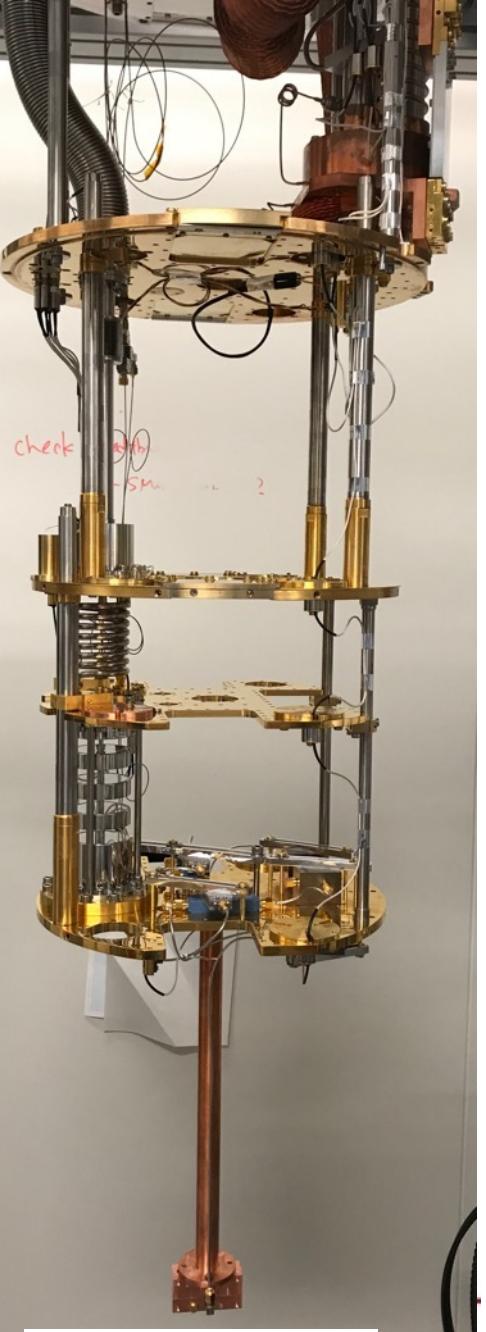
Yuimaru Kubo |

Setup photos

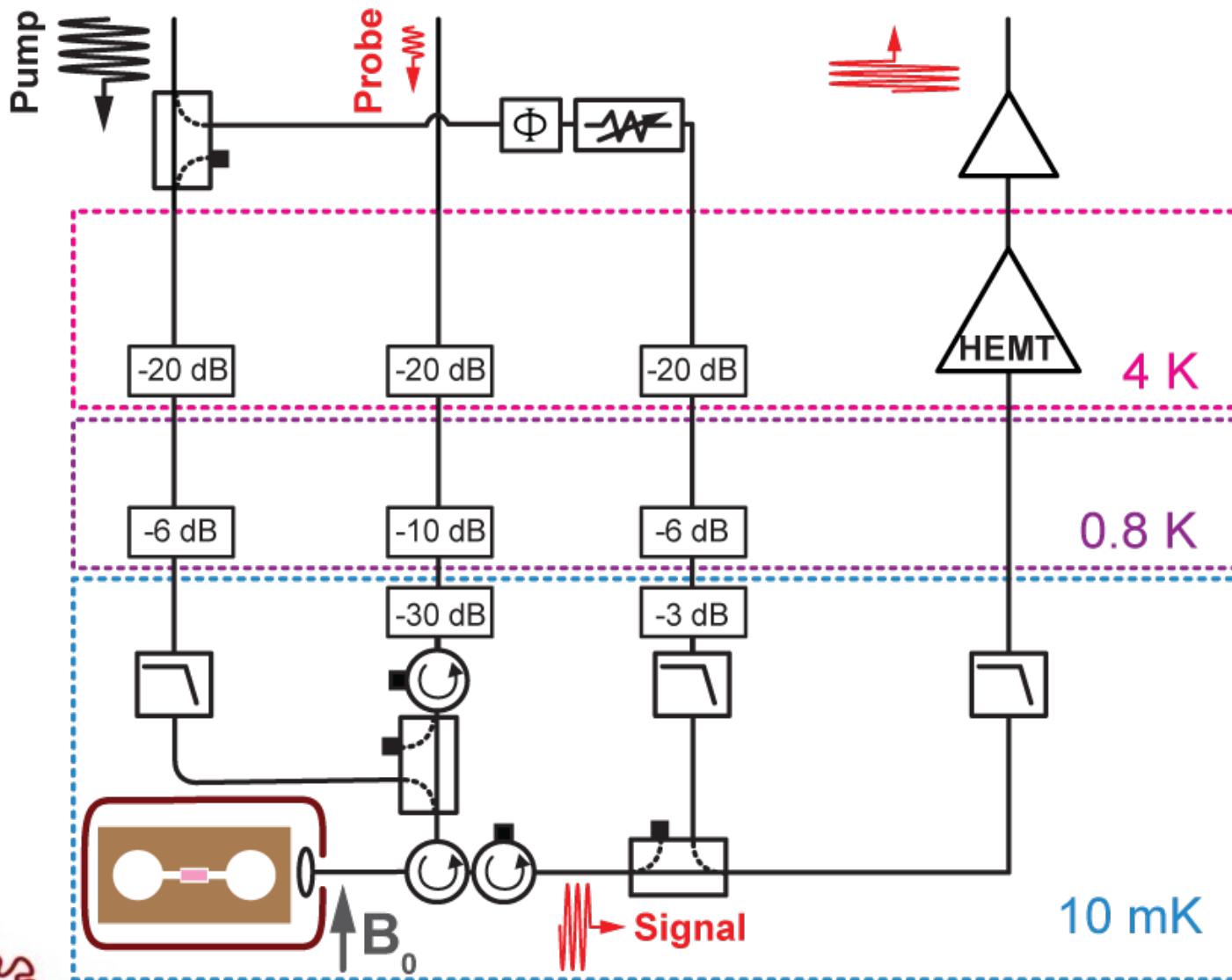
4K

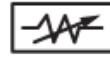
0.8 K

10 mK

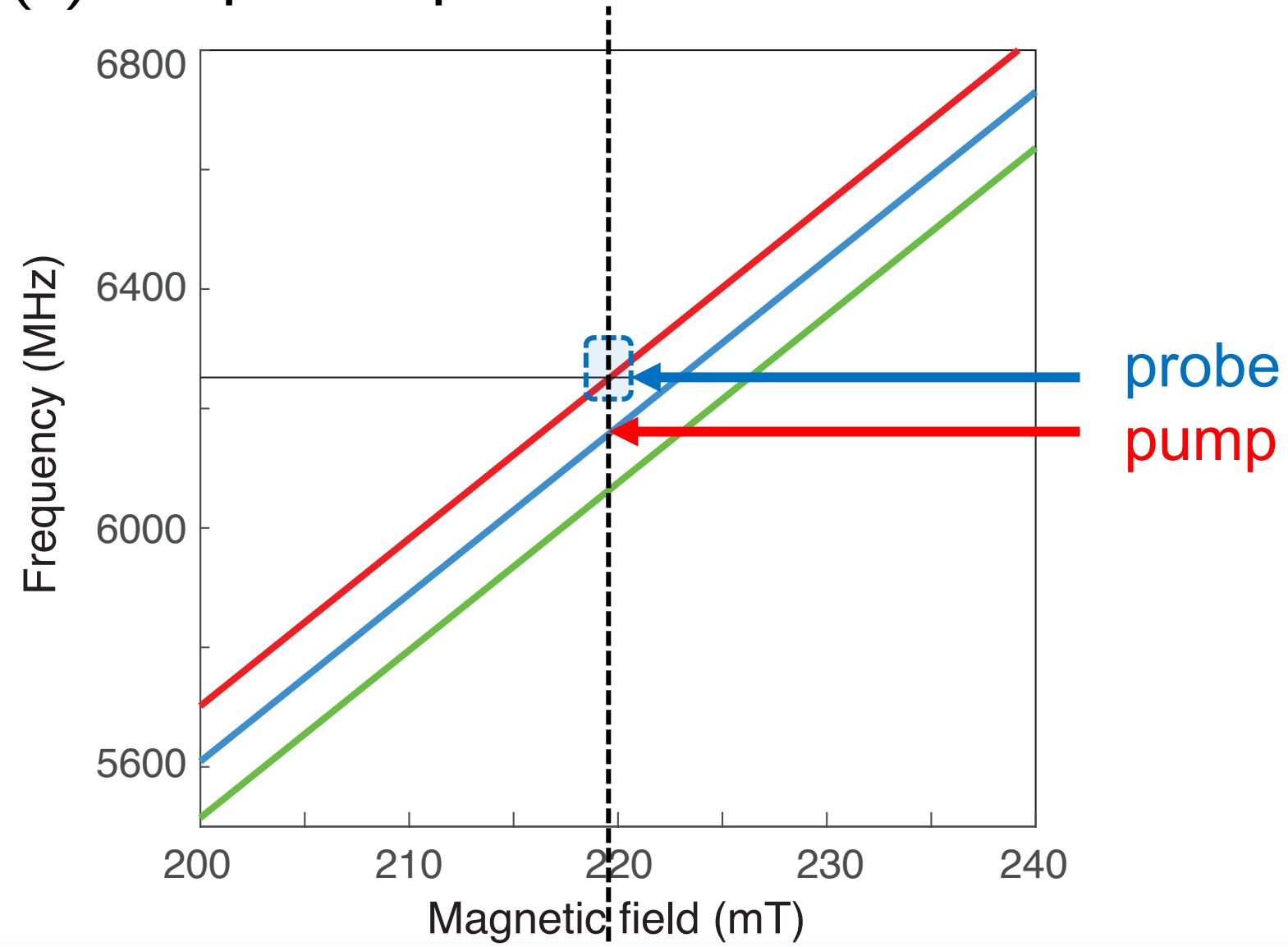
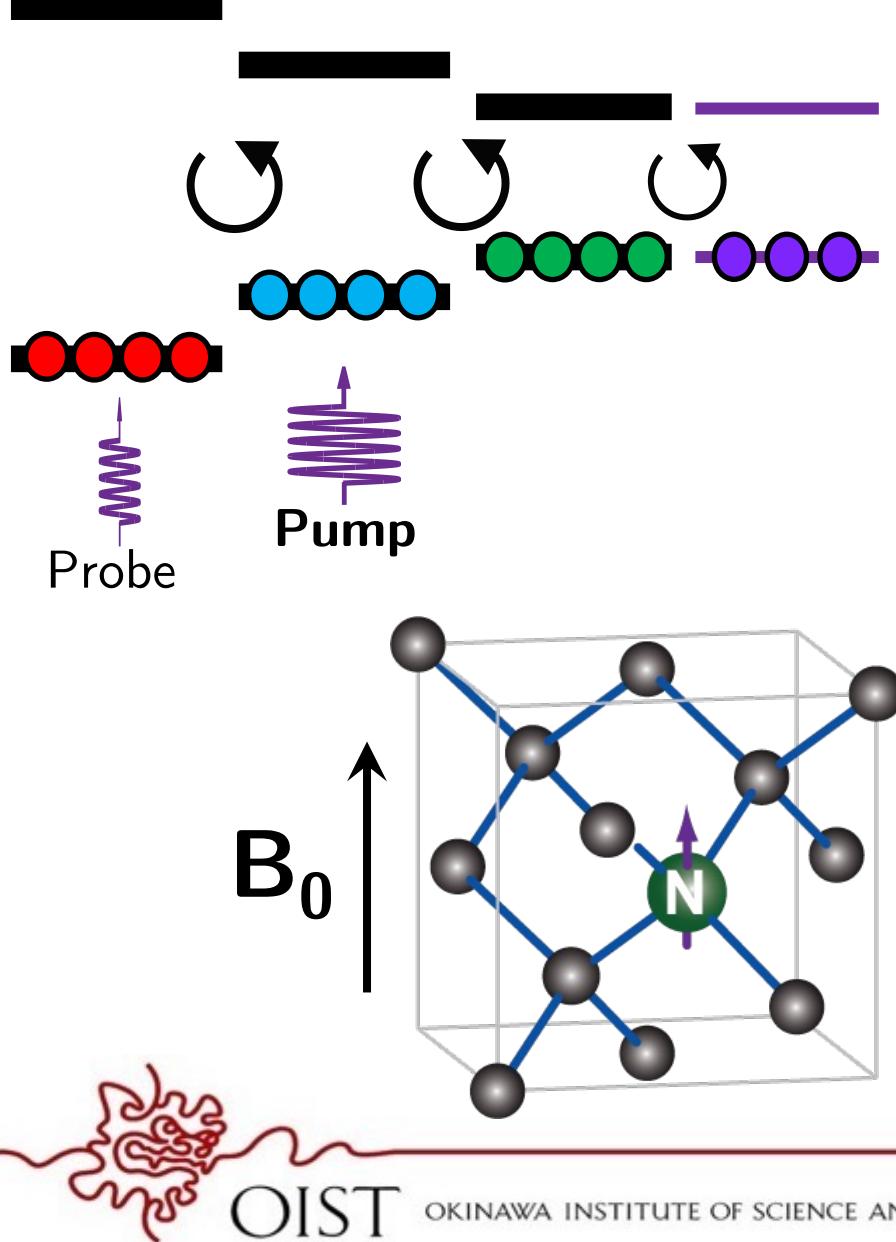


Setup



-  Directional coupler
-  Variable attenuator
-  Phase shifter
-  High electron mobility transistor amplifier
-  Circulator (isolator)
-  Low pass filter
-  Loop-gap resonator
-  Copper enclosure
-  Diamond

Mase(r) amplifier protocol

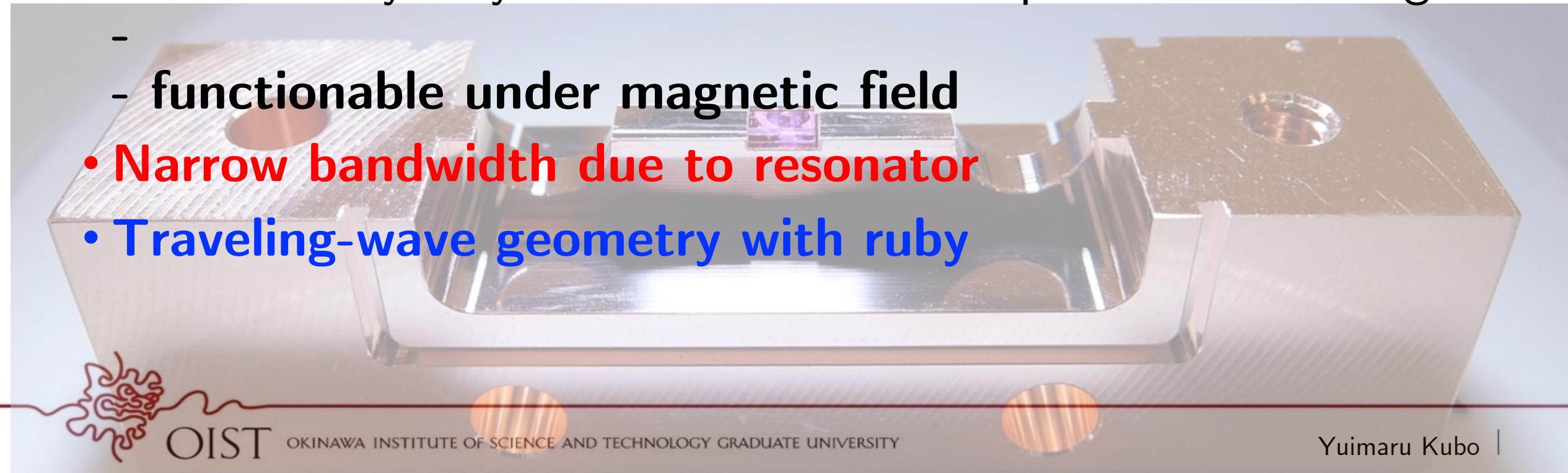


P1 maser amplifier

To be published...

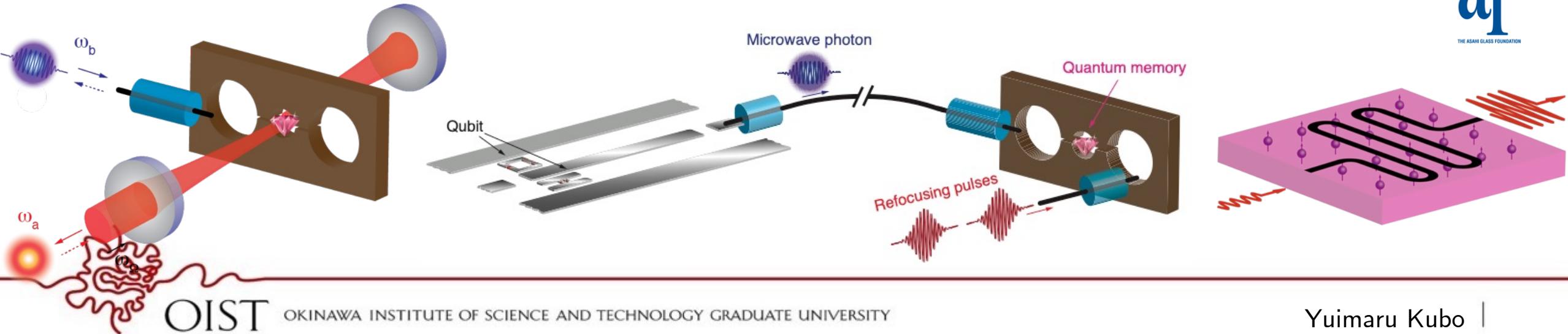
Summary and perspectives

- **Maser amplifier at millikelvin temperature** by means of P1 centers in diamond
- **(Near-)quantum-limited noise and huge dynamic range**
 - >Potentially very useful for microwave quantum technologies:
 -
 - **functionable under magnetic field**
- **Narrow bandwidth due to resonator**
- **Traveling-wave geometry with ruby**



On-going projects at OIST

- **Quantum microwave-optical photon transducer**
with spins in diamond
- **Spin ensemble quantum memory**
for superconducting qubits
- **Spin-based ultra-low noise microwave amplifier**



Acknowledgements

- Grants:



科 研 費
KAKENHI



Science of Hybrid
Quantum Systems



公益財団法人
住友財团

日本私立学校振興・共済事業団
Promotion and Mutual Aid Corporation for Private Schools of Japan

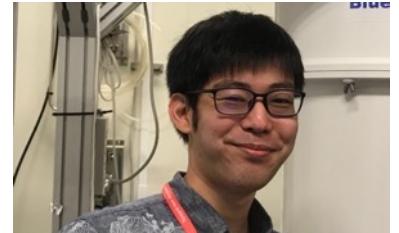
af
THE ASAHI GLASS FOUNDATION



公益財団法人 平和中島財團
+++ HEIWA NAKAJIMA FOUNDATION HOMEPAGE +++

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(ruby maser,
Now @NPL, UK)

Jason Ball
(diamond maser
Now @Tabor)

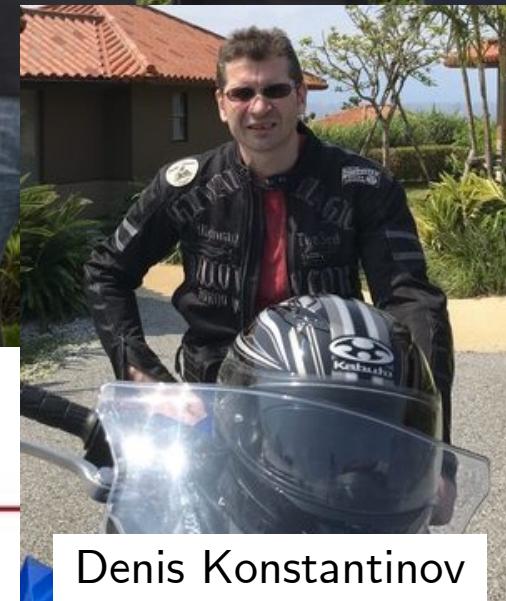


Tatsuki Hamamoto
(transducer)



Morihiro Ohta
(diamond maser)

Denis Konstantinov
(diamond maser)

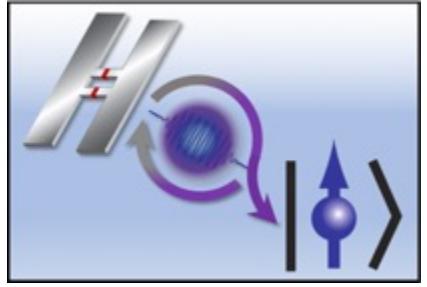


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Hybrid Quantum Device Team

